Requirements for a next generation nuclear data format

OECD/NEA/WPEC SubGroup 38

(Dated: April 1, 2015)

This document attempts to compile the requirements for the top-levels of a hierarchical arrangement of nuclear data such as is found in the ENDF format. This set of requirements will be used to guide the development of a new set of formats to replace the legacy ENDF format.

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I. INTRODUCTION

It was realized as far back as the Manhattan project that collections of nuclear cross-section data were needed. These collections evolved into the “Barn Book” (Hughes [1955]) first sponsored by the United States Atomic Energy Commission. By 1963, there were many nuclear data libraries such as the United Kingdom Nuclear Data Library (UKNDL) (Parker [1963]) from Ken Parker at the Atomic Weapons Research Establishment, Atomic Energy Authority in Aldermaston, UK; the fast reactor data library from Joe Schmidt at the Institute for Neutron Physics and Reactor Technology, Nuclear Research Center, Karlsruhe, Germany; the NDA library from Herb Goldstein at Nuclear Development Associates, in New York; and the Evaluated Nuclear Data Library (ENDL) from Bob Howerton at the Lawrence Radiation Laboratory in Livermore, California. Each laboratory developed its own storage and retrieval schemes for its data and in many cases libraries were hard-wired into simulation codes. As a result, reactor designers and other data users could not use new cross-section data, even though in some cases the data were available for five or more years (Goldstein [1968]). Furthermore, dissimilarities in the internal formats of each lab kept data users from reconciling differences in calculated values for the same reactor configurations.

There was a need for a common mechanism for intercomparison between these systems. Following a discussion between Henry Honeck (Brookhaven National Laboratory), Al Henry (Westinghouse) and George Joanou (General Atomics) at the Colony Restaurant in Washington, D.C, the Reactor Mathematics and Computation (RMC) Division of the American Nuclear Society (ANS) was requested to sponsor two meetings to discuss a plan to develop this mechanism. Honeck, the chairman of the Division’s sub-committee on Evaluated Nuclear Data Files, held these meetings. This effort culminated in a meeting of 18 representatives from 15 US laboratories in New York City on July 19, 1963 to review cross section libraries and discuss means for exchanging these libraries. A sub-committee was appointed to meet in Hanford on September 18-20, 1963 to examine library formats in more detail. The conclusions of these meetings were:

- There was a need for a standard format for evaluated nuclear data.
- The format should be as flexible as possible so that existing libraries could be translated into the standard format and so that the format can be extended to meet future needs.
- This standard format would be the link between a data library and the processing codes.
- It was also suggested that a center be created and tasked with the development and maintenance of the new format called the Evaluated Nuclear Data File (ENDF). This center would also collect and distribute data.

A preliminary version of the ENDF format was sent for review and comment. At the final meeting at Brookhaven on May 4-5, 1964, the 22 attendees discussed changes to ENDF and settled on a final version. The description of the system (referred to as ENDF/A [1]) was documented in the report BNL-8381 (Honeck, 1965). The initial ENDF/A library contained an updated version of the UKNDL library and evaluated data from a number of different laboratories. As ENDF/A did not contain full evaluations, there was also a need for evaluated nuclear data to be used for reactor design calculations. The description of this system (referred to as ENDF/B) grew out of ENDF/A and was documented in the report BNL-50066 (Honeck, 1966). Whereas the format of ENDF/A was flexible, allowing data centers to produce and accept data in a variety of representations, the format of ENDF/B had to be simple and mathematically rigorous to facilitate the development of the supporting infrastructures including data processing, integration and plotting.

Nearly 50 years later, we are revisiting the format and its specifications in order to modernize it. The fundamental need for data has not changed in this time, but the computational tools we have available are much more advanced and our physics understanding has advanced as well. The kinds of data we need to store have also grown markedly. In November 2012, a Working Party on Evaluation Cooperation SubGroup, WPEC-SG38, was formed (WPEC Subgroup 38 [2012]) to coordinate the modernization of the ENDF format and supporting infrastructure. At that meeting, the following seven tasks were organized (WPEC Subgroup 38 [2013a,b]):

<table>
<thead>
<tr>
<th>Requirement 1: WPEC Subgroup 38 tasks</th>
</tr>
</thead>
<tbody>
<tr>
<td>1.1 Low level data containers</td>
</tr>
<tr>
<td>1.2 Top level reaction hierarchy</td>
</tr>
<tr>
<td>1.3 Particle property hierarchy</td>
</tr>
<tr>
<td>1.4 Visualization, manipulation and processing tools</td>
</tr>
</tbody>
</table>

ENDF/A referred to both a format and to a library stored in that format.
This document attempts to compile the requirements for a hierarchical arrangement of nuclear data, such as is found in the ENDF format, that addresses WPEC Subgroup 38 Tasks 1.1, 1.2 and 1.3. This set of requirements will be used to guide the development of the specifications for a new structure to replace the legacy ENDF format. The requirements do not define the format used to store the data. Instead, the structure can be stored in any nested hierarchical meta-language, ensuring that future users and developers of nuclear data will be able to store data in whatever medium is available in the future.

In this document, we will commonly refer to the Generalized Nuclear Data (GND) format and the Fudge code system, two projects initiated at LLNL (Mattoon, 2012; Pruet, 2006). GND is a hierarchical nuclear data format that is the prototype for the system WPEC-SG38 is creating. Fudge is the first code framework that can interpret and manipulate GND formatted data. Neither the GND examples nor the names or arrangements of data in the figures of this document are “set in stone”; all are expected to evolve in the process of developing the new format and infrastructure.

The authors of this document have discussed the requirements broadly with members of the nuclear data community and attempted to capture the ideas and needs from a wide range of applications and perspectives. With a high degree of confidence we can say we did not capture all of them. One can surely find one more case of “what about ...?”. However, we believe this document reflects a comprehensive view of our communities’ “best practices” and provides the map to an ENDF modernization that will stand the test of time. There are significant risks to building anything by committee: the joke goes that an elephant is a mouse built by committee. History will judge the success or failure of this effort by the adoption and longevity of its product.

A. Scope of data to support

A nuclear reaction data library evaluation describes an incident particle (called a projectile) impinging on target. The target may be a single atom or atomic nucleus, or a collection of atoms with which the projectile reacts. The projectile may be the traditional n, p, d, t, 3He, α, γ, or e− or any other single (composite) particle (e.g., 12C, muon or pion). The ultimate goal for a new structure is that it must support storing self-consistent, complete evaluations along with their derived data. The derived data includes things like transfer matrices and energy depositions along with the parameters (e.g., group boundaries and fluxes) used in their derivation. Storing derived data in a standardized structure facilitates inter-laboratory comparison of processed data. Additionally, the format must support uncertainties, generally given as covariance tables, on all these quantities.

The data to be stored are typically a balance between what an evaluator can provide and what a particular application needs. Therefore, it is useful to look at the most common use cases:

- **Particle Transport:** For transport, cross sections for all reactions that are energetically possible over a given range of incident energy $E$ must be stored, along with a list of outgoing reaction products with the products’ multiplicities and probabilities for outgoing energies and angles. Data may be parametric (for example, Watt spectra for storing energy distributions or resonance parameters for storing both cross sections and distributions), or they may be given in tabular form. Some of these data are temperature-dependent and the temperature must be denoted. Additionally, different solution methods (deterministic or Monte Carlo transport) have different derived data requirements:
  - **Deterministic transport:** In addition to cross sections, deterministic transport codes require transfer matrices (which store the double-differential cross section for each reaction product as a Legendre expansion where each term is averaged over incident and outgoing energy ranges), as well as energy and momentum deposition cross sections. Although these quantities are derived from the cross section, multiplicity and distribution data, the transfer matrices in particular are computationally intensive to calculate, so the new structure needs to be capable of storing them for re-use and exchange.
  - **Monte Carlo transport:** A single cross section for a given incident energy is not always sufficient for Monte Carlo simulations. In particular, the rapidly fluctuating cross section in the unresolved resonance region can often be better described as a probability distribution of cross sections. The new data format therefore needs to support storing cross section probability tables. The new format may also need to support storing other cumulative distribution functions (CDFs) for sampling outgoing particles.

- **Transmutation/Isotope Burn-Up:** For isotopic production and depletion, cross sections and optionally outgoing spectra for chosen reactions that are energetically possible over a given range of incident energy $E$ are needed. In addition, the decay half-lives and product branching ratios must
be known for all nuclei to enable the time-dependent calculation of the isotope inventory.

- **Astrophysical network calculation:** In an astrophysical network, one needs isotopic accretion and depletion cross sections. These data may be averaged over Maxwellian neutron spectra to simulate the neutron flux in an astrophysical environment. Astrophysical networks also involve reactions with charged particle projectiles – not just neutrons – and may require detailed knowledge of charged-particle interactions in plasmas.

- **Web Retrieval:** For archival, there are no completeness requirement as the data will not be used in applications. The data will most likely only be visualized. That said, visualization requires a pointwise representation for many data types.

- **Uncertainty Quantification (UQ):** UQ applications encompass all of the above use cases. The defining difference is the need to also specify uncertainties/covariance on aspects of data. One can then either generate statistical realizations of these data (if one is adopting a Monte Carlo approach to UQ) or one can use the “Sandwich Formula” (if one is using a deterministic approach).

For all use cases, data will need to be documented. Therefore we will need facilities for documenting data clearly, concisely and as machine and human readable as possible. In addition, version information for the format, the documentation, the evaluation itself, the codes used in evaluation, etc. all need to be stored.

**B. How to use these requirements**

This document is a list of the requirements for the new format. It is not the specifications of the format. At times the details of the requirements may constrain the actual specifications so much that the requirements may seem to be specifications. Other times the requirements will be broad and may be left to interpretation or may have several possible implementations.

In this document, requirements are called out and uniquely numbered so that they can be clearly referenced in later work. A requirements list is formatted as follows:

**Requirement 2: An example**

2.1 Don’t be evil
2.2 Respect the user
2.3 Let the evaluator express themselves clearly

In the process of developing these requirements, we have had discussions\(^2\) among ourselves and with other members of the nuclear science community. We have captured many of these discussions in this text so that future users of the data and formats can understand our reasoning. Much of this discussion has been already integrated into the requirements. However, there were other questions that required further discussion and these are formatted as follows:

<table>
<thead>
<tr>
<th>Discussion point:</th>
</tr>
</thead>
<tbody>
<tr>
<td>Users didn’t get this point.</td>
</tr>
<tr>
<td>Resolution:</td>
</tr>
<tr>
<td>We added an example to clarify it.</td>
</tr>
</tbody>
</table>

Also, during the effort to create this and its companion documents, it became clear to us (the authors and other contributors) that we often do not agree on the meaning of a concept. Given that the ENDF format manual (Trkov 2009) is several hundred pages long and encapsulates decades of accumulated knowledge, it is no surprise that we are all not intimately familiar with all of its intricacies. Therefore, we provide a glossary in appendix C.

We also tend to use XML in examples and to denote elements in the hierarchy. This is only a matter of convenience as the data should be serializable in any nested hierarchical meta-language (e.g., HDF5, ROOT, JSON, Python classes). We denote major nodes/elements in the XML element-like notation (e.g., `<element>`) and attributes of these nodes/elements without the brackets (e.g., `attribute`).

Finally, we note that as we list and discuss requirements, we encounter suggestions for test cases for either data or supporting codes. This are marked up as follows: **SUGGESTED TEST:**. This is an example test.

**C. Main requirements**

At the highest level, there are some overarching goals that we believe a new format should achieve. These goals represent the most broadly held beliefs that guide all subsequent requirements. The main goals are:

**Requirement 3: Main**

3.1 The hierarchy should reflect our understanding of nuclear reactions and decays, clearly and uniquely specifying all such data.
3.2 It should support storing multiple representations of these quantities simultaneously, for

\(^2\) And they were sometimes contentious
Some of these goals may seem contradictory. For example, allowing multiple representations while at the same time eliminating redundancy appear to be conflicting goals. However, derived data are not redundant: they reflect the choices of the processor and needs of particular applications (e.g., group boundaries and fluxes) in addition to the original evaluated data.

It is up to each evaluated data project to determine specific requirements, such as the completeness, of the data to be stored in a particular library. Similarly, it is up to the data processor to decide how to process the data, but the resulting processed data need to be stored in a common structure to facilitate exchange and comparison. Crafting a structure that is capable of balancing all of the goals in requirements 3 is the task at hand.

D. Hierarchal structures

As implied in the first requirement for the new format, data shall be stored in a hierarchal structure. The design choices made in the development of the original ENDF format enabled the community to shoe horn data onto punch cards at the expense of obfuscating even the simplest things such as determining when one floating point number ends and another begins. The inability to read a data file without a deep understanding of the ENDF format leads to issues with quality assurance. Nevertheless, there is a common sense arrangement of data that reflects both our understanding of particle transport and reaction physics. Capturing and encoding that arrangement within the formal structure will vastly improve the usability of the data for those with little knowledge of our formats. This is an obvious benefit to the community at large, enabling more people to work more effectively and with greater assurance that they are correctly using the correct and best data for their applications.

Hierarchal storage defines itself by means of nodes existing as parents, siblings and children spanning logical levels. We require a means of identifying these nodes and do so by providing them names. These names are, to some degree, arbitrary and thus likely to be contentious. Care should be taken to ensure that names align themselves with generally accepted definitions. At the higher levels, these names should convey the physical essence of the data, for example crossSection or emissionDistribution. At the lower levels, they should convey the type of data being stored, for example points defining a piecewise continuous curve or vectors defining bands in a sparse array.

There are two significant questions regarding the naming of nodes: the choice of language and the character set used to encode them. These are pragmatic choices. The character sets available across storage systems vary greatly. Choosing a restricted set will enable a broader adoption of this structure into working implementations. As many of these storage systems allow only ASCII characters, or only a limited set thereof, the allowed language and character set need to be clearly defined.

The format must store data and information qualifying the data. The qualifying information are typically known as meta-data. There are values that are clearly data; for example cross sections, resonance parameters or emission distributions. There are also values that are clearly meta-data; for example the number of point-wise energy/cross section pairs, the units for values or the frame of reference for a particle emission. Hierarchal storage systems typically allow attribute names/values to be associated with a node. In general, the format should store meta-data as attributes to nodes and data within the node itself.

At the bottom of such structure are the actual data. These are general numeric—real, integer, complex—values with documentation, annotation or meta-data provided as text. At some point these values must be laid out sequentially in a storage system; that is, they must be written into a file for archival storage or exchange. At this lowest level, it is vital that the beginnings and ends of each component of the data be clearly delineated and that the necessary mathematical transformations and inherent physical meaning of the data be defined by the structure containing these data. Anyone who has ever struggled to determine where the data they need are stored inside a large data file will understand just how important this requirement is.

Done well, complexity reveals itself nothing more than elegant simplicity. Herein each data structure should itself consist of its basic building blocks: well-defined, well-formed, universal in specific application, and built from such lower level blocks as needed down to some final atomic component. Humans are by nature organizational animals, seeking to find order in chaos.

Requirements related to the hierarchical structure are:

**Requirement 4: Hierarchal structures**

4.1 Data shall be stored in a hierarchal structure
E. Complications

As we develop requirements for a new nuclear data hierarchy, we naturally encounter thorny issues that must be dealt with before specifications can be authored. Here we discuss several issues would, if not addressed early on, frustrate the development of the specifications.

Differences in opinions and advancement in understanding can lead to a desire to restructure existing data layout. As we consider solutions to these issues, we must strike a balance between keeping the legacy solution – that is, the ENDF-102 (Trkov, 2009) with which the nuclear data community is already familiar – and the desire to reach for new solutions. We also comment that we will need to strike a balance in how deeply to nest the hierarchy since some storage schemes perform better with a flatter hierarchy (e.g., HDF5) even though a deep hierarchy may make sense for organizing the data more clearly.

1. Is it a material property or a reaction property?

Some kinds of data can be viewed as reaction-independent properties of a target material. A case in point is the gamma branchings from an excited nuclear state of a nucleus. An excited nucleus decays via a gamma cascade through the lower levels of the nucleus in accordance with the tabulated gamma branchings, independent of whether it was formed in a neutron induced reaction or fission or any other process that leads to the same compound system.

Given this, we view all “particle/material properties” as data that are independent of the excitation mechanism. This includes (but is not limited to):

- For atomic nuclei:
  - target mass
  - number of neutrons, protons (and maybe even hyperons!)
  - nuclear level schemes (energies, spins, parities, ...)
  - level lifetimes, level widths
  - gamma and decay branching ratios from particles emitted during the de-excitation of excited states of a nuclear level
  - emission spectra from these nuclear decays

- For elements:
  - Z
  - Chemical symbol and name

- For atoms/ions:
  - mass
  - atomic shell properties (binding energies, spins, parities, ...)
  - X-ray decay and branching ratios from excited states
  - level lifetimes and widths
  - emission $e^-$ spectra from the internal conversion of gamma rays emitted from nuclei
  - charge state

- For composite materials (as encountered in thermal neutron scattering):
  - target density (at STP)
  - target stoichiometry

These lists may be amended as needed in the discussion below and a deeper discussion of them will be presented in the requirements for the material properties database. Indeed, it was recognized at the December 2013 WPEC meeting that, in order to ensure consistency of masses, Q values, levels and gammas within an evaluation, an external database is needed to perform this role library-wide. This database addresses main requirement 1.3 and is covered in section III.

A separate particle properties database also provides us a mechanism to store the ENDF/B-VII.0 and ENDF/B-VII.1 Decay and Atomic Relaxation sublibraries (Chadwick, 2006, 2011).
2. Different optimal representation in different physical regions

There are different optimal representations of data in different physical regions. For example, at low energies (i.e., in the resolved resonance region), neutron scattering is best described with an R matrix approach, whereas above the \((n,n')\) threshold it is best described via tabulated data as shown in Figure 1. This implies for example that

- Different physical regions may require us to change our concept of what a target is (e.g., fast neutrons see a single nucleus while thermal neutrons may (in)coherently scatter off many atoms in a material)
- Different macroscopic environments require us to change our description of microscopic data (e.g., Doppler broadening due to temperature effects)
- Different incident energies affect what particles are produced (e.g., pre-equilibrium, multifragmentation, particle production, spallation)

This fact was already recognized in the design of the legacy ENDF format and is a reality we too must confront [Trkov 2009]. Hence, we not only must consider different optimal representations (e.g., resonance parameters) in different physical regions, but we also must consider

- Different alternate representations (e.g., resonance parameters versus reconstructed pointwise cross sections)
- The matching (and potentially overlap) between representations
- A mechanism to “glue” them together, especially in cases where the concept of a target or reaction changes dramatically (e.g., thermal neutron scattering on molecules transitioning to high energy neutron resolving the nuclei in the atoms of the molecule)

3. Ensuring consistency

As we design the format(s) and supporting infrastructure, it is important to maintain internal consistency of the data. Within an evaluation, we must ensure at the very least consistency between:

- Cross section sum rules
  - Summing to the total cross section
  - All \((n,n')\) cross sections sum to total inelastic (ditto for other similar reaction types); similar to \(MT=3\)
- The sum of the prompt nubar and all delayed nubars must equal the total nubar in a fission reaction
- Masses, Q values, thresholds, upper energy bounds on secondary distributions
- Normalization conditions (on probability distributions, multiplicity tables, etc.)
- Energy and momentum balance
- Energy ranges of tables within a reaction and between different physical regimes
- Gamma branchings (that is, an excited state should have the same gamma-decay paths open no matter how it was produced)
- Resolved and unresolved resonance regions and the fast reaction region
- Original and processed data

Between evaluations, we must also ensure consistency between

- Fission product yields and decay data linkage
- Fast reaction region and the particle production region if they are stored in separate evaluations as is the case in for example JENDL and the JENDL-HE high energy library.
- Masses and other material properties
- Covariances and mean values between data common to both the Neutron Standards [Carlson 2009] and CIELO projects [Chadwick 2014].

Some of these can be handled with a simple hyperlink. Others may require capability within external processing/manipulation infrastructure. In the following discussions, we will point to features of the hierarchy that enable maintaining consistency.

Material properties are a special case where inconsistencies may need to be tolerated. In particular, when dealing with legacy ENDF evaluations, there is no guarantee that the same masses or level schemes are used consistently throughout an evaluation. As such, we may need to override any external material properties database with local versions within an evaluation.

4. Legacy data

More than six decades of work have gone into the creation of the nuclear data libraries and formats. Many very old data files are still in production and are needed for specific applications. These data must be supported until such time as they can also be updated.

The following requirements reflect the need to support legacy data:
FIG. 1 Cartoon of energy regions in a neutron induced reaction. The high energy region labeled “(smooth) fast” is usually handled in a very different way than the low energy regions where the R matrix approach applies.

5. Special cases

The ENDF/B-VII.0 and ENDF/B-VII.1 libraries (Chadwick, 2006, 2011) contain 14 separate sub libraries covering a variety of reaction data types. All of these must be covered by the format whose requirements we are drafting. However several of the data have special requirements that require further discussion:

- Particle production (Section IX)
- Radiative capture (Section X)
- Fission (Sections XI and XII)
- Thermal scattering law data (Section XIV)
- Charged particle reactions (Section XIII)
- Atomic data (Section VIII)

II. COMMON MOTIFS

In the sections to follow, we will describe several repeated patterns in the proposed data hierarchy. These repeated patterns or motifs include documentation elements (see subsection II.A), lists of reaction products and the sub elements of a <product> specification (see subsections II.C, II.D II.E II.F) and more complicated constructs detailing what data is derived from what other data (see subsection II.B). We note two other recurring “patterns”: covariance data (described in section VI) and Fission Product Yields (FPY). FPY data exists for both induced reactions (for example, the Neutron Fission Product Yield sublibrary in ENDF/B-VII) and as decay products from spontaneous fission (in both the ENDF/B-VII Decay and Spontaneous Fission Product Yield sublibraries). FPY format requirements are detailed in section XII. We comment that the motifs described here are also needed for the particle properties data.

A. Documentation

The documentation for an evaluation or part of an evaluation is in a way the most essential piece of information.
From it, we must be able to determine who performed the evaluation and how they did it. This is essential both for attributing credit (and blame;) and for debugging problems in an evaluation. This also aids in addressing requirement 1.6. Figure 2 illustrates the required contents of a documentation markup. We now discuss the elements of this markup before listing the requirements for a <documentation> element.

Each part of an evaluated data file may be evaluated separately, creating a “frankenevaluation” made up of stitched together evaluation parts. Therefore, we must allow <documentation> elements at many different levels in our data hierarchy. Furthermore, as this <documentation> element details a subset of information in an evaluation, the <documentation> element should contain both version tracking information for that subset as well as an unique identifier to reference the subset of information. The unique identifier should be a valid Digital Object Identifier (DOI). A DOI is a character string that uniquely identifies a piece of information and is associated with the URL where the information may be found. We must have a place for a DOI in the hierarchy because recent initiatives from both the United States and the European Union have imposed an additional legal burden on scientific data:

- European Union now issues Digital Object Identifiers (DOI) on documents produced with EU funding.
- In the United States, U.S. DOE Scientific and Technical Information Program (STIP), a collaboration working to increase the availability and transparency of the results of DOE-funded research and development, is working to make US generated data available through Office of Science and Technical Information (OSTI) (OSTI, 2015) through DataCite.org [DataCite 2015]. As part of this effort, the National Nuclear Data Center is authorized to assign DOIs to new datasets.

Therefore, since each node in the tree representing the data hierarchy can have its own URL, each <documentation> element should be allowed to have its own Digital Object Identifier (DOI).

We illustrate how DOIs can be used to reference independent parts of an evaluation in Figure 3. In this figure, each box represents one discrete part of a schematic evaluation. The coloration of the boxes is meant to represent the independent contribution from different publications, all “green” data come from the document with DOI #0, all “cyan” data come from the document with DOI #2, etc.

Moving down from the top element in figure 2, we note the <metadata> element. This element is currently undefined but can be adapted for individual data projects needs. For example, if an evaluation is meant to be indexed in the Nuclear Science References (NSR) database 

(Pritychenko, 2011), NSR keywords can be placed here. Alternatively, keywords for web searching may be placed here.

Continuing, the elements <title>, <abstract>, and <listOfAuthors> are clear. Next, the <versions> and <copyright> elements are other data project specific markups. The <versions> element allows a data project to store its own version information in whatever format the data project requires. The <copyright> element allows the data project to store the copyright notices for certain parts of their evaluation.

**Discussion point:**
It was suggested that a “model only” evaluation be flagged with some form of metadata. However, the form of such a flag is not agreed upon.

It is desirable to have enough detail in the documentation to reconstruct the evaluation exactly as the evaluator has produced it (complete with renormalizations of fitted data, etc.). Therefore we require markups for <writeup>, <listOfInputs> and <listOfEXFORSets>. The writeup itself is given in the <writeup> element and should be essentially free text (supporting internal formatting such as HTML or LaTeX). The <listOfInputs> element allows the evaluator to store their TALYS or EMPIRE input decks within the evaluation. The <listOfEXFORSets> is, as the name implies, the list of EXFOR datasets used in the evaluation. The IAEA’s EXFOR web application allows both on-the-fly data renormalization (example in Figures 4 and 5) and covariance generation (example in Figure 6). Therefore, both require spots in the hierarchy.

Finally, both <EXFORSet>s and the writeup bibliography (in the <bibliography> element) require a markup for bibliographic references. The easiest solution here is to adopt an accepted format such as BibTeX XML.

**Discussion point:**
How should a bibliographic entry be formatted?

**Resolution:**
In principal it should be shared with EXFOR. BibTeXXML (Gundersen, 2015) was suggested as a viable choice.

**Requirement 6: Documentation**

6.1 The format shall have metadata element.
6.2 The format shall have a markup for the DOI (likely including a link attribute)
6.3 The format shall have title element (using a <text> element)
6.4 Have markup for the evaluation date (authors
FIG. 2 Basic structure of a `<documentation>` element.

6.5 The format shall have a markup for the library acceptance date (library maintainers make this up), this is a legacy ENDF requirement.

6.6 The format shall have an abstract element (optional, using a `<text>` element)

6.7 The format shall have a markup for authors (names, affiliation, email, etc.). The corresponding author must be denoted.

6.8 The format shall have a markup for storing the input decks from codes used by the evaluators to prepare the evaluation (`<listOfInputs>` and `<inputDeck>` elements)

6.9 The format shall have a markup for the evaluation version

6.10 The format shall have a copyright element (for optional copyright notices).

6.11 The format shall have a markup for referencing the EXFOR datasets used in the evaluation. This markup should include a link to the
original data and details of any data rescaling or evaluator generated covariances. These can be done using V. Zerkin’s EXFORWeb tools [EXFOR 2015]. Examples are shown in Figs. 6–5.

6.12 The format shall have an element to store the writeup itself (using a <text> element)

6.13 The format shall have markup for the bibliography.

B. What data are derived from what other data?

According to main requirement 3.2, we require a mechanism to specify what data are “original” and what are derived. To accommodate this, we must allow the storing of the original and derived data, sometimes at the same level in the hierarchy. Derived data must point back to the original data with some key or a link. We note that while some data are derived from only one other element, some data are derived from several disparate pieces of information within an evaluation.

There are many cases where a capability to link original and derived data would be useful. Here are a few that come to mind at the time of this writing:

- Doppler broadened data at a temperature $T > 0\degree K$ should link to the $0\degree K$ data.
- Grouped and pointwise data
- Angular distributions converted between pointwise angular tables and Legendre moments
- Any (and all) parameterized data converted to pointwise
- Average energy deposited, average forward momentum deposited and KERMA factors are all derived using product distribution data and energy balance of all the particles emitted in a reaction
- Changes in interpolation schemes (e.g., log-log to lin-lin)
- Resonance data converted to pointwise
- Resonances with smooth backgrounds (this is allowed in ENDF, and we argue below that it should be deprecated in the new hierarchy).
- Cross section (and other data) may have uncertainty data and should be associated with any correlation matrices. Grouped or other representations of the data need to point back to both the uncertainty data and the mean value data.
- Monte Carlo realizations of a data set should point to the mean value and the associated covariance

Fig. 7 illustrates what we need in practice for a mid-level container. The top level of the mid-level container may contain within it several representations of the underlying data. In this cartoon, the different representations are stored in a <listOfDataRepresentations>
element. The mid-level container must then have a link to the actual dataset in `<listOfDataRepresentations>` that stores the original or “official” data. This simplifies navigation since we don’t have to query which data version is the original. Each contained dataset has a flag to denote whether the set is “official” or not and all derived data contain a link to the data from which it was derived from. In this example, set #2 is derived from set #1 (it might be from a group collapse for example) and set #1 is derived from set #0, the original set (set #1 may be the grouped version of set #0).

This concept was pioneered in GND. In GND, the different versions of the same data are each encapsulated within a `<form>` element. The form element either is the data container itself or the lowest level of the top-level hierarchy before encountering the actual containers holding the data. GND links derived data to the original data with a `nativeData` attribute.

Figure 8 shows another example of how original and derived data can be handled. In this figure, we show the original cross section given as the mean values of the cross section and a covariance matrix. From that, one can derive a second representation of a cross section with uncertainties as well as a correlation matrix.

We now summarize the general requirements for a mid-level container that supports linking between original and derived data:

### Requirement 7:

**Linking between original and derived data**

1. A data containing element shall have a `<listOfDataRepresentations>` that contain the different version of the data.
2. A data containing element shall have one specific implementation of a data table plus additional attributes.
3. Derived data shall have a `<listOfDerivedFromLinks>` with one or more `<derivedFrom>` links to the original data.
4. Derived data should be placed as close to the original data as possible.
5. A `<derivedReactions>` branch should be present to store trees of entire derived reactions.
C. Product list elements

Both particle decays and induced reactions produce reaction products. In addition, in legacy ENDF there is the possibility of “orphaned products”, that is, products that are not correctly associated with their parent reaction. Therefore we will need an element that encapsulates lists of reaction products. To support ENDF, we will require three different kinds of product lists:

- A `<reactionProducts>` element to list the reaction `<product>`s of an induced reaction. In GND, a `<reactionProducts>` element is referred to as `<outputChannel>`.
- A `<decayProducts>` element to list the daughter particles into which a parent `<product>` decays.
- An `<orphanedProducts>` element to list `<products>` that are not associated with any reaction, as is found in older evaluations (especially for gamma products).

Although each of these product lists share nearly the same structure, we assume that it is easier to work with unique names as it helps to denote the context in which they are used. Whether they have separate names or the same name is up to the authors of the specification document. Each list of reaction products has additional individual requirements that we list here. Their use in either the context of induced reactions or the decay of particles are discussed later in this document.

In Fig. 9 we show the common structure of all of these lists. From the list of products it should be possible to reconstruct the reaction designator in the `<reaction>` element (see section IV.C).

**Requirement 8: `<reactionProducts>`**

8.1 List of `<product>` elements
8.2 The `kinematicType` (e.g., two-body, uncorrelated). Elastic reactions and all resonance reactions using the R matrix formalism are two-body reactions. GND refers to these by the name `<genre>`. Table I lists allowed kinematic types.
8.3 A flag to denote whether to use relativistic or non-relativistic kinematics when handling this channel
8.4 Optionally the Q value of the reaction for backwards compatibility with the legacy ENDF format

**Requirement 9: `<orphanedProducts>`**

9.1 List of `<product>` elements
9.2 The `kinematicType` (most likely uncorrelated). Table I lists allowed kinematic types.
9.3 A flag to denote whether to use relativistic or non-relativistic kinematics when handling this channel
9.4 Optionally the Q value of the reaction for backwards compatibility with the legacy ENDF format

**Requirement 10: `<decayProducts>`**

10.1 Intended mainly for use in particle property databases
10.2 List of links to decay products or the decay `<product>`s themselves
10.3 The `kinematicType`. Table I lists allowed kinematic types.
10.4 Q. For new evaluations, this is optional but should sync with material properties. For legacy evaluations it must be retained.
10.5 `lifetime`, a material property but it may be useful to repeat it here, so it is optional

D. `<product>` elements

In the legacy ENDF format, outgoing product data are stored in a variety of ways, potentially leading to
confusion for the user:

1. Outgoing neutron energy distributions are stored in an MF=5 file and angular distributions in an MF=4 file. Alternatively, both may be stored in an MF=6 file. Double differential neutron data can only be stored in an MF=6 file. Neutron multiplicities from fission are stored in MF=1, MT=452, 456, and 455 files.

2. Outgoing charged particle data are stored exclusively in MF=6 files.

3. Outgoing gamma data can be stored in MF=6 files or in a combination of MF=12 (for multiplicities and discrete level energies), MF=13 (production cross sections), MF=14 (angular distributions) and MF=15 (energy distributions). Additionally, delayed gamma data from fission are stored in MT=1, MF=460.

4. Additionally, the energy released from fission is stored in MF=1, MT=458 and is not associated with the produced particle.

We would like to simplify and unify these options into a simple product (and daughter) elements.
TABLE I Kinematic types for an attribute for <reactionProducts>, <decayProducts> and <orphanedProducts> elements.

<table>
<thead>
<tr>
<th>kinematicType</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>two-body</td>
<td>only two products are emitted per step in a reaction, the products are correlated, and only the center-of-mass angular distribution is needed in order to calculate the double-differential distribution</td>
</tr>
<tr>
<td>uncorrelated</td>
<td>the products are treated as being uncorrelated from each other, and a complete double-differential distribution is required for each product</td>
</tr>
<tr>
<td>correlated</td>
<td>proposed</td>
</tr>
<tr>
<td>fission</td>
<td>proposed</td>
</tr>
</tbody>
</table>

defined in the previous section. The product element must support storing (at least) a multiplicity and outgoing distributions.

The <product> structure is illustrated in Figure 10. We note that, in this figure, the <product> element has both documentation as well as all of the derived and evaluated distributions and multiplicities corresponding to the reaction or decay product.

Often these products further react either because they are an intermediate state (as in a breakup reaction), or metastable or unstable (as in decay data). To enable this, we allow the <product> element to have both <reactionProducts> and <decayProducts> within. This enables multistep, breakup or decay reactions in an (hopefully) obvious way. Strictly speaking, if one has an external (or even internal to the evaluation) particle properties database, the <decayProducts> element could be placed with the particle properties rather than in a <reaction> element.

Each <product> should have:
Requirement 11: <product>

11.1 The particle’s identity. For a transportable particle this refers directly to the particle. For a particle (nucleus or atom) with excitation levels, this refers to both the particle and the excitation level.

11.2 The multiplicity

11.3 An optional ENDF conversion flag (deprecated) to tell the user whether the original ENDF product data was in MF=6 or a combination of MF=4,5 or MF=12,13,14,15 data and to aid in reverse translation.

11.4 All specified outgoing particle distributions for that particle. These are mid-level container elements, see requirements 7. This can include things not pictured in figure 10 such as the mean energy and forward momentum of the particle to enable heating calculations.

11.5 If the product can subsequently decay, then a <decayProducts> element should be allowed.

11.6 If the product is an intermediate state in a breakup or multi-step reaction, then a <reactionProducts> element is allowed.

11.7 Optional transfer matrix for group-wise deterministic calculations

Discussion point:

The nesting of product lists inside a <product> list leads to schema recursion. This is allowed in XML and other hierarchical format, but may lead to complicated schema coding and should be avoided if possible.

Resolution:

Noted. One solution is to provide multiple product list element tags with names that depend on the nesting depth.

E. <distributions> and <distribution> elements

The <distributions> contains all of the outgoing probability tables associated with a reaction product. Each probability table is contained inside a <distribution> element corresponding to some variation of \( P(\mu, E'|E) \) that can be used for a transport application. More than one distribution may be stored inside the <distributions> (i.e., both original and derived distributions may be stored as separate <distribution> elements).

Several different types of distribution data are possible. These include angular distributions \( P(\mu|E) \), used for two-body reactions where knowledge of the outgoing angle of one product is sufficient to reconstruct all reaction kinematics for both products\(^3\), and double-differential distributions \( P(\mu, E'|E) \), used to give the full probability density for products coming from more complex reactions. Many evaluations use a special case of \( P(\mu, E'|E) \) by making the simplifying assumption that the outgoing \( \mu \) and \( E' \) distributions are <uncorrelated> and can be stored separately as \( P(\mu|E) \cdot P(E'|E) \) (e.g., MF 4 and MF 5 in ENDF-6 rather than in MF 6).

In GND, the <distributions> element contains one or more components, each of which corresponds to a

\(^3\) For a two-body reaction, the double differential distribution is related to the angular probability distribution by \( P(\mu, E'|E) = \delta(E' - E'(E, \mu))P(\mu|E) \).
type of distribution (angular, angular-energy double differential, uncorrelated, etc.). Each component in turn (except for the <uncorrelated> component) contains one or more forms, which store the distribution data. The <uncorrelated> component is an important exception: instead of forms it contains two other components, <angular> and <energy>, each of which in turn contain forms.

The organization used in GND has some advantages: it starts by storing the type of distribution (at the component level) and then the details of how that distribution is being stored (at the form level). However, this two-level organization has drawn some criticism, partly due to confusion over the special treatment required for <uncorrelated> distribution components, and also partly due to its limited capability for handling different forms of distribution in different energy regions (like ENDF, GND permits mixing Legendre distributions at low incident energy with pointwise angular distributions at higher incident energy, but other combinations are not supported).

Unfortunately, in the opinion of the GND authors there is no way to avoid treating <uncorrelated> as a special case since it must contain both \( P(\mu|E) \) and \( P(E'|E) \); neither by itself is sufficient to describe the distribution. Some other possible data hierarchies should be considered, however. For example, instead of nesting forms inside of component elements, both the type and the form of distributions could be given as metadata, as in \(<\text{distribution type}="\text{angular}" \text{ form}="\text{pointwise}">\). Also, a general-purpose <piecewise> distribution form could be defined, containing two or more distribution <region> elements each of which applies over a specified range of incident (and / or outgoing) energies.

![Diagram](image)

**FIG. 11** Overview of a <multiplicity> element.

would be helpful in storing processed data at various temperatures for Monte Carlo transport where one only heats the cross sections. One could then generate the heated cross sections and store the cross sections in evaluations at different temperatures and connect them with the metaEvaluation markup. To reduce the massive redundancy in the outgoing distributions (they never get heated), all the distributions in the heated evaluations could then link back to the zero temperature file’s distributions. In fact, it may be more economical to have the entire <reactionProducts> element link to another.

**Resolution:**
Agreed

**F. Multiplicities**

A <multiplicity> element would be used by atomic scattering and nuclear reaction data. It is analogous to ENDF’s MF=12 (for gammas) or MF=1, MT’s 452, 455 or 456 (fission \( \bar{\nu} \)'s), but simplifies and unifies the structure by defining one common scheme for multiplicity specification.

**Requirement 13:** <multiplicity>

13.1 The <multiplicity> element is only necessary for non-constant multiplicity. Constant multiplicity could be stored as a product attribute or with a constant markup. Storing multiplicity as an attribute in one place and an element in another leads to a more complex format and should be avoided. Therefore, <multiplicity> should be an element.

13.2 A non-constant <multiplicity> element consists of at least one representation contain-
an <interp1d> with a dependent variable (the multiplicity itself) given in units of number of emitted particles and an independent variable (projectile’s incident energy) in units of energy. The first energy point could be (real or effective) threshold or the lowest energy supported by the encapsulating evaluation. The <multiplicity> is assumed to be zero outside of the specified energy region. If the <multiplicity> is variable and given as a non-integer (as is common for fission $\nu$), it is up to the code using the data to interpret the data correctly.

13.3 Need specification for $P(\nu|E)$ for fission neutrons.
13.4 These are mid-level container elements (see requirements 7).

Discussion point:

For consistency, should we require that all multiplicities be given in <product> elements as an element and eliminate the idea of storing constant multiplicities as an attribute? This would make for simpler coding of an API and clearer data files at a small cost of verbosity.

Resolution:

Agreed

Discussion point:

Should we allow multiple names for the same element? Physically, <promptNubar> is just the average multiplicity for the prompt neutron product, so it could be stored just like any other product multiplicity. However, <promptNubar> is easier to search for...

Resolution:

No, it leads to a more complex format.

Discussion point:

(Actually a continuation of the previous discussion point): in GND/XML, converting from metadata to a unique element only takes up a little more space. However, when translating to HDF5 each element is converted into a unique ‘group’ which takes up a minimum of about 1.3 kB, so the difference between metadata and element can become significant.

Resolution:

This is a serious issue with one particular potential implementation (HDF5). Therefore it will be up to the developers of the API for the data to deal effectively with this issue.

III. PARTICLE AND/OR MATERIAL PROPERTIES DATABASE

Task 1.3 of WPEC Subgroup 38 (SG38) is to define a database hierarchy for handling particle information needed for nuclear reaction evaluations and transport codes. The hierarchy must be general enough to describe relevant particles, including mass, charge, spin and parity, half-life, decay properties, and so on. In a way, the hierarchy encapsulates the needs of the ENDF decay and atomic relaxation sublibraries as well as the RIPL mass and level tables. Particle databases built with this hierarchy are meant to serve as central locations for particle information that can be linked to from codes and other databases. It is hoped that the final product is general enough for use in other projects besides SG38.

While this is called a “particle database”, the definition of particle (as described in Section 2) is very broad and could include materials in the thermal scattering law sublibrary as well as atoms and ions. What we refer to as a “particle” includes the projectile and product(s) as well as what the ENDF-6 format refers to as a material or MAT. The database is meant to be general enough to include not only fundamental particles like quarks and leptons, but also composite particles like mesons, baryons, atoms, nuclei and even excited states. Under this definition the list of possible particles becomes quite large.

To help organize them the database will need a way of grouping similar particles together into “families” with similar attributes.

The following list of requirements is a summary of the larger discussion of requirements and specifications in the “Requirements and specifications for a particle database” document (WPEC Subgroup 38 2015a). The key words “shall” and “should” will be used to differentiate between requirements and recommendations:

**Requirement 14: Particles**

14.1 Each ‘matter’, ‘particle’ and ‘alias’ instance in the database shall have a unique id used to identify and refer to it. Only these classes shall have ids (for example, no id is given to the mass or spin, only to the instance itself)

14.2 Every particle shall contain at least the following properties: mass, charge, spin, parity and half-life (which may be ‘stable’). How-
ever, some of these properties may be inherited from higher in the hierarchy rather than being listed explicitly (see requirement 4).

14.3 The database shall support storing uncertainties with all particle properties. Uncertainties may be given either in the form of a central value with uncertainty (for example, mass = 54.938 +/- 0.729 amu) or as a list of multiple possible assignments (for example, spin = 3/2, 5/2 or 7/2). If multiple assignments are listed, the database shall require that one assignment be explicitly listed as the ‘recommended’ value.

14.4 The database shall use nesting and inheritance where possible to reduce redundancy by grouping similar particles together. For example, the database should support grouping isotopes together inside an element, such that all isotopes inherit the same nuclear charge Z from the element.

14.5 The database shall support defining ‘families’ to classify similar particles. Each particle family may have additional required data elements (beyond the list in requirement 2). For example, a ‘lepton’ family may be defined, where each lepton requires a lepton number in addition to mass, charge, spin, parity, etc.

14.6 If several different decay modes are open, the database shall permit storing each mode along with its probability. Decay modes may be grouped: for example, if a particle is subject to beta decay, gamma emission or internal conversion, gamma and IC decays should (shall?) be grouped together as electromagnetic decays.

14.7 Within each decay mode, all direct decay products shall be listed explicitly along with their particle ids. Only decay products emitted directly by the current particle shall be included in the list of products. If any of those products can also decay, their decay properties shall be accessible by looking them up through the particle id.

14.8 The database shall support storing different charge states of atoms. For example, it must be able to differentiate between an alpha particle and a 4He atom.

14.9 The database shall support storing documentation sections inside (at least) each particle and each property within that particle.

14.10 The database shall support a bibliography section. Each item in the bibliography shall include a unique citation label that can be used to refer to it from any documentation section.

14.11 The database shall support a section that defines a list of aliases for particles. For example, the id Am242_m1 could be an alias for Am242_e2

One important function of the particle database will be to provide an easy way for codes and external databases to look up any particle stored inside. In order to make this access as simple as possible, the database will include a unique name (or “id”) for every particle that it stores. Users can then access a specific particle either by providing a full path to it or simply by using its id. Suggestions from experience with GND:

- **Aliases**: A limited number/scope of aliases for commonly used particles, such as “e” for electron or “α” for alpha or “n” for neutron. Also to associate a level of an isotope with an isomer.

- **Compounds**: c_String_Describing_Material can be used to specify say H in ZrH or the phase of the material. Useful for Thermal Scattering Law data

- **Elements**: Sym0 (e.g., Fe0 or C0), useful for atomic data

- **Isotopes**: SymA (e.g., Fe56)

- **Nuclear levels**: SymA_eN (e.g., V51_e1 for the first excited state of 51V or SymA_c for continuum).

- **Excitations of an atom**: Sym0_eN (e.g., V0_e1 for the first electronic excited state of 51V or Sym0_c for continuum).

**SUGGESTED TEST**: Test that all particle id’s are unique since they are keys and we must avoid key collisions.

**Discussion point:**

Regarding requirement 14.3 what if a particle has several possible assignments that are correlated, as in ‘2+’ or ‘3-’ (i.e., spin AND parity together)? Should we always treat spin/parity together (as in <Jpi spin="2" parity="-"/>)? Also, in addition to supporting multiple assignments should the database support flags such as ‘firm’, ‘tentative’, etc?

**Discussion point:**

Regarding requirement 14.7 this might cause problems for handling correlated decays. Also, how should we handle a case where we know a particle decays, but do not know its final state with absolute certainty? For example: imagine a nucleus with two beta-decays and the second has short half-life. What if we dont know levels populated in the first decay? In addition to giving explicit decays to final states,
do we also need to support a ‘spectrum’?

**Discussion point:**
Do we need to explicitly add requirements for bands and flags to indicate how complete a level scheme is?

**Discussion point:**
Do we need to be able to explicitly store correlated decays, such as the angular correlations between $^{60}$Ni gammas following $^{60}$Co beta decay? The angular correlations arise because the two M2 transitions happen back-to-back so the second decay is not truly independent of the first. However, the correlations can be predicted if we know the M2 nature of each gamma. Perhaps we can avoid this problem by adding more information to each gamma and requiring user codes to make their own correlations if needed.

**Discussion point:**
Requirement 14.6 deals with grouping decay modes together. One advantage: this makes it easy to support familiar ideas like the internal conversion coefficient (ICC) that gives the ratio between gamma decay and IC. However, we need some thought about how we store decay probabilities. Conceptually the simplest is to store each as a float and require all probabilities sum to one, but historically evaluators have taken different approaches. In ENDF for example, most probable gamma is assigned ‘intensity’ of 1.0, and other gamma intensities are given as a ratio to that one. Should we support that?

**Discussion point:**
Check this again: is the most intense gamma really 1.0?

**IV. ONE EVALUATION**

The top level of data files in all major libraries is the “evaluation”, consisting of one target material and one projectile and all the data that goes with the reactions between this pair. This arrangement is familiar to the nuclear data community and should be embraced going forward.

Because of the different kinds of evaluations, what happens below the uppermost node in the hierarchy can differ from sublibrary to sublibrary. There are three main classes of reaction sublibrary that concern us:

- **Thermal scattering law data:** neutrons reacting with such low energy that the de Broglie wavelength of the neutron is too large for the neutron to resolve individual nuclei (in principle other particles could do this too)
- **Atomic scattering data:** electron and photon interactions with atoms
- **Nuclear reaction data:** any projectile impinging with enough energy to interact with an atomic nucleus. This collection of data can include resonance data which is arguably different enough from fast reaction data to merit its own discussion.

An ENDF-like decay sublibrary is discussed in the context of a material properties database.

For the purpose of discussion, we name the top-level element of an “evaluation” `<evaluation>`. In GND, the equivalent concept is called a `<reactionSuite>`. The `<evaluation>` element is expected to be the root XML node of a file.

In all cases, the data can be arranged using a consistent set of rules. Because of the different optimal representations in the fast and resonance regions, we must distinguish between evaluated data broken out by reaction and tabulated one at a time (using the `<reactions>` element) or evaluated data tabulated in a parameterized form such as in the resonances region (using the `<resonances>` element). We also need spots for covariance data that is not associated with particular reactions (e.g. cross–reaction covariances). Furthermore, we must provide markup that supports transport and other applications through `<derivedReactions>` and `<derivedTransportData>` markups to stored derived data and a `<styles>` markup for things like default group structures and such. We mention the list of function definitions that are detailed in section XVI as the `<functionDef>` markup would only rarely be used and used only in “unofficial” evaluations. Finally, at the top level we store a local particle properties database in the `<particles>`. This is meant to override default values given in some external reference database. This top level arrangement is shown in Figure 12.

**Requirement 15:** `<evaluation>`

15.1 Require one projectile (e.g., “n”)
15.2 Require one target material (e.g., “Fe0”)
15.3 Require the version of the data format
15.4 Require the library designator (i.e., a name string that says “ENDF/B” and a version string that says “VI.1”)
15.5 File-wide specification of the Lorentz frame of the incident energy of the projectile
15.6 Optionally support other data the library maintainer needs for proper data manage-
FIG. 12 Top level arrangement of an <evaluation> element. Only the documentation element is required in an <evaluation>, but <reactions>, <resonances>, and <covariances> are expected in nearly all (neutron induced) reaction evaluations. The <styles>, <particles> and <functionDefs> elements are used primarily to override or (re)define default behaviors. Finally, the <derivedReactions> and <derivedTransportData> elements are nearly exclusively for processed data.

15.7 Require a temperature attribute: for low enough energy projectiles, this is a crucial piece of information. For neutrons, Doppler broadening is important to determine effective reaction rates and to get self-shielding corrections. For astrophysical applications, the temperature of the plasma is needed to handle charge screening properly.

15.8 Require ELow and EHigh attributes to specify the energy range of validity of this evaluation.

15.9 Require an activationFlag attribute to signal whether the data in this evaluation is meant for activation or for particle transport. The two applications have very different completeness requirements that, in the XML variation of a format, can be enforced by checking against an XSD file.

15.10 Require a file-wide <documentation> element. In GND this is handled with a styleInformation attribute.

15.11 Optionally a material database to override defaults with values local to the evaluation (the <particles> element, described in reference [WPEC Subgroup 38, 2015a]).

15.12 Optionally a place for evaluation–wide default style information such as group-structures, fluxes, etc. (see the <styles> element description IV.F)

15.13 Optionally a place for covariance data (see the <covariances> section for more detail VI)

15.14 Optionally a <reactions> element (more on <reactions> in the subsection IV.B)

15.15 Optionally a <resonances> element (more on <resonances> in the subsection IV.H)

15.16 Optionally a <derivedReactions> element (see subsections XV and IV.G)

15.17 Optionally a <derivedTransportData> element to store application specific
data that cannot fit within a single `<derivedReaction>`. We discuss the derived transport data in section XV.

Discussion point:
Should we also be defining higher-level organization, such as ‘library’ or ‘projectile’? Different institutions may have very different ideas of how these higher levels should be organized, but we may still be able to standardize to some degree.

Resolution:
Agreed, see section V

Discussion point:
As the resonances and the fast regions are two distinct physical representations of data in two different energy regions, it might make sense to require that they NOT be together in the same evaluation and that users use the `<metaEvaluation>` markup to combine them. This simplifies bookkeeping and ensures that users understand that they are different things that they must combine themselves. However, this is a change from the ENDF mindset and would complicate translation of the outgoing particle distributions in the resonance region. Many ENDF forms don’t support the calculation of outgoing angular distributions from the resonance parameters so evaluators must provide these tables.

Resolution:
This idea is too much of a change from the ENDF arrangement and will likely lead to confusion.

Discussion point:
Within an evaluation in a particular sublibrary, one must ask whether to arrange the data per-energy or per-reaction. For data to be per-energy or energy-major, we mean that all data (reactions, cross sections, distributions, etc.) for one incident energy are collected together in one parent element. For data to be per-reaction or reaction-major, we mean that all data (cross sections, distributions as a function of incident energy) for one reaction are collected together in one parent element. A per-energy arrangement is particularly convenient for Hauser-Feshbach (and other) modeling codes because one normally computes one energy at a time and loops over incident energies to assemble the full evaluation. An energy-major arrangement has certain benefits and drawbacks:

- **Energy-major benefits**
  - Natural output of a reaction model such as EMPIRE or TALYS
  - Energy-major is natural for sampling in Monte Carlo transport: in a collision in a Monte Carlo calculation, one samples between the reaction cross sections for different channels at a fixed incident energy
  - One can see at a glance what channels are open and compete with one another

- **Energy-major drawbacks**
  - Very difficult to plot say a cross section as a function of incident energy
  - Can be difficult to compare to experimental data taken at different energies
  - Although most Monte Carlo transport codes sample on a per-energy basis, the codes are all written assuming a reaction-major arrangement of data and therefore would require major refactoring to reap the benefits of an energy-major arrangement. Besides, if an incident energy falls between two tabulated energies, all sampled quantities must be interpolated.
  - Difficult to diagnose unphysical discontinuities as a function of incident energy
  - Even a single resonance in a resonance regions span many energies so an energy-major arrangement seriously complicates resonance reconstruction
  - Hard for deterministic codes to use
  - Not familiar to users as legacy ENDF data are stored with the reaction-major arrangement
Resolution:
The arrangement of data must reflect the users’ needs first, not the evaluators. Therefore, we argue that the benefits of an energy-major arrangement do not outweigh the drawbacks and so we recommend maintaining the ENDF-style reaction-major arrangement. However, denoting the energy range of validity of an evaluation coupled with the <metaEvaluation> concept allow an evaluator to achieve the effect of an energy-major arrangement by having only one incident energy in an evaluation. It was suggested that someone develop a tool to combine these one-energy sized evaluations into a complete reaction-major evaluation. This tool would then be reusable for data generated using any Hauser-Feshbach code.

Discussion point:
Do we need an index or directory outlining the evaluation such as what Red Cullen produces with DICTIN (part of PREPRO)?

Resolution:
No, the hierarchical arrangement of data should make this unnecessary

Discussion point:
It was suggested that we allow <particles>s at other levels in the hierarchy to, say, allow per-reaction mass of target specifications. This would make it easier to re-use legacy evaluations that perform well in applications but yet use older values of fundamental data.

A. Reaction designation

Common sense requires that we have a unique method for identifying a reaction that is preferably human readable and easily parseable by machine. Examples from GND (Ref. [Mattoon, 2012]) are shown in Table II.

We note that GND’s scheme is more general than ENDF’s MT designator and this scheme does not meddle MF and MT (as happens in the fission reactions in ENDF). In GND’s scheme, the reaction designator is not always unique but is derivable from the reaction products (and their decay products if this is a breakup reaction). However, the user does have the ability to define their own reactions and add qualifiers to existing ones.

<table>
<thead>
<tr>
<th>GND reaction label</th>
<th>MT</th>
</tr>
</thead>
<tbody>
<tr>
<td>( n + Pu^{239} \rightarrow n + Pu^{239} )</td>
<td>2</td>
</tr>
<tr>
<td>( n + Pu^{239} \rightarrow n + Pu^{239} \text{[compound elastic]} )</td>
<td></td>
</tr>
<tr>
<td>( n + Pu^{239} \rightarrow n[\text{multiplicity:}2] + Pu^{238} )</td>
<td>16</td>
</tr>
<tr>
<td>( n + Pu^{239} \rightarrow n[\text{multiplicity:}3] + Pu^{237}_{e1} )</td>
<td></td>
</tr>
<tr>
<td>( n + Pu^{239} \rightarrow n + Pu^{239}_{e1} )</td>
<td>51</td>
</tr>
<tr>
<td>( n + Pu^{239}<em>{m1} \rightarrow n + Pu^{239}</em>{e} )</td>
<td>91</td>
</tr>
<tr>
<td>( n + Pu^{239} \rightarrow Pu^{240} + \gamma )</td>
<td>102</td>
</tr>
<tr>
<td>( n + Pu^{239} \rightarrow Pu^{240}_{e1} + \gamma )</td>
<td></td>
</tr>
<tr>
<td>( C^{12} + Pu^{239} \rightarrow C^{12}<em>{e2} + Pu^{239}</em>{e1} )</td>
<td></td>
</tr>
<tr>
<td>( n + Be^{7} \rightarrow (Be^{8} \rightarrow He^{4}[\text{multiplicity:}2]) )</td>
<td></td>
</tr>
</tbody>
</table>

TABLE II Example of reaction labels in GND. ENDF MT numbers are listed when possible. Some GND reactions have no MT equivalent. From Ref. [Mattoon, 2012]

Whatever scheme we use, users should immediately be able to tell what the target, project, and all reaction products are. We should also be able to define aliases for commonly understood names like “fission” and “spallation”. Additionally, we should be able to describe some reactions with additional qualifier which could allow us to say break “elastic” scattering into “compound elastic” and “shape elastic”. This would allow us to better represent the reaction distributions since neutrons from “compound elastic” reactions are typically isotropic in the residual nucleus rest frame while “shape elastic” neutrons are typically forward–peaked.

In what scheme is finally agreed on for the reaction designators it should follow these recommendations:

**Requirement 16: Reaction designation**

16.1 Should be shared/agreed upon with EXFOR
16.2 Should not be limited to simple targets (we need to denote thermal neutron scattering data)
16.3 Support aliases for things like “elastic”, “total_fission”, “capture” and “spallation”.
16.4 Support need to distinguish input vs. output channels
16.5 Allow uncorrelated particle emission
16.6 Support processes with non-constant multiplicities
16.7 Support sequential processes (esp. 2-body)
16.8 Support qualifiers such as “compound_elastic” and “shape_elastic” which allow evaluators to split up reactions with common final products but different reaction mechanisms and/or kinematics.

**SUGGESTED TEST**: Test that all reaction id’s are
unique since they are keys and we must avoid key collisions.

B. The collection of reactions: <reactions>

Below the <evaluation> markup, most of the data are stored in the <reactions> and <resonances> branches. The <reactions> element is pictured in Figure 13. To understand the arrangement of data here and in the <reaction> and <derivedReactions> elements below, it is useful to have a mental model for particle transport.

For a neutral particle (such as a neutron or a gamma) with energy $E$, the mean free path ($\lambda_{\text{mfp}}$) is defined as

$$\lambda_{\text{mfp}} = \frac{1}{\sum_x n_x \sigma_{x,\text{tot}}(E)},$$

where the sum goes over each material $x$, and $\sigma_{x,\text{tot}}$ and $n_x$ are the total cross section and number density respectively of material $x$. The mean free path determines the transit distance and time between “hard” nuclear collisions. The total cross section may be tabulated in a <derivedReactions> element.

For charged particles, the total cross section does not exist because of the Coulomb singularity in, for example, the elastic scattering reaction. Coulomb scattering is “soft” in that the Coulomb force always acts to gently nudge the projectile at all distances. To implement Coulomb scattering in practice one divides up scattering events by angle relative to the center-of-mass momentum of the target and projectile. At small angles, one uses condensed history treatments (Seco, 2013). At large angles, one uses the large-angle Coulomb scattering approximation which treats the reaction as a “hard” inelastic collision. This mixture of treatments means the data cannot be heated.

Regardless of the projectile-target combination, once we decide that a “hard” collision will occur, we can proceed as follows:

- Determine which material was hit by assigning probabilities proportional to $\sigma_x(E)$
- For the material, decide what reaction occurred by assigning probabilities proportional to the ratio of each partial cross section to the total cross section for that material
- Once the reaction is decided, determine what particles will be emitted
- Loop over emitted particles
  - If the multiplicity is not constant, sample the number of emitted particles using the energy-dependent multiplicity which may be given as a distribution (e.g., $P(\nu)$).

- Depending on the kinematics (two-body or uncorrelated), sample the emitted particle’s energy and/or angle
- If particle or reaction product decays, follow the decays ...

- Where possible, use energy/momentum conservation to determine the recoil energy and momentum

C. The <reaction> element

With a reaction major arrangement, there is one common motif in the three different sublibraries described in Section IV which for a lack of a better name, we call the <reaction> element. This element denotes one reaction that can be sampled in a Monte Carlo code. In it, we specify reaction <crossSection> and the outgoing particle distributions for all emitted particles.

We comment that data computed from reconstructed resonances are placed in the appropriate sub-elements of a <reaction>. This reconstructed data would of course have hyperlinks connecting the derived data back to the data in the <resonances> branch.

<table>
<thead>
<tr>
<th>Requirement 17: &lt;reactions&gt;</th>
</tr>
</thead>
<tbody>
<tr>
<td>17.1 An optional &lt;documentation&gt; (not shown)</td>
</tr>
<tr>
<td>17.2 A list of &lt;reaction&gt;s with actual reaction data inside</td>
</tr>
<tr>
<td>17.3 An optional (and deprecated) &lt;orphanedProducts&gt; branch. This allows for the storing of older evaluations. In many old ENDF evaluations, gammas were not associated directly with their source reaction, but were lumped in MT=3 or MT=4. This destroys energy balance and should be deprecated. This branch provides a spot for these unassociated gammas and other particles.</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th>Requirement 18: &lt;reaction&gt;</th>
</tr>
</thead>
<tbody>
<tr>
<td>18.1 An optional &lt;documentation&gt;</td>
</tr>
<tr>
<td>18.2 The reaction data itself using at least one of these schemes:</td>
</tr>
<tr>
<td>18.2.1 Option #1, breaking out cross sections and outgoing distributions</td>
</tr>
<tr>
<td>18.2.1.1 A &lt;crossSection&gt;</td>
</tr>
<tr>
<td>18.2.1.2 A &lt;reactionProducts&gt; element listing the reaction &lt;product&gt;s. From the list of products it should be possible to reconstruct the reaction designator unless the reaction is denoted by some kind of reaction</td>
</tr>
</tbody>
</table>
FIG. 13 A possible arrangement of inclusive and exclusive reactions in the `<reactions>` element. Note that there are three options for the double differential cross section data: a) store $d\sigma(E)/dE'd\Omega$, b) store $d\sigma(E)/d\Omega$ if the outgoing energy is fixed by kinematics or c) store both $\sigma(E)$ and $P(E',\mu|E)$ separately. In some cases the total cross section $\sigma(E)$ is not defined and one must use one of the other options.

The distributions, etc. (and even the cross section) may have sublibrary or class specific representations.

**Discussion point:**
How does one implement breakup and/or multi-step reactions? The proper use of ENDF’s LR flag scheme is complex. It is used for light element breakup, (n,gf) reactions and reactions which lead
to unstable residuals (e.g., isomers) but whose half-lives are large enough that the residuals must be accounted.

**Resolution:**
This is handled by the `<decayProducts>` element in section II.D

**Discussion point:**
A `<channel>` element could be used by decay data, atomic scattering, thermal neutron scattering and nuclear reaction data to denote the reaction in finer detail than is possible with a simple reaction designator. Within nuclear reaction data it would be used for fast reactions and the resonance region differently. A selection of kinematic types for the channels is given in Table I.

**Resolution:**
The idea is incompletely formed but adds an extra, unneeded, requirement on the class structure of the code reading this data.

**Discussion point:**
We can optionally store the Q value and threshold energies, but these are derivable if one knows the identity of the initial and final state particles. Requiring that they be given in a channel potentially introduces an internal consistency error if the values are not kept in sync with any external material property database.

**Resolution:**
Putting in the Q values in the `<reaction>` is useful, but we shouldn’t take the values seriously. That said, legacy ENDF evaluations specify the Q values for each reaction and these values may be chosen by the evaluator to get the correct threshold behavior of a reaction even if the Q value does not match that derivable from the incoming and outgoing particle masses. Furthermore, often the masses in ENDF are not stored with enough precision to obtain an accurate Q-value. So, if one intends to translate a file back into ENDF, one should specify the Q-value.

**Discussion point:**
Further discussion: What about saying “Default is to derive the Q-value from masses. However, for supporting legacy evaluations we include a deprecated option to specify the Q-value”?

**Discussion point:**
Reconstructed cross sections and angular distributions could go here too

**Resolution:**
Agreed. This also provides a location in the hierarchy for other derived data used in processed data files.

**D. Differential cross sections: \(d\sigma(E)/d\Omega \text{ and } d\sigma(E)/d\Omega dE\)**

There are many cases where it is more convenient to write two-body scattering data as \(d\sigma(E)/d\Omega\) rather than as a separate cross section \(\sigma(E)\) and angular distribution \(P(\mu|E)\) where \(d\sigma(E)/d\Omega = \sigma(E)P(\mu|E)\). Cases include:

- **Thermal Scattering Law (TSL) data**, see Section XIV
- **Large Angle Coulomb Scattering (LACS) data**, see Section XIII
- **Photo-atomic data described with the Klein-Nishina (KN) formula**, see Section VIII.A

Indeed, in the case of large angle Coulomb scattering (LACS) data, the singularities in the Rutherford cross section prevent us from integrating to find the total cross section \(\sigma(E)\). Therefore, we must provide a facility for flagging a reaction as a special parameterized two-body reaction and a facility for storing \(d\sigma(E)/d\Omega\).

**Requirement 19: `<dcrossSection_dOmega>`**

19.1 The actual implementation (which depends on the nature of the described data).
19.2 An optional `<documentation>`
19.3 A flag to denote whether to use relativistic or non-relativistic kinematics when handling this channel
19.4 These are mid-level container elements (see requirements 7).

In the case of Thermal Scattering Law (TSL) data (see Section XIV), it is more convenient to write two-body scattering data as \(d\sigma(E)/d\Omega dE'\) rather than as a separate cross section \(\sigma(E)\) and energy-angle distribution \(P(E',\mu|E)\) where \(d\sigma(E)/d\Omega dE' = \sigma(E)P(E',\mu|E)\). Therefore, we must provide a facility for flagging a reaction as a special parameterized two-body reaction and a facility for storing \(d\sigma(E)/d\Omega dE'\).
Requirement 20: `<dcrossSection_dOmega_dE>`

20.1 The actual implementation (which depends on the nature of the described data).
20.2 An optional `<documentation>`
20.3 A flag to denote whether to use relativistic or non-relativistic kinematics when handling this channel
20.4 These are mid-level container elements (see requirements 7).

E. Integral cross sections: \( \sigma(E) \)

A `<crossSection>` element would be used by atomic scattering and nuclear reaction data. It is analogous to ENDF’s MF=3 or 23.

Discussion point:

It was suggested to give cross sections as ratios, for example ratios to total? This was seen as a way to eliminate sum rule failings.

Resolution:

However, this would intentionally introduce synchronization troubles and require rewriting a lot of code to take advantage of. In addition, we have substituted one sum rule (summing to a cross section) with another (summing to 1), so we really haven’t gained anything.

Requirement 21: `<crossSection>`

21.1 A `<crossSection>` element is either:
   21.1.1 At least one version of the data containing an `<interp1d>` element with a dependent variable (the cross section itself) given in units of area and independent variable (projectile’s incident energy) in units of energy. The first energy point could be (real or effective) threshold or the lowest energy supported by the encapsulating evaluation. A `<crossSection>` is assumed to be zero outside of the specified energy region
   21.1.2 A link to the resonance region that one must reconstruct in order to retrieve the cross section data tables
21.2 These are mid-level container elements (see requirements 7).
21.3 A `<crossSection>` element may have `<documentation>`.
21.4 An optional `<PURR>` table (see section XV)

Requirement 22: `<styles>`

22.1 Shall provide a mechanism for denoting on a per–transportable particle basis, energy group structures.
22.2 Shall provide a mechanism for denoting on a per–transportable particle basis, angular group structures.
22.3 Shall provide a location for storing the Doppler-broadened target temperature
22.4 Shall provide a mechanism for uniquely naming and referring to style information.
22.5 Shall provide a mechanism for storing fluxes for flux weighting of projectiles.
22.6 Should provide a way to name a style or group of styles.

F. Defining data styles

An evaluation may contain one or more variations of processed or derived data in addition to the original evaluated data. Processed data are usually processed in the similar way (for example, in the same group structure or with the same flux weight), so it makes sense to define “styles” of data. For processed and evaluated data, the styles include, but are not limited to, target temperature, group structure (in energy and angle) and particle flux weights. Each style should be uniquely named so that it can be referenced in the lower level data structures.

G. Derived reactions

There are many instances of derived data that we wish to store in an evaluation but which do not fit neatly in the `<reactions>` or `<resonances>` tree. For example:

1. Inclusive reactions such as:
   (a) Total cross section (a requirement for neutrons), allows evaluated version with covariance and the total may not exactly equal the sum of partial cross sections
   (b) ENDF sum rules (e.g., MT103 vs. MT600 – MT649)
   (c) Lumped cross sections for lumped covariances (should be deprecated!)
2. Gas production cross sections
3. Damage cross sections
4. Activation cross sections
5. Particle and isotope yields

The total cross section is especially interesting: it is simple to measure and evaluators can both fit the data
Fig. 14 Layout of the `<derivedReactions>` element.

and model it with optical model calculations with high fidelity. The evaluation of the total cross section is usually quite solid with high quality covariances. However, there is no guarantee that if one sums up the partial cross sections of an evaluation that one will retrieve the total cross section. Therefore, while most evaluations will provide a total cross section (and therefore place it in the `<derivedReactions>` branch, a processing code will discard the evaluated total cross section and recompute the total from the sum of partials in order to ensure unitarity.

In GND, inclusive reactions are encoded in a `<summedReaction>` element. This element includes a cross section (and this may be connected to covariance data). Additionally, there is a list of links to the reactions which are meant to be summed together to match the cross section data in the element. This element is used to implement all of the ENDF sum rules in Section 0.4.3.11 of the ENDF format manual [Trkov 2009].

The requirements for derived reactions are:

**Requirement 23: `<derivedReactions>`**

23.1 Reactions shall have common sense reaction designators (e.g., “total” or “absorption”).

23.2 A (likely derived) cross section which includes list of links to the reactions whose cross sec-
tion is meant to match the cross section tabulated here. They are mid-level containers (see requirements 7).

23.3 Optionally `<reactionProducts>` list with `<product>`s to store yields.

**Discussion point:**
Is this a good place to put ENDF’s 6 time-group beta delayed fission data? That data are “derived”, but the derived data are so far removed from the original that it may be impractical to link to it all – we would have to reference a huge part of the FPY and decay sublibraries for each actinide!

**Discussion point:**
Should we also include weights in the list of links? This would add flexibility, but it is not clear how they would get used outside the case of lumped covariances. Incidentally, lumped covariances should be depreciated as they are not very useful.

**SUGGESTED TEST:** As `<derivedReactions>` have lists of links to the reactions they are derived from, we should check that the links are valid and, for example, `<crossSection>`s sum as promised.

**Discussion point:**
Below, all reactions are lumped together in `<reactions>` and `<derivedReactions>` elements. Should we also group ‘production’ reactions (i.e., the sum of all reactions that produce product ‘x’)? Activation libraries could make heavy use of production cross sections.

**Resolution:**
The `activationFlag` should be used to denote an evaluation comprised of activation data. The production cross sections themselves are derived from cross sections in the `<reactions>` branch and stored in the `<derivedReactions>` branch. See subsection IV.G for more detail.

**Discussion point:**
Do we allow production cross sections?

**Resolution:**
A mis-filed production cross section is very dangerous. In the simplest case, a production cross section is a product of a particle’s multiplicity and the cross section for that particle’s production. For an (n,2n) reaction, this means a potential factor of two error. For gamma production, this could mean a factor of a 100. Because the units on a production cross section and a regular cross section are the same so there may be no way to tell if one mis-filed a production cross section, leading to crazy energy balance bugs. However, with special markups, this should be avoidable.

**Discussion point:**
Further discussion: should we have an explicit `<productionReaction>` element, so that production cross sections can be given unambiguously?

H. Resonances

The resonance region is a surprisingly complicated part of any evaluation. A resonance is a sharp feature in a cross section caused by a target capturing a projectile and forming a compound nucleus. The compound nucleus later decays by particle emission. For charged particle incident reactions, there are typically very few resonances in the observed cross sections because the resonances are suppressed by Coulomb effects. On the other hand, in a neutron induced reaction there can be thousands of observed resonances.

In ENDF and in our proposed format, the parameters defining resonances are stored, rather than the reconstructed cross sections. This is done partly for a more compact presentation of the data and partly for a more exact representation of the data. In this section, we describe resonance data as in ENDF’s MF=2, MT=151. Our proposed hierarchy is given in Figure 15. These data describe resonances that are observable in neutron cross sections for $E = 0$ eV → 100’s keV (or higher for charged particles). In ENDF, these data are only used for neutrons, but should be legal for charged particle reactions and even photonuclear data.

To understand the hierarchy of resonance data, it is helpful to understand a little about R matrix theory. In it, the universe is divided into the inside and outside of a spherical box. Inside the box is the reaction zone, where all the interesting nuclear (or other) reaction business occurs (see Figure 16). We have little chance of modeling what goes on in the box correctly without a lot of work. Outside the box we write all incoming and outgo-
FIG. 15 Our proposed resonance data hierarchy.

ing relative two-body scattering states in a basis of analytic wave functions, usually taken to be free ones. We then match wave functions on the box boundary. This matching is done in a clever way involving Bloch surface operators on the box boundary and from this we arrive at a Green’s function of the projected Bloch-Schrödinger equation, also known as the R matrix. The elements of the R matrix are:

$$R_{cc'} = \sum_{\lambda} \frac{\gamma_{\lambda c} \gamma_{\lambda c'}}{E_{\lambda} - E}$$  \hspace{1cm} (2)

The factor $\gamma_{\lambda c}$’s are the reduced widths for channel $c$, $E_{\lambda}$ becomes the resonance energy (it is a pole in the Laurent series expansion of the Green’s function) and $\lambda$ is the resonance (pole) index. The channel index $c$ contains all the quantum numbers needed to describe the two-particle state and all of those quantum numbers are described in the <channel1> and <spinGroup> element markups below. The channel index may refer to an incoming or an outgoing channel.

With the R matrix, it is possible to compute exactly the channel-channel scattering matrix $U_{cc'}$:

$$U_{cc'} = e^{-i(\phi_{c} + \phi_{c'})} \sqrt{P_c} \sqrt{P_{c'}} \times \left\{ [1 - R(L - B)]^{-1} [1 - R^*(L^* - B)] \right\}_{cc'}$$  \hspace{1cm} (3)

where the logarithmic derivative of an outgoing channel function is

$$L_c \equiv a_c \frac{O_c'(a_c)}{O_c(a_c)} = \left[ r_c \frac{\partial \ln O_c}{\partial r_c} \right]_{r_c = a_c}$$  \hspace{1cm} (4)

and we write

$$L_c = S_c + i P_c.$$  \hspace{1cm} (5)

The penetration factor is $P_c = \Re L_c$ and the shift factor is $S_c = \Im L_c$. Both take their names from their function
in the simple complex square well scattering model. \( \varphi_c \) is the phase factor

\[
\varphi_c \equiv \arg O_c(a_c) = \arctan \frac{\Re O_c(a_c)}{\Im O_c(a_c)}
\]

The constant \( B_c \) is the so-called “boundary parameter” which must be specified to correctly compute the scattering matrix, but is not always clearly given. We note that the non-trivial nature of even the hard-sphere penetrability, shift and phase give rise to the potential scattering behavior of the cross sections even when there are no resonances.

With the scattering matrix, one can compute all channel cross sections, the total cross section, and all angular distributions. The angle integrated cross section can be written as sum over all entrance channels \( \alpha = \{ \alpha J_1 \ell \} \) and exit channels \( \alpha' = \{ \alpha' J_2 \ell' \} \) that lead from partition \( \alpha \) to \( \alpha' \):

\[
\sigma_{\alpha \alpha'} = \pi \lambda_2^2 g_c |\delta_{\alpha \alpha'}|^2
\]

The total cross section for channel \( c \) is

\[
\sigma_c \equiv \sum_{\alpha \alpha'} \sigma_{\alpha \alpha'} = 2 \pi \lambda_2^2 (1 - \Re U_{cc})
\]

The factor of \( g_c \) is the probability of getting the correct \( J \) from the spins of the collision partners (according to Frohner) and is \( g_c = (2J + 1)/(2\ell + 1)(2\ell' + 1) \).

The Blatt-Beidenharn equation (Blatt, 1958) is used to construct the \( d\sigma_c/d\Omega \) for two body channels in the center-of-momentum. In the ENDF formatted libraries, the Blatt-Beidenharn equation is usually used for elastic channels. Although \( d\sigma_c/d\Omega \) can be written as a Lorenz covariant quantity, we will write the outgoing dependence on angle in their pair center of mass frame and the incident energy in the laboratory frame.

For particles with arbitrary spin, we have

\[
\frac{d\sigma_{\alpha \alpha'}(E)}{d\Omega} = \frac{1}{k^2(2i + 1)(2\ell + 1)} \sum_{s, s'} \sum_{L=0}^{\infty} B_L(\alpha s, \alpha' s'; E) P_L(\mu)
\]

and

\[
B_L(\alpha s, \alpha' s'; E) = \frac{(-)^{s-s'}}{4} \sum_{J_1, J_2, \ell_1, \ell_2} \sum_{J_1', J_2', \ell_1', \ell_2'} \tilde{Z}(\ell_1 J_1 \ell_2 J_2 s L) \tilde{Z}(\ell_1' J_1' \ell_2' J_2' s' L)
\]

\[
\times (\delta_{\alpha \alpha'} \delta_{\ell_1 \ell_1'} \delta_{\ell_2 \ell_2'} - U_{\alpha \ell_1 s, \alpha' \ell_1' s'}(E))^* (\delta_{\alpha \alpha'} \delta_{\ell_2 \ell_2'} \delta_{s s'} - U_{\alpha \ell_2 s, \alpha' \ell_2' s'}(E))
\]

\[
= \frac{(-)^{s-s'}}{4} \sum_{\ell_1, \ell_2, s, l_1, l_2, s', l_1', l_2'} \tilde{Z}(\ell_1 J_1 \ell_2 J_2 s L) \tilde{Z}(\ell_1' J_1' \ell_2' J_2' s' L)
\]

\[
\times \delta_{s s'} \delta_{\ell_1 \ell_1'} \delta_{s s'} \delta_{\ell_2 \ell_2'} \delta_{s s'} \delta_{\ell_2 \ell_2'} (\delta_{c_1 c_1'} - U_{c_1 c_1'}(E))^* (\delta_{c_2 c_2'} - U_{c_2 c_2'}(E))
\]

where

\[
\tilde{Z}(\ell_1 J_1 \ell_2 J_2, s L) = \sqrt{(2\ell_1 + 1)(2\ell_2 + 1)(2J_1 + 1)(2J_2 + 1)(\ell_1 \ell_2 00, L0)W(\ell_1 J_1 \ell_2 J_2, s L)}
\]

and \( W(\ell_1 J_1 \ell_2 J_2, s L) \) is a Racah coefficient.

We use the notation \( \sum_c = \{ \alpha \ell s J \} = \sum_{\ell} \sum_s \sum_J \). The ENDF manual uses the notation \( \sum_c = \sum_{\ell} \sum_s \), so it
needs an extra sum over $J_1$ and $J_2$ (Irkov, 2009). This is detailed in several places including (Blatt, 1958; Descouvemont, 2010; Froehner, 2000; Lane, 1958). Given the mathematical completeness of the theory, it is no surprise that we mostly just view the R matrix parameters as simple fit parameters and then essentially get all of this for free.

**Requirement 24: <resonances> element**

24.1 Optional documentation
24.2 A list of the channels referred to in this evaluation. Traditional ENDF SLBW, MLBW and Reich-Moore formats support only capture, elastic, fission, total and a catch-all competitive channel. The R matrix formalism can support any two-body final state.
24.3 Optionally one or more resolved resonance region (RRR), although multiple RRR is deprecated
24.4 Optionally an unresolved resonance region (URR)
24.5 Upper and lower energy limits delineating the range of applicability of the resonance region.

Finally, we comment that the R matrix approach works for any two-body reaction, relativistic or not, as long as the incoming and outgoing relative states can be clearly defined. In the nuclear data community we often forget this fact. As a result, the ENDF format never really had the ability to represent charged particle data in an R matrix inspired form, reducing the quality and scope of data available to several communities who need it:

- Inertial Confinement Fusion community needs all sorts of charged particle incident data
- Astrophysical community needs the ($p, \gamma$) reaction among many others
- For Nuclear Resonance Fluorescence, need to support ($\gamma, \gamma'$) data
- FRIB Support: all of the FRIB experiments will be with beams of rare isotopes. Since you can’t make a target of neutrons, nearly all of the beam-target experiments will be charged-particle reactions, probably on hydrogen, deuterium and helium targets. There has been some advances in the application of R-matrix methods to these reactions.

**Discussion point:**
By allowing the possibility of creating an evaluation in which one specifies the URR without an RRR, we allow the possibility of an evaluation in which we do not know any resonances, but only the average parameters (perhaps via systematics). This was attempted in the original ENDF 240Am evaluation which has since been superseded.

**FIG. 17 Our channel element.**

1. Designating channels

As discussed above, a channel is the partition (the target and projectile) and the set of quantum numbers needed to completely designate the incoming state. Our scheme is shown in figure 17. Therefore, at the very least a <channel> element must contain the quantum numbers of the state and the target and projectile. Because we anticipate reconstructing the resonances back into pointwise cross sections for plotting among other things, we should also connect to the <reaction> element (or at least its designator). For backwards compatibility with ENDF, we also need the MT and Q value.

Now, because resonances are a specific quantum state in a compound nucleus, not all channels can couple to a particular state. Therefore it makes sense to group together channels with similar quantum numbers into “spin groups” as in ENDF. These leads to denser tables of resonances and so more efficient storage. As in ENDF, we recommend adding a <spinGroup> markup containing all <channel>s with a common $J$ and $\Pi$.

**Requirement 25: <spinGroup>**

25.1 The format shall have a spot for the $J$ an $\Pi$ quantum numbers
25.2 The format shall have a list of <channel>s

Both the RRR and the URR should share the same master channel list. This aids in reconstruction since the number and kind of channels do not change with energy unless a threshold opens up. Also, to aid in reconstruc-
tion and complete specification of the “box” that defines the R-matrix, we must specify the boundary parameter $B_c$ and the box radius $a_c$. We comment that in some older evaluations, the scattering length $a_c$ is energy dependent and this is used as another mechanism to adjust the potential scattering contribution to the resonances.

Conventional R-matrix approaches for neutron channels assume that the outgoing waves are free waves hitting a hard-sphere. Newer approaches to the R matrix use alternative outgoing waves, for example Gurbich’s (2014) uses optical model distorted waves. In addition, some channels (outgoing gamma and fission in particular) may use phenomenological shifts, penetrabilities and phases. Therefore we should allow the evaluator to override the phase, shift and penetrability factors.

**Requirement 26: `<channel>`**

**26.1** The reaction designator; for resonances, this also specifies the “partition” (see Lane and Thomas (1958) IV.A. It is expected that this designator maps correctly onto one in the `<reactions>` list, otherwise there may be problems when reconstructing resonances. The reaction designator can also specify the excited state of the residual nucleus or the other ejected particle.

**26.2** The ENDF MT, if applicable (deprecated)

**26.3** All `<channel>`s must have a link to a `<reaction>` in the outer `<evaluation>`’s `<reactions>` tree so that it is explicit where the data for this channel will be placed after the resonances are reconstructed. This `<reaction>` could be some sort of empty placeholder.

**26.4** All quantum numbers needed to uniquely specify the reaction, this is needed for resonances as well. In particular, the spin $s$ of the channel, the orbital angular momentum $l$, the total angular momentum $J$ and any other quantum numbers. This eliminates a degeneracy implicit in the ENDF format and helps with the quality assurance of the data.

**26.5** Configurable channels to denote whether corresponds to actual two-particle final state or effective one (as in fission or competitive channels). Only two-body channels can be used to compute angular distributions; need to be able to flag “effective” channels

**26.6** List the spin $s$ for each resonance (resolves an ENDF ambiguity).

**26.7** Boundary parameter $B_c$

**26.8** Channel radius vs. true channel radius

**26.9** Sign of reduced width

**26.10** Optionally the Q value of the reaction as it can help the evaluator. As this can be derived from the reaction products, it is strictly not needed as it is redundant.

**26.11** To override the defaults, optionally specify

- **26.11.1** phase $\varphi_c(E)$
- **26.11.2** shift $S_c(E)$
- **26.11.3** penetrability $P_c(E)$
- **26.11.4** hard-sphere radii $a_c$ (with potential dependence on energy). Likely need to be able to break it into multiple regions so that, for example, the RRR can have a constant one while the URR can have an energy dependent one.

**SUGGESTED TEST:** Check that spin groups and channels are consistent between the RRR, URR and channel specifications.

**Discussion point:**

We will need tests to ensure consistency between the `<channel>`s, the `<reaction>`s in the `<reactions>` and between the `<channel>`s and the `<RRR>` and `<URR>` columns.

**Discussion point:**

Is our approach enough to handle ($\gamma$,γf) and/or fission reactions through class II states?

**Discussion point:**

Would configurable ignored or collapsed channels (like γ ones in Reich-Moore approximation)? Or is the Reich-Moore approximations on photons the only one that makes sense in practice?

**Discussion point:**

Would user-definable (possibly fake) quantum numbers be useful? We would then need to define whether to combine using angular momentum adding rules (for Blatt-Biedenharn) or incoherently. For deformed nuclei, the $K$ quantum number may require this feature.

**Discussion point:**

The channel wish list is very big. That said, a `<channel>` has all of the attributes of the `<reaction>` element. Does it make sense to completely separate the `<channel>` concept of the `<reaction>` element?
Resolution:
No, our <reaction> concept is very broad compared to the specific nature of the <channel> concept. Nevertheless, some codes that use these data may find it useful to design a parent class from which the channel and reaction classes might inherit.

Discussion point:
For charged particle channels, it is important to specify the correct mass and ionization state of the target atom so that electron screening and target recoil can be properly accounted for. How do we handle targets in a plasma environment?

2. Resolved Resonances

The ENDF format supports several different approximations to the collision matrix from R matrix theory in Eq. (3). In addition, we should support the addition approximations discussed by Fröhner [Froehner 2000] and in CALENDF [Ribon 1986]:

1. Pure potential scattering with either hard sphere or tabulated energy and/or $\ell$-dependent scattering radius. Allows cross section and angular distribution calculation.


3. ENDF style SLBW. Allows only cross section calculation. (deprecated)

4. Multi Level Breit-Wigner (MLBW). Allows cross section and angular distribution calculation. CALENDF refers to this approximation as the Multi-Niveau Breit-Wigner (MNBW) format [Ribon 1986].

5. ENDF style MLBW. Allows cross section and angular distribution calculation for elastic reactions.


These approximations all require the same resonance energies $E_\lambda$ and resonance widths $\Gamma_\lambda = 2P_\lambda \gamma_\lambda^2$.

FIG. 18 Sample table of resonance parameters.

| $E_\lambda$ | $\Gamma_{\text{tot}}$ | $\Gamma_{\text{c0}}$ | $\Gamma_{\text{c1}}$ | $\Gamma_{\text{c2}}$ | ...
|------------|-----------------|-----------------|-----------------|-----------------|...
| 1.23       | 9.433           | 0               | 2.33E-03        | 7.1             | ...
| 1.46       | 4.833           | 0               | 2.33E-03        | 4.6             | ...
| 3.45       | 1.78            | 1.78            | 0               | 0               | ...

Inevitably, either because of deficiencies in the approximation use, because of resonance missed in the evaluation process, or simply because of the tails of distant resonances, we must take action to correct the reconstructed cross section and angular distributions. This is covered in detail in the next section (section IV.H.4). However, it is possible to make some corrections directly to the R-matrix. The easiest approach is just to add fake resonances to the resonance region just outside the energy range of the RRR. The tails of the fake resonances serve to correct the shape of the R-matrix inside the energy range of the RRR. This trick is commonly used to recreate the $1/v$-like behavior of the neutron capture cross section below thermal energies. The use of fake resonances in this way must be flagged. Alternatively, one can add
a background R matrix to mimic the effects of distant resonances (replacing $R_{cc'} \rightarrow R_{cc'} + R_{cc'}^{\text{back}} \delta_{cc'}$). One can do this with a complex-valued energy dependent function for $R_{cc'}^{\text{back}}(E)$ or with parameterizations provided by Fröhner (Froehner, 2000) or given in the SAMMY manual (Larson, 2006).

**Requirement 28: Background R matrix**

28.1 The `<backgroundRMatrix>` element shall have an option flag to denote the background R-matrix correction scheme, equivalent to ENDF’s KBK flag.
28.2 For KBK=0, only need dummy resonances so no other information needed
28.3 For KBK=1, the `<backgroundRMatrix>` contains a set of complex-valued `<interp2d>` tables for $R_{cc'}^{\text{back}}(E)$, one for each channel
28.4 For KBK=2, the `<backgroundRMatrix>` contains a set of eight parameters for each channel. These parameters are needed for SAMMY’s logarithmic parameterization (see (Larson, 2006; Trkov, 2009)).
28.5 For KBK=3, the `<backgroundRMatrix>` contains a set of four parameters for each channel. These parameters are needed for Fröhner’s parameterization (see (Froehner, 2000; Trkov, 2009)).

**Discussion point:**
Although the Adler-Adler approximation is not currently used in major data libraries and its use in the ENDF format is deprecated, the approximation offers some practical advantages for efficient Doppler broadening. Similarly, the multipole expansion of resonances could offer similar gains. If the use of either increases, we will revisit the inclusion of features that support these formats.

**Discussion point:**
There are two complementary approaches to expressing the R-matrix: Kapur-Peiers and/or Wigner-Eisenbud. Both approaches use different boundary parameters $B_c$. They are mathematically equivalent, but the RRR approximations in ENDF all use Wigner-Eisenbud formulation. Should we support Kapur-Peiers as well?

**Resolution:**
No, because in Kapur-Peiers, one sets the boundary constant $B_c = L_c$. This leads to a complex pole $E_\lambda$, forcing us to mix data types (complex vs. float) in the `<table>` element in the `<RRR>` element.

**Discussion point:**
Follow-up on the previous point: so far the low-level data containers have no requirement for complex data types. Do we need to add them? The only place they show up in GND right now is inside large-angle Coulomb scattering, but there the data are divided into two separate arrays (one real, one imaginary).

**Discussion point:**
Fröhner (Froehner, 2000) suggests storing width amplitudes $\gamma_{\lambda c}$ instead of widths. This avoids sign confusions, they are not energy dependent and they do not vanish at threshold.

**Resolution:**
To do so would mean that we’d need to have an excellent grasp on what the penetrabilities really mean for $\gamma$ (we know what they are for neutrons and charged particle channels). We’ll still need $\Gamma_{\lambda c}$ for fission and competitive channels since there is no notion of penetrability in those cases. We’d also need to know the relativistic version of the penetrabilities. One could tabulate effective penetrabilities in the `<channel>` such that $\Gamma_{\lambda c}$ comes out right.

**Discussion point:**
Channel major arrangement or maintain resonance major arrangement? What I mean is, are the rows in the “table” mean one row/resonance with all the channels as columns as in ENDF? Or do we switch to having a list of channels at the top with a list of resonances associated with each channel? Either way the matrix $\Gamma_{\lambda c}$ is sparse. Perhaps channels that are present purely for the purpose of getting the potential scattering part of the cross section correct should be skipped.

**Discussion point:**
Should we flag fake resonances beyond what is done to correct the R-matrix? Namely, should there be
additional flags for resolved resonances in the RRR that appear as real resonances.

Discussion point:
Do we need to clarify rules for the resolved and unresolved region widths for threshold reactions.

3. Unresolved Resonances

Above the RRR, we know that there are still distinct resonances, but they are too close together to be experimentally resolved. In this Unresolved Resonance Region (URR) we only know statistical properties of the resonances. Therefore, what is stored is not the resonance parameters, but ensemble averages of them: averaged first over ensembles of imagined resonances, then over the width distributions of the resonances. The widths are assumed to be distributed according to a $\chi^2$ distribution with a channel dependent number of degrees of freedom. In the ENDF format (Trkov, 2009), the resonances are assumed to be in the SLBW approximation before averaging leading to the particular parametric form of the cross sections in the ENDF manual. However, CALENDF and other codes can use other parameterizations. Therefore a flag denoting the approximation of the R-matrix to use should be given.

**Requirement 29: Unresolved resonance region (URR)**

29.1 Upper and lower energy limits delineating the range of applicability of the unresolved resonance region.
29.2 Need number of degrees of freedom associated with each channel in the channel listing
29.3 Need a `<table>` of URR parameters, organized by `<spinGroup>`. This table must include columns for incident energy, mean level spacing, average widths for all channels.
29.4 ENDF assumes SLBW, allowing the construction of average cross section and PURR tables. This is a somewhat arbitrary restriction that is removed in CALENDF (Ribon, 1986). This URR format should allow all approximations that are supported for the RRR.
29.5 An `<axis>` and interpolation details to determine how to interpolate in incident energy among the average parameters.

Both the PURR module in NJOY (MacFarlane, 2012) and the PURM module in AMPX (Dunn, 2002) can compute cross section probability distributions $P(\sigma_x|E)$ for all $x \in [\gamma, \text{el, tot, f, ...}]$. Upon reconstruction, the PDF should get placed in the appropriate `<reaction>` as deserved data.

4. Correcting cross sections and distributions with background data

Many of the resonance approximations have shortcomings, in either the shape of the reconstructed cross sections or in the fact that outgoing angular distributions of emitted particles are untrustworthy or nonexistent. Indeed, one can never reconstruct the angular distributions of any emitted particles from fission resonances. Therefore, we need multiple schemes to amend or fix there reconstructed data tables from the RRR and URR. In figure 19 we provide an overview of things to fix reconstructed data.

In figure 19 we show a `<backgroundReaction>` element that contains all the things needed to fix the reaction linked to in the `<linkToReaction>` attribute. Additionally, the `addOrReplaceFlag` instructs the person processing the resonance data whether to add the cross section within `<backgroundReaction>` to the reso-
nances or to replace the resonances entirely (this is useful for some older data). To accommodate for missing or mishandled outgoing particle data, we also add a `<reactionProducts>` element. Since several of the ENDF resolved resonance approximations DO NOT support angular distributions, the ENDF format provides a mechanism to store the “background” angular distributions separately. These “background” angular distributions are naturally go in the `<reactionProducts>` element of the `<backgroundReaction>` element. Upon the reconstruction of the resonance region into pointwise data, the reconstructed data should point back to both the resonance data and these background data.

**Requirement 30: The `<backgroundReaction>` element**

30.1 Background cross-sections should not given for (n,tot) and should be associated with the actual reaction to which the background cross-section is added. Otherwise we are adding a potential source for inconsistencies (background partial cross-sections not summing to the background total cross-section) or double counting.

30.2 The background cross sections must be associated with the resonance region for which they are the background. The obvious place is in the `<crossSection>` element of the `<backgroundCrossSection>` element.

30.3 The `<backgroundReaction>` shall have a flag to denote whether the cross section data within the element is meant to be added to the reconstructed resonances or replace them.

30.4 The `<backgroundReaction>` shall have a link to the reaction that these data will go to upon reconstruction.

30.5 The `<backgroundReaction>` shall have an option `<reactionProduct>` element to store all of the outgoing particle data not reconstructed from the resonances associated with this reaction in the resonance energy range.

We feel that the background cross sections and related data should be kept with the RR, not the high energy file so that the association is explicit. This is a different arrangement than in legacy ENDF where background cross sections are kept with the fast region cross section.

**SUGGESTED TEST:** Check for consistent energy ranges between RRR, URR and background reaction data.

**Discussion point:**
Primary gammas are the gammas produced when a neutron captures into a compound nucleus and the nucleus decays via gamma emission to a discrete state below the neutron separation energy. This state de-excites by a gamma cascade. As the neutron moves up in incident energy, the primary gamma’s energy appears to move with it. In ENDF, primary gammas are a mess, with two separate implementations (MF=12 and MF=6). Both ENDF approaches are kludges. A multi-step R matrix approach could handle it.

**Resolution:**
Although Lane and Thomas provide a mechanism for doing this [Lane 1958], GND provides a `<decayProduct>` markup and a particle database, both which can easily accommodate primary gammas and their related gammas.

**V. GROUPING TOGETHER EVALUATIONS**

There are many instances where we will need to group together evaluations either to make a collection of things such as an entire sublibrary or to connect evaluations to make a new “effective” evaluation (see subsection V.B).

**A. Defining a collection of evaluations**

There are many cases where one simply wants to collect together evaluations and the like:

- When defining a (sub)library
- To collect web retrievals under one element
- To produce a valid XML file containing more than one evaluation (the equivalent of an ENDF TAPE).
- To provide an index of a set of files including say associated particle definitions and covariance data

To enable this, we must provide a top-most level markup as shown in Fig. 20.

The requirements for such an element are simple:

**Requirement 31: `<library>`**

31.1 An optional `<documentation>`
31.2 A list of evaluations or links to evaluations
31.3 A list of covariances or links to covariances
31.4 An optional particle database or link to a database.

**B. “Effective” or “Meta” evaluations**

There is a relatively common need to “glue” together evaluations to make new “effective” or “meta” evalua-
FIG. 20 Grouping evaluations, covariances and/or particle properties.

- In LANL’s MCNP code system, the xsdir file allows one to connect the thermal neutron scattering data with the neutron nuclear reaction data and even various high energy models such as CEM. See, for example, Figures 21 and 22.

- The LLNL transport codes AMTRAN and Mercury both allow one to define target macros to describe the material in a zone.

- ORNL’s SCALE package contains a pre-built material composition database.

- At AECL, there is another, similar, facility to connect thermal neutron scattering data at different temperatures and even different phases of the tar-
get material.

There are other uses for being able to connect evaluations together:

- Defining elemental evaluations
- Grouping data on same target, but heated to different temperatures
- Defining generic fission fragments through a weighted average of fission fragment evaluations
- Putting together the parts of a TSL evaluation at fixed temperature, but including all the scatterers.
- Defining common material definitions. This helps answer the question “Which concrete?”

Ideally, these could be shared but rarely are because of the wildly different formats used by various projects. The need for “gluing” together evaluations is so common that we should seriously consider supporting it.

The idea of a `<metaEvaluation>` is straightforward. One uses a set of `<axis>` elements to define the grid in some parameter space one wishes to populate with evaluations. The `<axis>`’s could be temperature, incident energy, pressure, etc. The `<axis>` element defines the boundaries in the parameter space. The `<axis>` elements also define the interpolation scheme to be used in that parameter’s direction, but in practice the interpolation information will probably be ignored because each project defines their own rules for stepping up in temperature, etc. These ideas are illustrated in Figures 23 and

![Figure 21: Gluing together different models from different energy regions.](image)

**Figure 22** Gluing together thermal neutron scattering with the higher energy nuclear reaction data. Note that nesting `<metaEvaluation>`s can make the implantation of this quite simple.

**Requirement 32: `<metaEvaluation>`**

32.1 An projectile attribute to define what projectile this `<metaEvaluation>` is only valid for (say TSL+fast gluing only for neutrons).

32.2 `<axis>` elements to define the grid in which the evaluations will be inserted

32.3 `<referredEvaluation>` which links to an `<evaluation>` or another `<metaEvaluation>`.

32.4 Outside of parameter ranges in `<axis>` tags, the `<metaEvaluation>` does not exist

32.5 `<metaEvaluation>` only valid for listed projectile

32.6 Need tests to make sure every region in `<axes>` covered by a `<referredEvaluation>`.

**Discussion point:**

Is it possible to use say atomic weights instead of stoichiometricFraction to specify fractional composition of a material? This would simplify use in several transport code input decks.
FIG. 23 Sample `<metaEvaluation>` specification, in this case for water. This file requires another `<metaEvaluation>` to specify the composition of dissociated water into the elements hydrogen and oxygen. These then require other `<metaEvaluation>`s to specify the elemental composition of H0 and O0 in terms of their isotopics.

Resolution:
No, this should be done at the `<evaluation>` level so the `nativeData` information is associated with the evaluation file itself and not somewhere else.

Discussion point:
Is there a need for a separate `<metaTarget>` concept to handle arbitrary projectiles so we needn’t maintain 7-8 different (but nearly identical) element specifications?

Resolution:
Good point. Maybe allowing `any` or `*` as a projectile would serve this purpose. Alternatively, we could make the `projectile` attribute optional and if it is not present then the `<metaEvaluation>` is valid for all projectiles. Either way the links to the actual evaluation become meaningless. This requires some thought. Perhaps the resolution is to pre-make the elemental `<metaEvaluation>`s for the standard targets with fake URLs. Then users can swap-n-replace them with the correct URLs for their own needs. However, if one of the `<axis>` elements covered incident energy, there is a question of how to handle Q values and different channels opening up.

Discussion point:
It was felt at the WPEC SG38 meeting in May 2014 that implementing `<metaEvaluation>`s amounts to “scope creep”, meaning that this capability was not included in the original requirements. That said, it was generally agreed that this was a useful idea.
Resolution:
The implementation of metaEvaluation specifications should be deferred until the main infrastructure is “in place” and we can devote time to creating the specifications and creating translators back and forth between LANL’s xs-dir format and other related formats at other data centers.

VI. COVARIANCE DATA

Users of nuclear data need covariance data to quantify uncertainty on the metrics of importance in their specific application. These metrics (such as $k_{eff}$ in a criticality calculation) may have a deep dependence on the underlying data. Our users actually use the covariances with deterministic group-wise methods (using the “Sandwich formula” below) or with Monte Carlo techniques. We must do what we can to simplify both modes of covariance use. Our users are also wise about their use of nuclear data and often want to just plot a data set and see if the uncertainties make sense. We must do what we can to simplify this as well.

The legacy ENDF format stores nuclear data covariance in with a very complex set of schemes: the ENDF manual (Trkov, 2009) take over 80 pages to describe seven distinct types of data. Arguably, there should be one “simple” format to govern them all, after all a covariance matrix is, at its heart, just a matrix.

That said, we must deal with covariances not just within an observable, but across observables and evaluations. These covariances can also be quite large, far exceeding the size of the evaluations to which they refer. To enable all of these use cases, we must allow the storing of covariance data in several places, either with a particular data set (for covariance within a dataset), within an evaluation (for cross-reaction or cross observable covariance), or in an external file (for cross material or cross library covariance). In Fig. 25 we illustrate these points.

In Fig. 25, we have one large covariance matrix in the center of the figure composed of several blocks, each consisting of either a self- or cross- covariance in a schematic evaluation. Along the diagonal in orange are the self-covariances. A self-covariance is the covariance matrix for the data within a single data element. These are stored with their corresponding data element, in this case a reaction’s <crossSection>. The cross-covariances between the different data elements are shown in yellow. These connect to two different sets and so are not stored with the cross section data. Rather, they are stored in their own data tree in a <covariances> container. These blocks are not covariances in of themselves. The blocks in white are a reflection of the yellow ones since the full covariance matrix is symmetric.

Requirement 33: Where to put covariance data

33.1 Self-covariance data should be stored in the <listOfDataRepresentations> of a data element
33.2 Cross-covariance data shall be stored in a <covariances> container, either within an evaluation or in an external file with <covariances> as the root node.
33.3 <covariances> containers shall:
   33.3.1 have a spot for optional documentation
   33.3.2 allow one or more covariances

When it comes time to use covariance data, one must take care to use all relevant data. As a specific example, suppose we want to compute the ratio of the capture cross section to the fission cross section. The two cross sections might be correlated in that they have a cross-reaction covariance for which we must account. Therefore we must correctly use this cross-reaction covariance in addition to the self-covariances for the fission and capture cross sections. To ensure that this connection is noticed in practice, we require that all data correlated through cross-covariances be linked as illustrated in Fig. 26.

Requirement 34: Cross-covariance linkage

34.1 Cross-covariance data shall have links to the element of the data correlated.
34.2 The correlated data element will have links back to the cross-covariance data.

Discussion point:
Should we enable storing confidence interval information? Confidence intervals can be computed for any PDF.

A. Covariance Definitions

Before proceeding to state the requirements, it is useful to review the properties of covariances and related objects.

When we measure a quantity $x_i$, we assume that we do not actually get the “true” value given by Nature, but rather one sample from a probability density function (PDF). We note that the vector of observables $\mathbf{x} = (x_0, x_1, ...)$ do not need to be limited to a single observable within a reaction or even within an evaluation. Depending on the nature of the observable(s), the PDF might be Normal or Log-Normal (Zerovnik, 2013).
FIG. 25 Relationship of self-covariance and cross-reaction covariances, illustrating where they are to be stored in data files. The central covariance matrix is divided up into rectangular blocks, corresponding to different observables. The self-covariances are stored with the data elements in the tree on the left and the cross-covariance parts are stored in a separate list on the right.

or something else. For our purposes, we will assume that the PDF is either Normal or Log-Normal since the Central Limit Theorem guarantees that in the limit of large numbers of samples the peak of any PDF can be well approximated by a Normal distribution. We also include Log-Normal as an option since it forces values of an observable to be positive definite but otherwise behaves like a Normal distribution (Zerovnik, 2013).

For a quantity \( x \), its PDF has an expectation value of
\[
\langle x \rangle = \int dx \text{PDF}(x) x
\]
and this would be stored in the ENDF file. The uncertainty on \( x \) is \( \Delta x \). We define:

- **covariance**:
\[
\text{cov}_{x_{ij}} = (\Delta^2 x)_{ij} = \int dx_i dx_j \text{PDF}(x_i, x_j)(x_i - \langle x_i \rangle)(x_j - \langle x_j \rangle)
\]
\[
= \langle x_i x_j \rangle - \langle x_i \rangle \langle x_j \rangle
\]
(13)

- **relative covariance**:
\[
\text{rcov}_{x_{ij}} = \text{cov}_{x_{ij}} / \langle x_i \rangle \langle x_j \rangle
\]
(14)

- **uncertainty**:
\[
\text{unc}_{x_i} = \sqrt{\text{cov}_{x_{ii}}} = \Delta x_i
\]
(15)

- **relative uncertainty**:
\[
\text{runc}_{x_i} = \Delta x_i / \langle x_i \rangle
\]
(16)

- **correlation**:
\[
\text{corr}_{x_{ij}} = \text{cov}_{x_{ij}} / \text{unc}_{x_i} \text{unc}_{x_j} = \text{cov}_{x_{ij}} / \Delta x_i \Delta x_j = \text{rcov}_{x_{ij}} / \text{runc}_{x_i} \text{runc}_{x_j}
\]
(17)

Here, the covariance is a real, symmetric, positive \( N \times N \) matrix. A covariance may be sparse or dense or even (band) diagonal.

In the ENDF format, covariances and relative covariances are used exclusively in its covariance formats in MF=31–40. Unfortunately, most users want either only the (relative) uncertainty or the (relative) uncertainty and the correlation. Because of requirements 3.2 and 3.6, we really should attempt to accommodate both options within the same evaluation.

Returning to our example in Fig. 25, the self-covariances on the left side could be expressed several ways. They clearly could be expressed as either covariances or relative covariances. They could also be expressed as a combination of (relative) uncertainty and correlation matrix. In principal, both options must be allowed. This is illustrated in section II.B in Fig. 8.

In Figure 8, the original data are “data version 0” and consists of an energy–cross section interpolation table. The main container, containing “data version 0” links to “covariance 0,0” where the covariances for these data are stored. From the data and the covariance, we can compute the uncertainty on the cross section data and
generate a “data version 1” which now has an additional uncertainty column. “Data version 1” links back to “data version 0” and “covariance 0,0”.

**B. Covariance of continuous functions**

The covariance data we wish to store are sometimes discrete parameters. However, more often they are continuous functions of a variable such as incident neutron energy. To pack the covariance of a continuous function into a matrix, we must discretize. In the ENDF format, this is usually done by grouping in all the independent variables.

For example in ENDF, a cross section $\sigma(E)$’s covariance is given group wise as $\Delta^2\sigma_{ij}$. The group boundaries can be thought of as forming a basis function expansion

$$
\Delta^2\sigma(E_1, E_2) = \sum_{ij} B_i(E_1) \Delta^2\sigma_{ij} B_j(E_2) \quad (19)
$$

where the basis functions are window functions:

$$
B_i(E) = \begin{cases} 
0 & E < E_i \\
1/(E_{i+1} - E_i) & E_i \leq E \leq E_{i+1} \\
0 & E > E_{i+1} 
\end{cases} \quad (18)
$$

Thus, the basis function encodes the interpolation rule (in this case group-wise). To write the continuous covariance on the cross section $\Delta^2\sigma(E_1, E_2)$, we write

In ENDF, one other discretization method is used – the Legendre function expansion of angular distributions. In principal, many other discretization schemes are possible (various orthogonal function expansions, different order interpolation schemes, etc.)

To encode the discretization of a continuous function and to define the packing of covariance data as a function

FIG. 26 Coupling between reaction data and the corresponding covariance. This figure is a continuation of the example in Fig. 25 highlighting the links between a data element and the cross-covariances associated with it.
of independent variable, we require an element similar to the <axis> element discussed in section VII. This container can also encode the packing rule for a list of discrete variables, such as resonance parameters. In keeping with the latest version of GND (Mattoon [2012]), we call this container a <grid> element:

**Requirement 35: <grid> elements**

1. Shall have an attribute denoting the number of elements
2. For grouped independent variables, shall have the bin boundaries
3. For discrete variables, shall denote the packing order
4. For other discretization schemes, shall denote the packing order of for example coefficients in an orthogonal function expansion

**C. Covariance of multi-dimensional functions**

In the previous section, we described how to discretize the covariance for a function of one independent variable. How might we handle a function with more than one independent variable? The easiest solution is to discretize in each independent variable separately. Each independent variable then has an associated <grid> element. We need only define the order we loop over elements in the individual grids to pack the covariance matrix.

We illustrate this with the case of the covariance on the Prompt Fission Neutron Spectrum (PFNS). The PFNS is a PDF \( P(E'|E) \) with a covariance \( \Delta^2 \sigma(E_1',E_1,E_2,E_2) \). Let us discretize this:

\[
\Delta^2 \sigma(E'_1,E_1,E'_2,E_2) = \sum_{ijlm} B_i(E_1)B'_j(E'_1)\Delta^2 \sigma_{ijlm}B_l(E_2)B'_m(E'_2) \tag{20}
\]

If we decide to say loop over \( i \) first then over \( j \), we can effectively pack the \( i,j \) dependence into rows in the covariance matrix. Similarly, we can pack the \( l,m \) dependence into columns.

So, we simply need a general rule for looping over the elements in a list of <grid>s. Here we state the simplest rule we can conceive. For a list:

```plaintext
<grid0>
<grid1>
<grid2>
...
<gridN>
```

Do the following:

1. Loop though highest index grid bins first (<gridN>)
2. Then the second last
3. ...

**D. Covariance and correlation matrices**

We are now in a position where we can define the structures that can be used to store self- and cross-covariance information. Both self- and cross-covariance information can be stored in <covariance> elements, either as absolute or relative covariances. However, while a self-covariance is a full-fledged covariance in its own right, a cross-covariance is only a part of a larger covariance. So, we can convert a self-covariance into an uncertainty vector and a correlation matrix and the result will not introduce more couplings to other data. On the other hand, a cross-correlation does require linking the cross-correlation matrix with the two separate uncertainty vectors, introducing additionally (and needless) complexity.

To store just the covariances, we define a <covariance> element whose structure is shown in Figure 27 panel (a). In this element, we note the presence of <row> and <column> elements. These elements connect the covariance to the underlying data being described and contain required <grid> elements described in the previous two sections. Because self-covariance data connects only one data set to itself, it has the same row <grid>s and column. On the other hand, cross-covariance data must have different rows and columns that depend on the data getting correlated. The requirements of a covariance are:

**Requirement 37: <covariance> and <correlation>**

1. The format shall have a spot for optional documentation
2. The format shall have a <row> element which includes a link to the original underlying data and the <grid> elements to decide the data to
Fig. 27 The `<covariance>` and `<correlation>` elements. As they share identical components, but have very different meanings, the element name must most likely be used to distinguish them.

similarly, when one chooses to adopt the uncertainty and correlation matrix option, the requirements for a `<correlation>` are identical to a correlation. We give them separate tags to avoid confusion. We comment that there is no good reason why the uncertainty and correlation need to have `<grid>`s and `<axis>` elements that line up bin-by-bin. All that matters is that the domains of the independent variables are the same.

**Suggested Test:** Since covariances link to data and data links to covariance, we will need to check the hyperlinks for consistency.

Within the covariance element is the matrix data itself and the matrix data element is shown in Figure 28. Its requirements are

**Requirement 38:** `<matrixData>`

38.1 Flag to denote whether this covariance is absolute or relative
38.2 Flag to denote that a covariance matrix represents a quantity that has a fixed normalization (e.g., a probability distribution such as \( P(\mu, E) \) that must integrate to 1.0).
38.3 The data as a `<matrix>` or a `<matrixSandwich>` (see below)

**E. Weighted sums of covariance**

How can the format allow an evaluator to break up the covariance into components, say statistical errors from a
fit and systematic errors arising from experimental normalizations? Or say, how can an evaluator write the covariance for the \((n,\text{tot})\) cross section as the sum of the \((n,\text{el})\) and \((n,\gamma)\) cross sections? Both cases are easily addressed by adding a `<weightedSumOfCovariances>` element, shown in Fig. 29. This element takes advantage of the fact that the sum of two covariance matrices is still a covariance matrix. A `<weightedSumOfCovariances>` element should be usable anywhere where a `<covariance>` element can be placed. Incidentally, the ENDF format also allows for this construction. The requirements for this element are

**Requirement 39:** `<weightedSumOfCovariances>

39.1 Numerical weights (in ENDF, these are just floats of a component

39.2 Either the components as `<covariance>`s or links to `<covariance>`s

By allowing us to sum up covariance data in this manner, we open ourselves up to new classes of data bugs. Therefore we suggest a few tests of these data:

**SUGGESTED TEST:** Testing for dead links is especially important here as one must always traverse all links in order to reconstruct a covariance.

**SUGGESTED TEST:** All linked covariance data must be convertible to a `<covariance>` element so that they can be added.

---

**F. The “Sandwich Formula”**

Often times we have a parameter \(f_i\) that we want the covariance on, but it depends on something else, say \(\vec{x}\) and it would be much more efficient to store the covariance on \(\vec{x}\) directly. A case in point is the RRR parameters. The reconstructed cross section from tens of resonance parameters may have thousands of energy points to achieve a reasonable accuracy.

For function \(\bar{f}(\vec{x})\), the sens\(_{ij}\) \(= \partial f_i(\langle \vec{x} \rangle) / \partial x_j\) and is called the sensitivity matrix. Assuming that

\[
f_i(\vec{x}) \approx f_i(\langle \vec{x} \rangle) + \sum_j \frac{\partial f_i(\langle \vec{x} \rangle)}{\partial x_j} (x_j - \langle x_j \rangle)
\]

(21)

is a good approximation to the variation of \(\bar{f}(\vec{x})\) around
where $\langle \vec{x} \rangle$, we can evaluate the covariance of $f$ using the “sandwich formula” is

$$\text{cov}_{ij} = \sum_{i'j'} \text{sens}_{ii'} \text{cov}_{x_{i'j'}} \text{sens}_{j'j}$$

(22)

The “sandwich formula” can be reframed in terms of the relative covariance

$$\text{rco}_{ij} = \sum_{i'j'} \text{rsens}_{ii'} \text{rco}_{x_{i'j'}} \text{rsens}_{j'j}$$

(23)

where

$$\text{rsens}_{ij} = \frac{\partial f_i(\langle \vec{x} \rangle)}{\partial x_j} = \frac{\partial f_i(\langle \vec{x} \rangle)}{\partial (\ln x_j)}$$

(24)

The “sandwich formula” provides the scheme for deterministic uncertainty propagation.

Discussion point:
In many cases, the sensitivity of model parameters can be precomputed. In this case, we may not need to store the sensitivity matrix itself. Should we allow this? It makes for smaller files, but shifts the burden of computing the sensitivities to the processing codes.

Resolution:
Yes, this is already the case for RRR parameters.

As an aside, the covariance admits an eigendecomposition into $N$ eigenvalues $\lambda_i$ with eigenvectors $\vec{v}_i$. The covariance can be diagonalized in the eigenbasis as

$$\text{cov}_{x_{ij}} = (\vec{v}_k) i \lambda_k (\vec{v}_k)_j$$

(25)

This is the “sandwich formula” again, where the eigenvalues play the role of the sandwiched covariance matrix and the eigenvectors play the role of the sensitivity matrix. Often times the effective rank of a matrix $N_{eff}$ is much smaller than the actual rank $N$ because many of the eigenvalues are sufficiently close to zero that they may be neglected. The process of taking the main eigenvalues is called principal component analysis (PCA). Thus the “sandwich formula” storage scheme can be used to efficiently pack covariance matrices even in the absence of underlying parameter dependencies by using PCA. we will flesh this idea out further in the example below in section VI.H.

To support the “Sandwich Formula”, we must define the structure of a <matrixSandwich> and a <sensitivity>:

**Requirement 40: <matrixSandwich>**

40.1 The underlying parameter <covariance>
40.2 A <sensitivity> for the rows of the covariance

40.3 If the block of the matrix is off-diagonal, a <sensitivity> for the column as well.

**Requirement 41: <sensitivity>**

41.1 Optionally a <documentation> element
41.2 The <matrixData> for the sensitivity matrix
41.3 A <column> with a link pointing to the <column> element’s <axis> of the underlying parameter covariance matrix. This also defines the packing of the sensitivity matrix since we want them to match up for the matrix multiplication.
41.4 A <row> that mimics the row one would get if we were storing the full covariance on the derived data. Therefore we need an <axis> element to determine the packing of the sensitivity matrix and a (possibly fake) link to the <form> of the derived data.
41.5 An option for precomputed sensitivity matrices in the resolved resonance region (to store the MT32 covariances in ENDF).

**G. Monte Carlo Sampling**

How can one use a covariance to generate realizations for a Monte Carlo approach to uncertainty quantification? Suppose we have some $\vec{x}$ with a Normal PDF $P(\vec{x})$ specified by the mean $\langle \vec{x} \rangle$ and covariance $\text{cov}_{x_{ij}}$. To find the expectation value of a function $f(\vec{x})$, we do

$$\langle f \rangle = \int d\vec{x} P(\vec{x}) f(\vec{x}) \approx \frac{1}{N} \sum_{R} f(\vec{x}_R) P(\vec{x}_R)$$

(26)

Here the sum is over realizations of $\vec{x}$ drawn from the PDF. To generate the realizations $R$, we use principal component analysis (PCA) again:

$$\vec{x}_R = \langle \vec{x} \rangle + \sum_{i} \xi_{iR} \vec{v}_i \sqrt{\lambda_i}$$

(27)

Where $\xi_{iR}$ is drawn from a (log) normal distribution.

Discussion point:
Should we support ensembles of evaluations or evaluation parts (like TMC or list-mode output)? Would need index of realizations maybe. Could this be handled using the metaEvaluation scheme?
FIG. 30 The <sensitivityMatrix> element which connects the small matrix inside a “matrix sandwich” to the external covariance matrix in Eq. (22).

Resolution:
One would need reasonable number of samples $N_{samp}$ for each of the $i$ directions. This means that one needs $(N_{samp})^i$ samples to effectively sample all of $x$’s PDF to reliably propagate uncertainty. Saving these samples is then not really an effective savings of space. However, with a nativeData scheme, one should be able to accommodate variations.

H. Examples of covariance data usage in this hierarchy

In order to illustrate some of the concepts we have developed to describe our implementation of covariance data, we turn to three examples.

In our first example, we demonstrate the storing of Prompt Fission Neutron Spectrum (PFNS) covariance. The situation is shown in figure 31. In this figure, the PFNS itself is stored in the fission <reaction> under the neutron reaction product. The PFNS has as a subelement the covariance data accompanying it. The <covariance> itself is a self–covariance, is relative and the data are normalized. The element has documentation and <matrixData>. Since the covariance is on a PFNS, and the PFNS is a function of both the incoming and outgoing neutron energy, there are two <grid> elements associated with the <row> element.

In our second example, we show how to use the sensitivity matrix approach and the eigen-decomposition of a covariance to compress the covariance matrix. This is illustrated in figure 32. In this figure, we focus on a generic self–covariance $\text{cov}_{x_i}$. The <documentation> and <row> data (including a <grid> specification and link to the original data) are shown on the left side of the figure. We focus on the covariance itself in a <matrixData> container. We diagonalize our generic self–covariance using equation (25), producing a list of eigenvalues $\{\lambda_k\}$ and associated eigenvectors $\{v\}$. As we observed in the discussion around equation (25), equation (25) is already in the form of the “Sandwich formula”. The sensitivity matrix is just the matrix of eigenvectors, $\text{sens}_{ik} = (v_k)_i$, and the inner covariance is the matrix of eigenvalues $\text{cov}_{kk'} = \delta_{kk'}\lambda_k$. The matrix of eigenvalues goes in the inner <covariance> element in figure 32. The <row> (and <grid>) of the <covariance> is especially simple since the inner matrix is just a matrix of eigenvalues. The sensitivity matrix (the matrix of eigenvectors) goes in the <sensitivityMatrix> element and links to the inner and outer <row> element. We comment that the ENDF format also allows this approach to storing a matrix, using the MF=30 as described in subsection 30.3.3. To our knowledge, no one has used the MF=30 format in a production environment.

In our final example, we show how we re-implement the ENDF format’s MF=32 parameterized resonance covariance data. This is illustrated in figure 33. In this example, we will use the <sensitivityMatrix> scheme again, but this time it will be a little simpler since the sensitivity matrix is known analytically and is given in the SAMMY manual [Larson 2006]. In figure 33, we again start from the upper left side with the outermost <covariance> element. In this case, the covariance matrix we refer to is the full covariance of all of the reconstructed pointwise cross sections. Because of that, the <row> element of the outer <covariance> points to several cross section data sets and has two <grid> elements, one for the incident energy grid and another to allow looping over the reactions themselves. Inside the <matrixData> and <matrixSandwich>, we again find the <sensitivityMatrix> and the innermost <covariance> containing the matrix of resonance parameter covariance. The <row> element in the innermost <covariance> element describes how resonance parameters are packed into the <covariance> element.
VII. REQUIRED LOW-LEVEL CONTAINERS

At the low-level, data are described in universal terms defined by mathematics and found in all of science. Data may be categorized as constants, lines, surfaces or more complicated multi-dimensional functions. One potential way to consider the split between high-level and low-level data containers is to define it as the boundary between physical and mathematical descriptors of the data. Given that the low-level containers are crafted with this intention, they should be flexible enough to serve any higher-level purpose. Requirements for low-level data containers are:

**Requirement 42: General low-level**

42.1 Distinct regions of data with different meta-data shall be clearly delineated.

42.2 General-purpose data types shall be defined that are compatible with commonly used computer languages as well as their common usage and libraries.

42.3 Containers should make efficient use of computer resources. This is a compromise between:

42.3.1 efficient use of memory - volatile (e.g., RAM) and non-volatile storage (e.g., disk drive).

42.3.2 ease of conversion to other forms.

42.3.3 time to convert to other forms.

42.4 The structures should reduce redundancy.

42.5 Data containers should be designed to be consistent with object-oriented programming. This includes the nesting of data containers when it make sense instead of defining a new
FIG. 32 Illustrating the use of the `<sensitivityMatrix>` approach to compress a covariance matrix using the principal components of the covariance determined from an eigenvalue decomposition.

FIG. 33 Illustrating the storage of resonance covariance data. This is our approach to the ENDF format’s MF=32.

42.6 The specification for a data container shall state whether that container is extensible. If the container is extensible, the specification shall define the process for extending the container.

42.7 A mechanism for storing units and labels for data should be specified for each container type.

Even at the low-level there are still competing goals that require different basic structures. Each of these
needs to be distilled down until all that is left is the essence of its type. For the code developer, this will be similar to the process of code refactoring (Wikipedia 2015) where one seeks to remove redundant coding. For the database developer, this will be similar to the process of normalization (Wikipedia 2015a) where one seeks to untangle dependencies. In the end, one should be left with sets of one or more objects that each provide a specific functional form.

One of the most universal examples of this is to define a curve by use of a set of points. Alternately, this same curve may be defined by its parametric equivalent. Each fulfills the requirement to define a curve. In the end, what is to be done should shape how data are stored, but specifics of how it should be done should not be unnecessarily restricted.

The discussion in this section is a summary of the larger discussion of requirements and specifications in the “General–Purpose Data Containers for Science and Engineering” document (WPEC Subgroup 38 2015b).

A. The lowest-level

In a hierarchal structure, nodes most often contain another collection of nodes. However, at some point one reaches the bottom. In the case of a structure used to store scientific data, the lowest-level containers must store numeric data. That is, integer, fractional, real, and complex values that describe, for example, interaction probabilities, angular states, decay constants or barrier penetrability. To provide definitions for a structural lexicon, collections of nodes belong in “container nodes” that form branches of related information with data in “leaf nodes” at the end of the tree.

There are very pragmatic arguments that must be considered in how large collections of data are stored in memory. To demonstrate, consider the extreme where a desire to clearly label each data point might lead to a format in which every item was individually wrapped. This tends to be enormously wasteful in ways that do not gain any clarity. It is easily seen by considering wrapping every point of a set of values that define a piecewise curve. Wrapping the vectors that represent the values for each point clearly saves considerable space for any non-trivial data set. Conversely, if groups of these points have different interpolation schemes, delineating them by interpolation region greatly enhances clarity.

Similar consideration must be given to the issue of reading and writing data to memory. The ENDF format is a marvel of efficiency considering the 80 character line length inherited from its punch card origins. Each line knows its exact place in a file, its material, data type and reaction and stores up to six associated values. However, to maximize the significant figures available for its eleven character numeric values, ENDF allows one to drop the “e” from an exponential leaving, for example, “1.2345+7” in place of “1.2345e+07”. Given the strong Fortran roots of our community and Fortran’s understanding and acceptance of this format, it provides the broadest compromise. However, it comes at the expense that today’s modern languages – C, C++, Perl, Python, etc. – must parse these values ad-hoc at considerable expense in extra read times. Any modern format must store numeric values only in the widest adopted standards.

In practice all data – scalar, vector or multi-dimensional – can be serialized and stored in a contiguous block of memory, that is a vector of numbers. In binary form, these data tend to be the same in use, for example active RAM, or archived, for example on disk. An enormous time saving is available for such data if their movement is done by large block reads and writes that forego the need of conversions. Even in textual form, transformation to and from a sequence of numbers may be done much faster when it does not contain extraneous text. As the bulk of the data we intend to store is numeric, providing a “lowest-level” container that is strictly a vector of numbers can provide enormous time savings.

Unfortunately, nodes tend to be expensive to create in practice. This has implications for scalar data. While a low-level vector array container provides clarity and space savings for sets of numbers, it may not be necessary for a scalar quantity. In fact, there are times when it can be a significant waste of space. Consider two XML examples of storing the mass of a nuclei: “<mass array='True' size='1' id='U235'>235.043</mass>” versus “<mass id='U235' value='235.043' />”. The first example adds a minimum of 20 extra bytes for each scalar quantity and actually obscures the value with unnecessary attributes. For other storage formats, for example HDF5, this can be even worse. However, there is a clear need to recognize data, as opposed to meta-data, that are stored as an attribute. This argues for a clear universal attribute name to indicate the scalar data value.

Discussion point:
Should scalar values stored as attributes use a common naming convention? In other words, should we require the temperature always to have the attribute name temperature?

Resolution:
There is no general rule and we must (re)visit this issue on a case by case basis.

One of the desired outcomes for the use of this new structure is to encourage more documentation and to encourage placing documentation as close to the rel-
evant data as possible. Therefore, elements in the <documentation> structure (to be elaborated on later) are welcomed at many levels inside the structure. Documentation nodes may contain many forms of text input including XML, HTML, LaTeX, or other markup languages to help provide more meaningful information. As the community gains experience, we expect to eventually settle on a few (or hopefully one) standards.

**Requirement 43: Lowest level data elements**

43.1 Each node shall be either a container or a leaf

43.1.1 Container nodes shall contain only collections of other nodes

43.1.2 Leaf nodes shall contain only data or be empty

43.2 Data should be stored in the smallest possible container, but no smaller

43.2.1 Singular scalar values may be stored as attributes where space savings is significant

43.2.2 Scalar values stored as attributes shall share a universal attribute name

43.3 Leaf nodes storing only numeric data shall be clearly marked for ease of parsing

43.4 Numeric data stored as text fields shall use only generally accepted forms and should be avoided as much as possible

43.5 Text nodes may store documentation or other information in free or fixed formats

**B. General data containers**

At its simplest level, all numeric data can be broken into arrays of numbers and meta data describing the mathematical interpretation of the data. These should be grouped together into some kind of general data container. These general data containers serve the purpose of defining the mathematical and functional forms that are themselves basic components of higher level physical data storage containers. This addresses requirement 3.4.

**Requirement 44: List of general purpose elements**

44.1 Floats (float), complex numbers (complex), integers (int) and strings (string)

44.2 List or vector (<list>), must specify type of object in the list

44.3 Array (<array>), must specify dimensions. May be banded, symmetric, etc.

44.4 Table (<table>), like a 2-dimensional array, except it may contain non-numeric data and permits columns to have labels, units and data-type information

44.5 Orthogonal function expansion, Legendre polynomials and spherical harmonics being the most obvious

44.6 One-dimensional interpolation tables (<interp1d>): interpolation table for univariate data (i.e., \(x vs. f(x)\))

44.7 Two-dimensional interpolation tables (<interp2d>): interpolation table for bi-variate data (i.e., \((x, y) vs. f(x, y)\))

44.8 Three-dimensional interpolation tables (<interp3d>): interpolation table for trivariate data (i.e., \((x, y, z) vs. f(x, y, z)\))

44.9 Axis elements (<axis>) to specify units, labels, etc.

44.9.1 Specify names of \(x, y, z, \ldots\) axes

44.9.2 Specify units for all axis elements

44.10 Interpolation details

44.10.1 Specify interpolation scheme(s) or group boundaries

44.10.2 If interpolation refers to a probability density function (PDF), we also must specify whether is Normal or Log-Normal (Zerovnik, 2013).

44.11 Free text (<text>) (more on this in subsection VII.C)

44.12 Hyperlinks (<link> attributes) (more on this in subsection VII.D)

**Discussion point:**

Do we need to support complex numbers as well? It might help in the resolved resonances and with atomic reactions. It would also likely complicate specifying and using covariances on complex valued data.

**Resolution:**

Yes, some variants of resonance formats (both processed and evaluated) require complex numbers. However, these variants are very uncommon and we do not anticipate defining requirements for them at this time. Nevertheless, support for complex numbers should be considered.

**Discussion point:**

Should an axis define a normalization condition? For example, for \(P(E'|E)\) data are normalized over outgoing energy but not over incident energy...
Resolution:
General agreement that this should be denoted, preferably in the low-level container. Doing so avoids the requirement of adding additional context to denote normalization constraints.

Discussion point:
It has been suggested by several members of the nuclear data community to include uncertainty directly into elements such as the `<interpd>` table. This would make plotting the uncertainty simpler at the expense of introducing an additional data synchronization problem between the mean values and the covariance data.

Resolution:
This idea is adopted since it moves the data and its uncertainty/covariance together. If the covariance or correlation accompanying the uncertainty is not stored with the uncertainty data, a link to the covariance or correlation is needed. See subsection II.B.

C. Text

There is a common need for unstructured and partially structured text in evaluation and other documentation. There may be other, less obvious areas. In the documentation, one must store author names and affiliations and do so in non-Latin alphabets. Therefore, we require full UTF8 support (unicode) so at very least author names can be written in own language. This also enables each data project to write their own documentation in their own native language. Since encoding needs to be stated separately from the text itself, a `<text>` element must never be stored in an XML attribute field. Finally, we note that since text may be formatted, the internal `<text>` format must be denoted. Given that the reference format for which we are detailing requirements in will be a subset of the XML markup, we strongly recommend against allowing raw XML for data formatting so as to avoid parsing problems.

D. Hyperlinks

Links (links are attributes in XML) are an important part of the new format(s) and allow the evaluator to refer to other elements within the file or even to elements in external files or databases. Examples of data which use links include:

- Distributions for one reaction product may be treated as the recoil from another product, requiring a link to the other product.
- Production cross sections may be listed as an energy-dependent multiple of another cross section, requiring a link to the other cross section.
- Covariances are occasionally stored in a separate file from the quantities they correlate. Links are necessary to associate the covariance with the correct data.

Because the data are stored hierarchically, the path within a document can be followed straightforwardly. It is useful to think of these paths as similar to paths in a Unix filesystem, but with the top level of a document referred to with a URL.

Requirement 46: Hyperlinks

46.1 The paths may be absolute so that they can refer to external documents or relative so that they can make in-document referrals.
46.2 The URLs of the schema (in the case of an XML version of the format). We comment that these may not be accessible by some computers “behind the fence” so these URLs may be viewed as “placeholders” that can be overridden in specific applications.

One can easily imagine that one is using a nuclear data library on a computer not directly connected to the internet so external links may not be available. In that case, it would be up to the user of the data to remap the URL’s to the actual location of the data files on their own computer system. Anything we can do in our API specifications to simplify this task would be appreciated by the user.

Requirement 47: Placeholder hyperlinks

47.1 There should be support for place-holder names in urls. For example, `<myElement href="$(NUCLEAR_DATA_PATH)/fluxes"/>` where the NUCLEAR_DATA_PATH is defined by a particular institution.

**SUGGESTED TEST:** Link checking for all xpaths and hrefs.
VIII. SPECIAL REACTION CASE: ATOMIC SCATTERING DATA

The atomic scattering and relaxation formats in ENDF provide places to store collision data for electrons and photons as well as the atomic de-excitation data. Although the ENDF format is not meant to be a comprehensive resource for atomic reaction data, it has found a niche for storing data needed for electron and photon transport in otherwise neutron (or other nuclear) dominated transport applications.

If the readers of this document are interested in these data, they should consult the Virtual Atomic and Molecular Data Center (VAMDC) at http://portal.vamdc.org/. The VAMDC is a consortium of atomic and molecular data centers covering a variety of data types including experimental, evaluated, bibliographic, structure and reaction data. The entire code system is detailed in http://www.vamdc.org/standards. The VAMDC data are stored in the VAMDC-XSAMS format (an XML format). The format is described at http://www.vamdc.org/documents/standards/dataModel/vamdcxsams/index.html. There is also a wiki describing the entire system (web infrastructure as well as format) at http://voparis-twili.obspm.fr/twiki/bin/view/VAMDC/PortalHelp.

Atomic scattering data in ENDF includes only electromagnetic (electrons and gammas) projectiles interacting with the electronic orbitals of an atom. These data are very similar to nuclear reaction data, but simpler in some ways and more complex in others. The ENDF atomic scattering data are given in reactions specified by MT=500-599 in the formats specified by MF=23, 26, 27, 28. In the ENDF/B-VII.1 library (Chadwick, 2011), atomic data are collected in three sub-libraries:

- **atomic relax relaxation** (NSUB=6): Details the de-excitation of excited atoms or ions following an excitation or ionization event. The excitation process either excites an electron to a higher shell or knocks it out completely, resulting in a vacancy. This library tabulates the cascade of events that can occur as the atom or ion relaxes to its ground state. These data are handled by the particle properties database(s) discussed in section III and detailed in Ref. (WPEC Subgroup 38, 2015a).

- **electro-atomic scattering** (NSUB=113) Provides cross section and particle production data for electrons scattering off neutral atoms. Reactions tabulated include elastic scattering, bremsstrahlung and inelastic scattering resulting in either atomic excitation or ionization. Because the recoils of the target atom is small (and it would be difficult to compute the response of the electron cloud and nucleus in any event), it is neglected. Therefore, the data are considered to be given always in the lab frame and all scatterings transfer zero energy to the residual atom.

- **photo-atomic scattering** (NSUB=3) Provides cross section and particle production data for photons scattering off neutral atoms. Reactions include elastic scattering and inelastic scattering resulting in either atomic excitation or ionization.

These data are given in a standard <reaction> element with the following additional requirements:
Requirement 48: Atomic reaction data

48.1 A standard <reaction> element whose outgoing particles are photons, electrons and/or a residual atom

48.2 A particle property database for atomic and ionic de-excitation (i.e., decay) data (WPEC Subgroup 38, 2015a)

48.3 Outgoing photons may optionally use form factors for coherent and incoherent photon scattering (see MF=27) in a <dCrossSection dOmega> element. This is detailed below.

48.4 Usual outgoing distributions, with

48.4.1 Electron and gamma multiplicity (yields)

48.4.2 Outgoing electrons or photons may use form equivalent to LAW=1 (continuum, used for bremsstrahlung and ionization) (same as MF=6, LAW=1), or

48.4.3 Outgoing electrons or photons may use form equivalent to LAW=2 (two-body elastic) (same as MF=6, LAW=2), or

48.4.4 Outgoing electrons or photons may use form equivalent to LAW=8 (energy transfer for excitation, used for excitation and bremsstrahlung), described in MF=26; if so use <interp1d> to tabulate the energy transfer $E_T(E)$ for LAW=8

48.5 The residual atom product element with a location for the fluorescence yield. This is typically a float with units eV/photonization

48.6 An optional documentation element

48.7 Any links to covariance (if applicable)

The ENDF system for neutron and photon production data allows two alternatives for storing angular distribution data. One is by probability per unit $\cos(\theta)$ vs. $\cos(\theta)$, and the other is by Legendre coefficients. Neither of these is a “natural” method for photons. The natural method would be atomic form factors or incoherent scattering functions. These are discussed briefly below.

A. Incoherent Photon Scattering

The cross section for incoherent scattering is given by:

$$\frac{d\sigma_{coh}(E)}{d\mu dE'} = S(q, Z) \frac{d\sigma_{KN}(E)}{d\mu dE'},$$  \hspace{1cm} (28)

where:

$d\sigma_{KN}/d\mu dE'$: the Klein-Nishina cross section (Klein, 1929) which can be written in a closed form.

$S(q, Z)$: the incoherent scattering function. At high momentum transfer ($q$), $S$ approaches $Z$. In the other limit, $S(0, Z) = 0$.

$q$: the momentum of the recoil electron (ENDF specifies this in Å$^{-1}$).

$$q = \alpha \left[ 1 + \left( \frac{\alpha'}{\alpha} \right)^2 - 2\mu \left( \frac{\alpha'}{\alpha} \right) \right]^{1/2}$$ \hspace{1cm} (29)

$\alpha$: $E'/m_0c^2$,

$E'_s$: the scattered photon energy,

$\mu$: $\cos(\theta)$.

The angular distribution can then easily be calculated. Values of $S(q, Z)$ are tabulated as a function of $q$. The user presumably will have subroutines available for calculating $q$ for energies and angles of interest and for calculating Klein-Nishina cross sections. The user will then generate the cross sections for the appropriate cases by calculating $q$’s, looking up the appropriate values of $S$, and substituting them in the above formula.

B. Coherent Photon Scattering

The coherent scattering cross section is given by:

$$\frac{d\sigma_{coh}(E)}{d\mu dE'} = \pi r_0^2 \left( 1 + \mu^2 \right) \times \left\{ \left[ F(q, Z) + F'(E) \right]^2 + F''(E)^2 \right\},$$  \hspace{1cm} (30)

where:
$q$: $\alpha [2(1 - \mu)]^{1/2}$, the recoil momentum of the atom (ENDF specifies this in $A^{-1}$).

$r_0$: $e^2/m_0c^2$, the classical radius of the electron.

$F'(E)$: the real part of the anomalous scattering factor.

$F''(E)$: the imaginary part of the anomalous scattering factor.

The quantity $F(q, Z)$ is a form factor, which can be easily tabulated. At high momentum transfer ($q$), $F$ approaches zero. In the other limit $F(0, Z)$ tends to $Z$. The anomalous scattering factors are assumed to be isotropic. In addition, they smoothly approach zero at 1.0 MeV and can be assumed to be zero at higher energies.

An alternative way of presenting the photon scattering data would be to tabulate incoherent scattering functions and form factors. Users could then provide processing codes to generate the cross sections from this information. The calculation is quite straightforward and allows the user to generate all this scattering data from a relatively small table of numbers. The incoherent and coherent scattering data should always be presented as scattering functions and form factors, respectively, whether or not data are included.

By their nature, spallation reactions are inclusive in that many possible reactions are summed together and not individually considered. Although spallation is a complicated sum of reactions and processes, we can easily accommodate it in the hierarchy described above. To use it, one simply declares the spallation data to have their own `<reaction>`. The total of all the cross sections of all processes comprising the spallation reaction are placed in the `<crossSection>` element and each tabulated particle is given its own `<product>` element with variable multiplicities the full energy-angle PDFs. As the hierarchy can already accommodate these data, additional requirements are minimal:

**Requirement 50: Spallation**

50.1 Should allow a reaction annotation or alias for “spallation”

**X. SPECIAL REACTION CASE: RADIATIVE CAPTURE**

Here we detail needs for radiative capture data, including $(n, \gamma)$, $(p, \gamma)$, or any other reaction in which the only reaction product is a residual nucleus and emitted gammas. These reactions may be treated using R-matrix theory and the data stored in resonance formats or they may be treated using other reaction models stored point-wise. If radiative capture is treated with the R-matrix formalism, often one works in the Reich-Moore approximation, but one could treat with full R-matrix using format consistent with requirements in this document.

Radiative capture data can be handled just like any other channel that produces gammas, with one exception: primary gammas. When a projectile is captured, a compound nucleus is formed. This compound nucleus decays to discrete level(s) and a gamma cascade ensues. As particle marches up in incident energy, the first gamma(s) in the cascade march up in energy as well, in accordance with this equation:

$$E'_\gamma = E - \frac{M_{\text{targ}} + m_{\text{proj}}}{M_{\text{cn}}} + E'_{\gamma 0}$$

(31)

where $M_{\text{targ}}$, $m_{\text{proj}}$ and $M_{\text{cn}}$ are respectively the masses of the target, projectile and compound nucleus and $E'_{\gamma 0}$ is the energy of the gamma for a projectile with zero energy. The primary gammas and accompanying cascade gammas are unique to an isotope and are a potential tool for isotope identification through active interrogation.
XI. SPECIAL REACTION CASE: FISSION

In many ways fission is just a regular channel, but physically it is a continuum of channels all lumped together for practicality. Thus, while it fits neatly in our top level hierarchy, at the lowest levels (the components and forms), there are many data types we would like to include. That said, fission can be treated straightforwardly as a regular <reaction>. In this section we will elaborate on how to do this and reiterate requirements that must be met to support fission data.

Discussion point:
The top level fission reaction should just be called “fission,” the name should not be overly complicated.

A. Introduction

At the top level, the fission cross section includes not only total fission but multi-chance fission when zero or more neutrons are emitted prior to fission, based on the relative branchings for neutron emission and fission. This process is distinct from pre-equilibrium neutron emission. While the total fission cross section includes up to fourth-chance fission, no further information is available regarding reaction products from the different chances. Therefore, the most reasonable thing to do would be to have one <crossSection> element for each component, each named in such a way as to make clear what are the different components of the total fission cross section (e.g., 'first_chance'). The various chances are required to sum to the total fission cross section.

Discussion point:
In ENDF, we noted that all n-th chance cross sections all had the same Q value as the total fission Q value. If this value isn’t used it should be deprecated.

The reaction products from fission can be separated by timing into prompt, those emitted immediately after scission while the hot, excited fragments are cooling by neutron and gamma emission, and delayed, those that are emitted after processes occurring on a slower time scale such as beta decay. The prompt products include only neutrons, gammas and fission products. (The fission product yields are discussed in the following section.) The delayed products include not only neutrons and gamma but also electrons and neutrinos arising from the subsequent decays of the initial fission products following prompt emission. The delayed neutrons are further separated into time groups according to ranges of decay half-lives.

Currently ENDF includes an energy-dependent average energy release for fission which includes energy deposition from all the prompt and delayed emission products. The fission products are represented by post prompt neutron emission fission fragment kinetic energies. The components of the generalized energy dependent Q values include prompt neutrons and gammas as well as delayed neutrons, gammas, electrons and neutrinos. These are currently stored in a single ENDF file and, in GND, the coefficients describing the energy dependence are tabulated according to, for example, promptNeutrons, delayedGammas in the reactionProducts element. It may be advantageous to replicate the prompt neutron, prompt gamma, and delayed neutron contributions to the energy release within each of the product elements as long as the values match those at the top level. Indeed, the energy release here should be equal to the pseudo-Q value for fission (MT=18 in ENDF).

Discussion point:
The right place for energy deposited is in the appropriate reaction product. Should there be an overall element called energyDeposition (or similar)?

Discussion point:
The Q value is constant while the energy release is now energy dependent (MT=458 in ENDF) so this value should be deprecated, as mentioned previously for n-th chance fission.

Both prompt and delayed neutrons are included under the neutron product element. Current information in ENDF includes the average neutron multiplicity, ν, for prompt and delayed emission; the prompt fission neutron spectrum, PFNS; and the delayed neutron spectrum, DFNS. The delayed component includes DFNS values for different incident neutron energies separated according to time constant. Only prompt gamma emission is included under the gamma product element. Usually only prompt gamma energy deposition is included. However, there is room for the gamma multiplicity and energy spectra that could be made more use of.

The PFNS is unit normalized so that multiplying the PFNS at one energy by the ν for the same energy gives the correct energy distribution. The average PFNS is typically generated by the Madland-Nix Los Alamos model and tabulated. These prompt neutron data should be augmented by the neutron multiplicity distribution, \( P(\nu|E) \). Recent model calculations have shown that the shape of the PFNS depends on neutron multiplicity. For some applications, \( P(E'|E,\nu) \) distributions may be useful but sampling from such distributions in lieu of an inline model of complete prompt fission events would be superior if possible at a later time. Meanwhile, the combination of \( P(\nu|E) \) with the PFNS at a given \( E \) should be
sufficient for most applications, particularly those involving large systems and average quantities. Covariances on the PFNS are currently allowed in ENDF. These are relevant between incident energies, outgoing energies, and between incident and outgoing energies.

Discussion point:
Right now spontaneous fission is stored under radioactive decay. What about making a spontaneous fission reaction? Yes, I know a lot of things are different – no energy dependence etc. but it is a close cousin to neutron-induced fission, especially with respect to fission product yields and if they’re moving over under a reaction called fission, why not this also? In addition, the fission reaction structure should be applicable to any other projectile inducing fission, including photofission.

Resolution:
Spontaneous fission is a decay mode of many actinides and is covered by the particle properties database discussed in section III and detailed in Ref. (WPEC Subgroup 38, 2015a). Spontaneous fission data would go in a <decayProducts> element.

C. Fission format requirements

52.1 Allow total fission cross section to be broken out by chance in the crossSection element with “1st_chance_fission”, “2nd_chance_fission”, etc. elements. Ensure sum rules are obeyed so that the total fission cross section is retained.

52.2 Allow Fission Product Yield data (see next section for a discussion).

52.3 Allow a <table> of fission energy release data in the reactionProducts element. If broken out according to ejectile, the values should match those in the table. Alternatively, allow ejectiles such as neutrinos and (anti-)electrons to be plain reaction products, but with a relaxed completeness requirement that only requires there to be energy release through that product.

52.4 Allow for emission of fission fragments (products), neutrons and gammas in separate product elements. Electrons and neutrinos could be allowed for if users specify need but so far the tabulated deposition in the fission energy release should suffice.

52.5 Allow energy-dependent prompt, delayed and total ν in the <multiplicity> element. Ensure sum rules obeyed for each energy.

52.6 Allow P(ν|E) for prompt neutrons.

52.7 Allow PFNS using tables or Madland-Nix model.

52.8 Break out delayed data by time group and put each group’s delayed ν with the groups DFNS and time constant.

52.9 Allow all emitted prompt neutrons and gammas to have variable multiplicities and energy-angle spectra.

52.10 Allow for energy-energy covariances in PFNS. There should be the possibility for incident energy, E − E; outgoing energy, E′ − E′; and incident-outgoing energy, E − E′, covariances.

52.11 Allow for fission reactions with other projectiles.

Discussion point:
Should we allow P(E′|E, ν) data for prompt neutrons? If we really want to do this, then we should only do 1 or 2 multiplicities near the average (at a given energy) since otherwise it becomes too unwieldy to try and generate decent spectra for outlying values of ν.

B. Existing ENDF format

The total fission cross section is in MF 3, MT=18. The various chances are contained in other MT sections: first chance, (n,f) in MT=19; second chance, (n,n′f), in MT=20; third chance, (n,2nf), in MT=21; and fourth chance, (n,3nf), in MT=38. The data in MTs 19, 20, 21 and 38 must sum to the data in MT=18. All neutron and gamma fission data are stored under MT=18.

The total number of neutrons per fission, ν, is stored in MT=452. The total multiplicity is further broken down into delayed, MT=455, and prompt, MT=456, neutron multiplicities. The ENDF format stores energy-dependent Q values in MT=458, implemented in ENDF/B-VII.1. These are all stored in MF=1. The energy release per fission in MT=458, aside from the neutrino energy, should equal that stored in MT=18-21, 38. Covariances of the average neutron multiplicity are found in MF 31.

The PFNS is stored with the cross section in MT=18. The temperature parameter, T_M, is found in MF 5 and the spectra are generated with the appropriate T_M values according to the Madland-Nix spectrum, LF=12.

Gamma yields are in MF=12 but fission gamma data are sparse with few energy points.

The way ENDF handles the fission product yields is described in the next section.
**Discussion point:**
On the subject of the 6 delayed time groups data, should we even attempt to connect the delayed data to the particle properties data and the Fission Product Yields? The 6 time groups are really effective time groups. The real process involves hundreds of individual beta decays, at least one for each independent fission product. These are averaged over in some fashion to produce the time groups.

**XII. SPECIAL COMPONENT CASE: FISSION PRODUCT YIELDS**

Fission Product Yields (FPY) are currently stored in their own sub library in the major evaluated data libraries (e.g., ENDF/B-VII.1), but conceptually they really belong in the description of emitted particles from the fission reaction. Because there are many different ways to induce fission, FPYs rightfully belong in a discussion of mid-level data structures.

**A. Introduction**

In the 2012 Working Party on Evaluation Cooperation (WPEC) meeting, two new subgroups were created: SG-37 to investigate Fission Product Yields (FPYs) and SG-38 to define a possible replacement for the ENDF nuclear data format. The Generalized Nuclear Data (GND) format is the main candidate for replacing the ENDF format and is under active development under auspices of WPEC/SG-38, lead by D. McNabb. GND is an outgrowth of earlier LLNL (US) project to replace LLNL’s own internal ENDF format and the initial focus of the GND project was to develop formats and tools for handling neutron and charged particle transport data. SG-38 is now looking toward other ENDF formats and data, in particular, fission product yield (FPY) formats.

In the May 2013 SG-37 meeting, many new theoretical and experimental results were presented and new evaluations and evaluation techniques were presented. The new evaluations provide extensive covariance data which cannot be accommodated in the ENDF format. However, users require these covariance data for performing uncertainty quantification in many applications. The concurrent development of the GND format allows us to address many shortcomings of the ENDF format and define a new format that can meet future needs of members of the SG-37 group.

Let us now discuss what data SG-37 intends to store in GND. The Independent Fission Product Yields (IFPY) are the fragments immediately after fission and de-excitation from prompt neutron and gamma emission while the Cumulative Fission Product Yields (CFPY) are the fragments after they undergo further rapid (beta and other) decays. The two yields are connected by the Q-matrix:

$$CFPY_i(E) = \sum_{ij} Q_{ij} IFPY_j(E)$$  \hspace{1cm} (32)

This implies that, in practice, only IFPY or CFPY along with the Q-matrix may be needed. There are likely situations where both are needed, especially when fitting multiple kinds of experimental data. That said, it is typically easier to measure the CFPY while applications typically prefer IFPY.

The Q-matrix is a sparse matrix derivable from knowledge of the fission fragment decays and. A. Sonzogni and R. Mills have codes that can compute the Q-matrix from an ENDF-formatted decay sublibrary. Although the Q-matrix is a derived quantity, it is derived from data potentially not associated with the FPYs tabulated (e.g., JEFF yields could in principal use ENDF/B decay data) so should be associated with the IFPY and CFPY. Also, implicit in the Q-matrix is a time-integration that sets the maximum decay time of parent nuclei of interest to the evaluator or application.

During the SG-37 meeting, deuteron-, alpha-, photonuclear- and other particle induced yields in addition to the traditional neutron- and spontaneous yields were reported. The ENDF format has provisions for all of these.

In the process of evaluating yields, one often derives covariance data relating the yield of an isotope/isomer as a function of incident energy and covariance data relating yields from different isotopes/isomers. In addition, as the Q-matrix is derived from decay data which also has uncertainties on branching ratios, the Q-matrix may also have covariance data. The branching ratios enter into the Q-matrix linearly so the covariance calculation is straightforward. The uncertainties on half-lives is typically not so important except in the few cases of a long lived product whose half-life exceeds the integration time used to compute the Q-matrix. In this case, uncertainty propagation is very non-trivial since the half-life dependence is strongly nonlinear.

**B. Existing ENDF format**

The ENDF format make provisions for storing the IFPY in MT=454 and CFPY in MT=459. Both FPYs use the same ENDF format and this format stores tables of (I, YI, dYI), with I denoting the isotope/isomer in question, YI the corresponding yield and dYI the uncertainty on the yield. The yields are given for several incident energies E with a rule for interpolating from one energy to the next.

In practice, the interpolation rule is poorly enforced. For neutron induced fission yields, four energies are typ-
ically given which correspond to group boundaries for “thermal”, “fission spectrum”, and “14 MeV” neutrons. In practice, the yields change slowly with incident energy so this has proven to be a problem only in a few applications.

The ENDF format does not provide a way to store fission yield covariances nor does it provide a way to store the Q-matrix.

C. Detailed FPY format requirements for GND

During the WPEC/SG-37 meeting, D. Brown presented some ideas on possible formats and began a dialog with members of WPEC/SG-37. As a result of subsequent conversations, D. Brown developed a list of requirements for a new FPY format. We expect this list to evolve somewhat as discussions continue.

Requirement 53: Fission Product Yields (FPYs)

53.1 Clear rules for interpolation rather than a few vaguely defined groups (e.g., “thermal”, “fission spectrum”, “14 MeV”). Do not implicitly include spectrum averages in values.

Discussion point:
ENDF’s energies really are group averages. Should the fact that they are group averages be advertised? Should we also put in the group flux somehow?

53.2 Clearly defined range of validity of evaluation that can be matched to other reaction data. This may be nothing else than the \(<\text{evaluation}>\)-wide \(E_{\text{low}}\) and \(E_{\text{high}}\).

53.3 Clear location in the GND reaction hierarchy

53.4 Any incident particle (or none)

53.5 Per isotope/isomer yield \(Y_i(E)\), identical format for IFPY and CFPY

53.6 Per isotope/isomer yield uncertainty \((dY_i(E))\), identical format for IFPY and CFPY

53.7 A spot in the Q matrix markup to denote a cut-off halflife. Decay modes with longer halflives were ignored in the construction of the Q matrix.

53.8 Facility to store per isotope/isomer covariance on yield \((\Delta^2Y_i(E,E'))\), identical format for IFPY and CFPY.

53.9 Facility to store cross-isotope/isomer covariance \((\Delta^2Y_{ii'}(E_i,E_{i'},E_{ii'}))\), identical format for IFPY and CFPY. Only IFPY’s may be correlated with IFPY’s and CFPY’s with CFPY’s, the Q-matrix couples the IFPY and
53.10 Facility to optionally store the Q-matrix which connect the IFPY and CFPY. The upper cut-off integration time used in the generation of the Q-matrix must also be stored.

53.11 Facility to denote which (if any) of IFPY and CFPY is a derived quantity.

Discussion point:
Q-matrix can be computed from the decay library. Is Q-matrix something we want to store? It can be a very stringent requirement but if we computed CFPY using the Q-matrices computed from the decay data of the same library, we could store only IFPY data (and related uncertainties and correlations). In this sense CFPY can be considered as a sort of “reconstructed” FPY data as well as cross sections in the resolved resonance region are reconstructed from the resonance parameters. Obviously, this procedure would rely on a complete and consistent decay library and related uncertainties.

Resolution:
We want to allow storing Q-matrix as an option, not a requirement. Similarly, we were not requiring the evaluator to provide both the CFPY and the IFPY. However, we did want the evaluator to have the option to store either the CFPY or the IFPY and then the Q-matrix. Then the user can reconstruct what they need for their application. In the event that the evaluator has some fancy pants Bayesian scheme :) that requires a simultaneous fit of some IFPY and some other CFPY, then that evaluator would have to store everything for the sake of internal consistency.

Discussion points:
Additionally, we would like to investigate the possibility of storing the covariance of the Q-matrix.

Discussion point:
It would be very useful to integrate a Q-matrix calculator into one or more processing/testing codes so that the IFPY and CFPY and the decay data can be brought in accord with one another.

Discussion point:
On this option, we have to keep in mind that, in general, there are files with about 1000 FPY data for about 4 incident (neutron) energies. I would prefer option of Fig. 35. To imagine thousands of elements. 


d. Discussion of possible implementations

During the WPEC/SG-37 meeting, one “strawman” format was proposed, and in discussion with C. Mattoon and B. Beck others were discussed. Here we summarize this discussion and provide pros and cons. We expect that the format will go through many iterations as we attempt to meet the above requirements while maintaining a coherent and (hopefully easy to understand) structured data format.

Figure 34 shows an example of where fission product yields could fit in the current GND reaction hierarchy. As fission products describe the emitted particles of a fission event, it is logical to place them in the fission <reaction>s <outgoingChannel> of the corresponding <reactionSuite>. The collection of all fission product yield data is assembled in a <fissionProductYields> section. The FPY section has an optional nativeData attribute that specifies which of the IFPY and CFPY is the original source distribution. As spontaneous fission is a decay mode, <fissionProductYields> sections can appear in a particle properties database as spontaneous fission decay products.

Within the <fissionProductYields> section, we imagine an <independentFissionProductYields> section for IFPY, a <cumulativeFissionProductYields> section for CFPY and possibly a <fissionYieldConversionMatrix> section to store the Q-matrix. We expect the markup for IFPY and CFPY be identical, as in the ENDF format. Figures 35 and 36 show two different possible arrangements for data in the IFPY and CFPY sections.

Figure 35 shows one option. Here the yield tables use a modified version of the GND <linear> markup. The <linear> markup is attractive for several reasons:

- The interpolation rule specification is well developed.
- Fudge, the main tool for manipulating GND data, has strong data structures for storing X-Y data, including linearization, plotting, etc.
- All data for one nuclide are collected together in a simple, readable way.

The GND’s <linear> markup is a general markup used for data consisting of X-Y pairs. In our case, we would like to add dY’s as well. The current <linear> markup also allows for only one <data> tag whereas we imagine one per nuclide.
FIG. 34 A sample GND `<reactionSuite>` demonstrating where the fission product yields could reside within a fission `<reaction>` section in the current GND format.

in a horizontal array as described in the option of Fig. 36 is a little bit impractical.

**Resolution:**
A 1000 x 4 table may be silly and unworkable. However that arrangement is the most ENDF-like, so we put it in as an option.

Figure 36 show another option for storing FPY. Here all data are stored in the GND `<table>` markup. This markup is quite general and compact. It can accommodate any number of isotopes simply by adding another column (or pair of columns if dY is included). We would need to add a provision for specifying an interpolation rule in energy as this is not already provided by the current `<table>` markup. With this, we would need to add quite a bit of coding to Fudge in order to generate plots and manipulate the yield data.

**Discussion point:**
About the format for FPY covariance data, it was thought that ENDF compact format developed and used to store large covariance matrices would be suitable for this problem. However, there is no such option proposed in this requirements document.

**Resolution:**
In GND and the new format there is agreement that there will only be one covariance matrix format and it will be clearer than what is in ENDF. For each dataset that has covariance data, there will be a link (with a URL) to its own covariance and any (and all) cross covariances with other datasets. It is hoped that this arrangement can be made practical for FPY’s so we don’t have 1000 mini-FPY tables, each with 1000 URL’s pointing to 1000 mini-covariance matrices.

The Q-matrix should be stored in its own section, here called `<fissionYieldConversionMatrix>`. GND already provides a `<matrix>` markup and it is natural to store the Q-matrix itself here. However we need to know how each row/column maps to a yield table. To solve this, in this example we provide the URL to the data for each row/column in the IFPY and CFPY tables. It is unclear at this time if this is the optimal way of referencing column and row elements and it depends on the way FPYs are stored in their corresponding data sections.

**XIII. SPECIAL REACTION CASE: LARGE ANGLE COULOMB SCATTERING (LACS)**

As we outlined in subsection IV.B, charged particles do not have a finite total cross section or angle integrated elastic cross section. Quantum mechanically, charged particle elastic scattering is a sum of Coulomb and Nuclear amplitudes:

\[ A = A_{\text{Coulomb}} + A_{\text{nuclear}} \]  

(33)
The Coulomb piece is analytic and well known. The nuclear piece must be evaluated. The cross section for elastic scattering is of course the square of the amplitude so the differential cross section has three terms:

$$\frac{d\sigma_{el}(E)}{d\Omega dE'} = \sigma_{\text{Coulomb}}(E) + \sigma_{\text{int}}(E) + \sigma_{\text{nucl}}(E)$$

(34)

The last two terms in this equation are traditionally lumped together in a “nuclear+interference” term. Note, this scheme is used for elastic scattering of electrons in the electro-atomic sublibrary in ENDF, but with no nuclear amplitude.

Whether the target and the projectile are identical or not, the Coulomb term is very singular:

$$\frac{d\sigma_{\text{Coulomb}}(E)}{d\Omega dE'} \propto \frac{\eta^2}{k^2(1 - \mu)^2}$$

(35)

Therefore, the elastic cross section diverges at small incident $E$ and small angles ($\mu \to 1$). One might think that, since this is analytic, we don’t have to store it and there is no problem. The problem is that since the Coulomb amplitude carries the square-root of these divergences, the interference term $\sigma_{\text{int}}$ in the total elastic differential cross section also carries divergencies.

The traditional workaround is twofold:

- Start the “nuclear+interference” data tables at some finite incident energy where nuclear effects become noticeable. This eliminates the incident energy divergence in the tabulated data.
- Cut-off the “nuclear+interference” term at small angles. At small angles, Coulomb scattering dominates and must be handled in particle transport separately with techniques such as condensed history. ENDF data uses $10^\circ$ as a cut-off (if remembered correctly, can’t be found in documentation so far).

Discussion point:
ENDF puts this data in MF=3 and MF=6, LAW=5. This leads to confusion since what is in MF=3 is not a partial cross section, but rather a kludge to get around the divergence. Indeed, the presence of these data in ENDF tempts one to try to heat it much like one does for neutron incident data.

Resolution:
We recommend putting these data in a special LACS $<\text{dcrossSection}\_d\Omega\text{mega}>$. 

Requirement 54: LACS

54.1 A $<\text{dcrossSection}\_d\Omega\text{mega}>$ for LACS data
54.2 A $<\text{form}>$ for “nuclear+interference” data
54.3 A location to denote the cut-off angle (since it may not be ENDF's default $10^\circ$)

XIV. SPECIAL REACTION CASE: THERMAL SCATTERING LAW

Thermal neutron scattering law (TSL) data describe the situation where the de Broglie wavelength of an incident neutron is so large that the neutron wave function cannot resolve individual nuclei but rather 'sees' the macroscopic material. The incident neutron cannot be absorbed by the material and may only (in)elastically scatter off of it. Thermal neutron scattering is typically formulated using the theory of Van Hove (VanHove 1954) which we detail here following the treatments in Refs. (Behr 2010) and (Cacuci 2010).

TSL data are given in sub-library 12 (NSUB = 12) in the ENDF6 format (using MF=7, MT=2 and MF=7, MT=4 data structures). This sublibrary provides dimensionless scattering kernels on a grid of dimensionless momentum and energy transfer to describe thermal neutron scattering. The effects of chemical binding of nuclides, dynamics and structure of materials determine the peculiarities of neutron scattering at low incident neutron energies ($E < 1 - 10$ eV).

The sublibrary is organized by a nuclide (scatterer) in a given material. For example, in the ENDF/BV-II.1 TSL sublibrary, we have data for Be in beryllium oxide, O in beryllium oxide, C in Graphite, etc. In some cases, only the most important scatterer in a material has the evaluation. For example, we have $\text{HinH}_2$, or hydrogen in the light water, but there is no evaluation for $\text{OinH}_2$, implying that usage of the free gas model for thermal neutron scattering by oxygen in the light water is an acceptable approximation. Some evaluations have the data at one temperature: for example, data for thermal neutron scattering by H in liquid parahydrogen ($\text{H}_2$, I = 0) are given at $T = 20.0^\circ$K only. However, many evaluations are given for a number of temperatures $T$. For example, $S(\alpha, \beta; T)$ data for $\text{UinUO}_2$ (U in uranium dioxide) are given at eight different temperatures.

Discussion point:
The materials in the current ENDF libraries could be organized using the metaEvaluation markup in section V.B

When using, for example, ENDF/B-VII.1 TSL data, it is expected that nuclear data processing codes can read $S(\alpha, \beta, T)$ data and generate differential cross sections, $d^2\sigma(E, T)/dE\,d\Omega$, as well as the integral data (such as, integral cross sections $\sigma(E, T)$, average scattering cosine $\bar{\mu}(E, T)$, average $E'$, etc.) in proper physical units (barn per eV per sr, barn, eV, etc.) for incident neutron energies $E$ and neutron scattering with the energy $E'$ and scattering cosine $\mu$.

Discussion point:
It was felt at the May 2014 Paris meeting, that we should consider focusing on storing only legacy ENDF data and possibly the phonon spectrum $\rho(\omega)$ and defer all subsequent discussions until the formation of a special WPEC subgroup that can specifically deal with TSL issues.

A. Theoretical Background

Working in the first Born approximation, neutron scattering off one class of $N$ scatterers $j$ with scattering
lengths \( b_j \) is

\[
\frac{d^2 \sigma(E)}{d\Omega dE'} = \frac{k'}{2\pi\hbar} \sum_{j,j'} b_j b_{j'} \times \int_{-\infty}^{\infty} dt e^{-\omega t} \left( e^{-\vec{\kappa} \cdot \vec{R}_j(0)} e^{-\vec{\kappa} \cdot \vec{R}_j(t)} \right)
\]

(36)

This expression includes both elastic and inelastic scattering. Here \( \vec{R}_j(t) \) is the positions of the \( j^{th} \) scatterer at time \( t \), and \( k \), and \( k' \) are the incident and outgoing neutron wave numbers. The momentum and energy transfers between the scattering neutron and the collections of scatterers are \( \vec{\kappa} = \vec{k} - \vec{k}' \) and \( \omega = E' - E \). The implicit dependence on the material temperature is suppressed.

In the limit of a large number of scatterers we may make the replacement \( b_j b_{j'} \rightarrow \langle b_j b_{j'} \rangle \) and further, assuming no correlation between the scattering lengths of different nuclei,

\[
\langle b_j b_{j'} \rangle = \langle b_j \rangle \langle b_{j'} \rangle = \langle b \rangle^2 \text{ if } j' \neq j
\]

\[
\langle b_j b_j \rangle = \langle b^2 \rangle \text{ for } j = j'
\]

(37)

Using this, we define the coherent scattering cross section as \( \sigma_{coh} = 4\pi \langle b \rangle^2 \) and the incoherent as \( \sigma_{inc} = 4\pi (\langle b^2 \rangle - \langle b \rangle^2) \). Now, define an intermediate function

\[
I(\vec{\kappa}, t) = \frac{1}{N} \sum_{j,j'} \left( e^{-\vec{\kappa} \cdot \vec{R}_j(0)} e^{-\vec{\kappa} \cdot \vec{R}_j(t)} \right)
\]

(38)

The intermediate function can be computed from condensed matter theory assuming that the dynamics (described in terms of vibrational eigenmodes or phonon-type spectra) and structure (e.g., a certain order or correlations in the positions of scatterers in space) of the medium of interest are well understood.

In terms of the intermediate function, the scattering kernel is

\[
S(\vec{\kappa}, \omega) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} dt e^{-\omega t} I(\vec{\kappa}, t)
\]

(39)

Both \( S(\vec{\kappa}, \omega) \) and \( I(\vec{\kappa}, t) \) may include all \( j, j' \) in the sum in Eq. (36) or they may be broken out into \( j = j' \) (self) and \( j \neq j' \) (distinct) contributions. This is useful because the incoherent cross section only contains the self correlation between an atom at time \( t = 0 \). We have

\[
\frac{d^2 \sigma(E)}{d\Omega dE'} = \frac{k'}{k} \left[ \sigma_{coh} S(\vec{\kappa}, \omega) + \sigma_{inc} S_{\kappa}(\vec{\kappa}, \omega) \right]
\]

(40)

where

\[
S(\vec{\kappa}, \omega) = S_{\kappa}(\vec{\kappa}, \omega) + S_{\kappa}(\vec{\kappa}, \omega)
\]

(41)

**Discussion point:**

Reaction annotations could be used to split the scattering kernel into “self” and “distinct” parts. This is useful for a model-based evaluation where both components can be computed separately.

In practice, one assumes that we may average over orientation of the scatterers such that we can replace the directional dependence of \( \vec{\kappa} \) with a directionless \( \kappa \) dependence. Also, one uses the scattering kernel rewritten in terms of the dimensionless variables \( \alpha \) and \( \beta \) so

\[
S(\alpha, \beta) = \frac{k_B T}{\hbar} S(\kappa, \omega)
\]

(42)

where \( \alpha = \hbar^2 k^2/2Mk_B T = (E' + E - 2\sqrt{EE'})/Ak_B T \) and \( \beta = \hbar\omega/k_B T = (E' - E)/k_B T \).

**Discussion point:**

Whether we use \( \alpha \) and \( \beta \) or \( \kappa \) and \( \omega \), we have reduced the parametric dependence of the scattering kernel to three. These are \( \kappa \), \( \omega \), and an implicit material temperature dependence. This makes storing the scattering kernel directly in ENDF feasible.

**Resolution:**

We should stick to storing the kernel in terms of \( \alpha \) and \( \beta \) for backwards compatibility.

**Discussion point:**

New experiments from NCSU/RPI/ORNL collaboration will directly measure the \( d\sigma(E)/d\Omega dE' \). This is equivalent to measuring the full scattering kernel. Storing the covariance on the full scattering kernel may be unfeasible. Storing the covariance on data using the approximations and distinctions below may be feasible.

**Discussion point:**

ENDF thermal scattering data can have a large dynamic range. To accommodate this, ENDF manual recommends the following:

For down-scattering events with large energy losses and for low temperatures, \( \beta \) can be large and negative. The main contribution to the cross section comes from the region near \( \alpha + \beta = 0 \). Computer precision can become a real problem in these cases. As an example, for water at room temperature, calculations using equation (7.6) for incident neutron...
tron at 4 eV require working with products like $e^{80} \times 10^{-34}$. For liquid hydrogen at 20 Kelvin and for 1 eV transfers, the products can be $e^{300} \times 10^{-130}$. These very large and small numbers are difficult to handle on most computers, especially 32-bit machines. The LLN flag is provided for such cases: the evaluator simply stores $\ln(S)$ instead of $S$ and changes the interpolation scheme accordingly (that is, the normal log-log law changes to log-lin). Values of $S = 0.0$ like those found in the existing ENDF/B-III thermal files really stand for some very small number less than $10^{-32}$ and should be changed to some large negative value, such as -999.

Should we preserve this capability or store the $d\sigma(E)/dE'd\Omega$ directly, avoiding these precision problems?

The scattering kernel can be divided and simplified further by taking advantage of the elastic limit ($\omega \to 0$) or by making several approximations

- In the Gaussian approximation the “self” part of the scattering kernel may be written in terms of the material’s phonon frequency $\rho(\omega)$ and computed using the LEAPR approach

- The short collision time approximation

\textbf{Requirement 55: Thermal scattering kernel}

55.1 Allow TSL data to be broken out into separate reactions as specified by the evaluator. Each reaction is treated independently for the purposes of neutron transport.

55.2 Denote the energy range for which these data are used. The $E_{\text{max}} = 5$ eV limit in ENDF is convention and has no general physical justification.

55.3 Allow reactions to be annotated by combinations of self, distinct, coherent, incoherent, tsl_elastic and tsl_inelastic labels. Because of this flexibility, care will need to be taken by evaluators to ensure that double counting does not occur.

55.4 All reaction data contained in $<$dcrossSection_dOmega_dE$>$ or $<$dcrossSection_dOmega$>$ elements, depending on the evaluators needs.

55.5 If the reaction data are be broken out by scatterer (e.g., HinH2O), the stoichiometric fraction of each class of scatterer must be given.

55.6 The scattering kernel $S(\alpha, \beta, T)$ can be given as an interpolation table or using one of the approximations or distinctions given below in subsections XIV.B.1-XIV.B.3

55.7 Coherent or incoherent cross sections are associated with their respective scattering kernels.

\textbf{Discussion point:}

Annotations might also be used to denote “tsl_elastic” and “tsl_inelastic” data as TSL data does not have the same two-body kinematics of higher energy data. When $\omega \to 0$, $E = E'$ in the lab frame and the center of mass frame is meaningless.

\textbf{B. Gaussian approximation of the self part of the scattering kernel}

By making the so-called Gaussian approximation to the scattering kernel [Behr 2010, VanHove 1954] the self part of the scattering kernel can be written in terms of the material phonon spectrum $\rho(\omega)$:

\begin{equation}
S_s(\alpha, \beta) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dt e^{i\beta t} e^{-\gamma(t)}
\end{equation}

where

\begin{equation}
\gamma(t) = \alpha \int_{-\infty}^{\infty} d\omega \rho(\omega) \left( 1 - e^{-i\omega t} \right) = \frac{e^{-\omega/2}}{2\omega \sinh(\omega/2)}
\end{equation}

This Fourier transform is coded in R.E. MacFarlane’s LEAPR module of NJOY [MacFarlane 2012] and can be used to generate the inelastic scattering kernel.

We note that if the full scattering kernel is well approximated by only the self term and in this Gaussian approximation, the entire scattering kernel can be specified with the phonon spectrum. This spectrum may have a discrete portion and/or a continuous portion. Nevertheless, this offers us a compact way to encapsulate the scattering kernel and it provides us with a two-dimensional object that we can specify covariance on.

\textbf{Requirement 56: Gaussian self-part scattering kernel}

56.1 An encapsulating element that specifies that these data is the self part of the scattering kernel only.

56.2 The phonon spectrum $\rho(\omega)$ as a discrete and/or continuous distribution.

56.3 Optionally, a link to the covariance on the phonon spectrum.
Following (Cacuci [2010]), we can expand the time-dependent part of the scattering kernel to arrive at the phonon expansion:

\[
e^{-\gamma(t)} = -\alpha \lambda_s \sum_{n=0}^{\infty} \frac{1}{n!} \left[ \alpha \int_{-\infty}^{\infty} d\omega P_s(\omega)e^{-\omega/2}e^{-i\omega t} \right]^n
\]

(45)

where \( P_s(\omega) = \rho(\omega)/2\omega \sinh(\omega/2k_BT) \) and the Debye-Waller coefficient is

\[
\lambda_s = \int_{-\infty}^{\infty} d\omega P_s(\omega)e^{-\omega/2}
\]

(46)

The \( n^{th} \) term in this expansion is identified with the number of phonons involved in the collision (Cacuci, 2010). This expansion allows us to arrive at two approximations given in the ENDF format manual that must be grandfathered into the new format: elastic coherent, elastic incoherent and the short collision time approximations. In the case of the elastic (in)coherent data, we set \( n \) to \( 0 \) which then forces \( \beta \to 0 \) and simplifies the phonon expansion to the first term, the zero phonon limit. Alternatively, in the limit of large \( n \), we arrive at the short collision time approximation (call this because the large number of collisions implies a short time for each individual collision).

1. Coherent Elastic Scattering

For crystals (polycrystalline materials), the information about the crystal structure is expressed in terms of the so-called Bragg edges (a discrete set of energies \( E_j \sim 1 \text{ meV} - 1 \text{ eV} \)) and a set of crystallographic structure factors \( s_j \) associated with \( E_j \) and a neutron scatterer in a crystal unit cell. In addition, one has to estimate the temperature dependent Debye-Waller coefficient \( W' \) (in the units of \( \text{eV}^{-1} \)). Then it is possible to generate the data structure that can be used to generate the contribution of coherent elastic neutron scattering into the thermal neutron scattering kernel, scattering cross sections, etc., for a given scatterer in the polycrystal.

Figure 38 compares the different elastic scattering prescriptions for two different forms of carbon. Here one can clearly see the Bragg edges in the elastic cross section.

In the early 1990’s, parameterized coherent and incoherent elastic scattering were added to ENDF format. Neutrons can only elastically scatter coherently off of regular substances such as crystals. The differential cross section for such scattering can be written (Trkov, 2009):

\[
\frac{d^2\sigma}{dE'd\Omega}(E,T) = \frac{1}{2\pi E} \sum_{i=1}^{E_j<E} f_i e^{-4W'E} \delta(\mu-\mu_i) \delta(E-E')
\]

(47)

The quantity actually given in the file is \( s_i(T) = S(E_i,T) \) which the ENDF manual states is conveniently represented as a stairstep function with breaks at the Bragg edges \( E_i \) using histogram interpolation. Here, we must store the structure factor \( s_i(T) = S(E_i,T) \) tables in ENDFs MF=7 (note these factors are given as a histogram in ENDf, hence the notation above).

Alternatively, this cross section can be written (Cacuci, 2010):

\[
\frac{d^2\sigma}{dE'd\Omega}(E,T) = \frac{\sigma_{coh}}{E} \sum_{i=1}^{E_i<E} f_i e^{-4W'E} \delta(\mu-\mu_i) \delta(E-E')
\]

(48)

The \( f_i \) are material dependent and related to the crystallographic structure factors.

Requirement 57: Coherent Elastic Scattering

57.1 An elastic channel reaction designator that includes the annotations tsl_elastic and coherent.

57.2 <dCrossSection_dOmega> element containing this data.

57.3 An list of <interp2d> elements containing the structure factor \( S(E,T) \) in the <distribution> element. The ENDF manual requires the interpolation in \( E \) to be a histogram and it is unclear whether there is a need to relax this requirement. The Bragg edges are the histogram boundaries. With this requirement, \( S(E,T) = S(E_i,T) \equiv s_i(T)/E \)

57.4 Alternatively, specify the Bragg edges \( E_i \) and
2. Incoherent Elastic Scattering

For partially ordered systems, the incoherent approximation to elastic scattering is given by

$$\frac{d^2\sigma}{dE'd\Omega}(E,T) = \frac{\sigma_b}{4\pi} e^{-2EW'(T)(1-\mu)} \delta(E-E')$$

(50)

where:

- $\sigma_b$ is the characteristic bound cross section (barns),
- $W'$ is the Debye-Waller integral divided by the atomic mass (ENDF data are given in eV$^{-1}$),

and all the other symbols have their previous meanings.

The integrated cross section is easily obtained:

$$\sigma(E) = \frac{\sigma_b}{2} \left(1 - e^{-4EW'} \right)$$

(51)

Note that the limit of $\sigma$ for small $E$ is $\sigma_b$.

3. Incoherent Inelastic Scattering in the Short Collision Time Approximation

In the short collision time limit, we have \cite{Trkov2009}

$$S(\alpha, \beta) = \exp \left[ -\frac{(\alpha-|\beta|)^2}{4\alpha T_{\text{eff}}(T)} - \frac{|\beta|}{T} \right] \frac{1}{\sqrt{4\pi\alpha T_{\text{eff}}(T)}}$$

(52)

**Requirement 58: Incoherent Elastic Scattering**

- **58.1** An elastic channel reaction designator that includes the annotations *tsl_elastic* and *incoherent*.
- **58.2** `<dCrossSection_d0Mega>` element containing these data.
- **58.3** An `<interp1d>` element collecting $W'$ the Debye-Waller integral divided by the atomic mass as a function of temperature.
- **58.4** The bound cross section $\sigma_b$, with units.
- **58.5** Optional link to `<covariance>` data on $W'$.
- **58.6** Optional link to `<covariance>` data of $\sigma_b$. This is a $1 \times 1$ matrix, but could be correlated with $W'$'s covariance.
- **58.7** Only one `<form>` of these data is currently possible.
- **58.8** These data must be given in the lab frame as the center of mass frame is meaningless for TSL data.

**Requirement 59: Short Collision Time Approximation**

- **59.1** Something to denote that the short collision time approximation is used.
- **59.2** The effective temperature $T_{\text{eff}}(T)$

XV. ADDITIONAL DERIVED DATA ELEMENTS FOR APPLICATIONS

One of the goals for the new structure is to be able to store processed or derived data for specific applications (see requirement 3.2). Here we attempt to list some of the more important cases and their requirements, but we realize that our list is far from complete.

Processed data are a representation of the data in a form required by transport (or other application) codes. This kind of data is often needed for intercomparison between labs. However, each institution generally processed evaluated data into a form mainly acceptable by applications managed by that institution and these forms depend on each institution’s unique needs. Despite this and the fact that needs of institutions and broader user communities evolve, some needs are apparent even at this stage in the development of a successor format.

The elements and requirements we describe below are located in a few distinct places in the data hierarchy:

- `<reaction>`: Store evaluated and any processed/derived data specific to a reaction.
- `<derivedReactions>`: Store derived reactions such as total, total inelastic, non-elastic computable from sum rules. Also they store any reactions `<reaction>`s created during any post-evaluation processing, to distinguish them from evaluated `<reaction>`s in the `<reactions>` branch.
- `<derivedTransportData>`: Store other derived data not tied to a specific `<reaction>` but is needed for transport, such as total transfer matrices. It is not to be confused with `<styles>` contents which define things like group structures and fluxes needed in the generation of processed data.

We comment that there are cases even in ENDF where there is incomplete information known (e.g., for fission we...
may only know the total emitted energy for some of the minor actinides) so these “derived” or “processed” data may be all we know. Therefore, in any given data set, the user is urged to check the <listOfDerivedFromLinks> within a data element to determine what is original and what isn’t.

A. General Transport Data

1. Product average kinetic energy and forward momentum

A product’s average kinetic energy $\bar{E}'(E)$ is defined as

$$\bar{E}'(E) = \int_0^\infty dE' \int_{-1}^1 d\mu E' P(\mu, E'|E)$$

It is convenient to store this quantity in a linear tabulated form as it is used to check energy balance.

A product’s average forward momentum $\bar{p}'_f(E)$ is defined as

$$\bar{p}'_f(E) = \int_0^\infty dE' \int_{-1}^1 d\mu \mu p'(E') P(\mu, E'|E)$$

where $p'(E')$ is the magnitude of the product’s momentum as a function of its kinetic energy $E'$ and the $\mu$ factor in the equation yields the projection of the product’s momentum along the direction of the projectile. It is also convenient to store this quantity in a linear tabulated form.

2. $\bar{\mu}_{lab}(E)$

In ENDF, the average forward scattering angle in the lab frame of a reaction product, $\bar{\mu}_{lab}(E) = \int d\mu_{lab}\mu_{lab}P(\mu_{lab}|E)$ is given in the MT=251 file. Although one might view $\bar{\mu}_{lab}(E)$ as a derived quantity, it can be measured experimentally and so may have covariance data associated with it. We note that the data typically given in evaluated files are in the center of mass frame, so care must be taken when performing the frame change.

Requirement 60: Average kinetic energy and forward momentum

60.1 Shall store the projectile’s averaged projectile kinetic energy and momentum within the appropriate <reaction>
60.2 Shall store the products’ averaged projectile kinetic energy and momentum within the appropriate <reaction>’s <product> element
60.3 The averaged projectile kinetic energy and momentum shall be stored as <interp1d> tables

3. CDF’s from PDF’s

Given a probability density function $PDF(E)$, one can define the cumulative distribution functions (CDF) as

$$CDF(E) = \int_{-\infty}^E dE' PDF(E')$$

Monte Carlo codes use these cumulative distribution functions (CDF)’s to convert uniformly sampled random numbers, $x$, on the interval $x \in [0, 1]$ to samples consistent with the underlying PDF through

$$E = CDF^{-1}(x)$$

For Monte Carlo transport, it is convenient to pre-compute the cumulative distribution functions (CDF) for each PDF. However, this is a fast calculation and pre-calculation is not essential, but may be advantageous. Therefore, we require a markup for CDF’s as derived data to be located in the same enclosing element as its PDF.

Requirement 62: CDFs

62.1 A markup for CDF’s as derived data, located in the same enclosing element as its PDF.

4. Probability tables in the URR

The unresolved resonance region parameters represent the average behaviors of resonances that cannot be resolved experimentally. Using techniques such as implemented in the PURR module of NJOY [MacFarlane, 2012] or the PURM module of AMPX [Dunn, 2002], one can convert these average parameters into the probability for a particular reaction cross section $\sigma_x$ as a function of incident energy, $P(\sigma_x|E)$.

Requirement 63: URR probability tables

63.1 The URR probability tables shall be stored in the appropriate <reaction>
63.2 The URR probability tables shall be stored as \texttt{<interp2d>} objects

**Discussion point:**
The GRUCON code [GRUCON] can compute the conditional probabilities \( P(\sigma_i|\sigma_j, E) \), properly accounting for the correlations in the probabilities of all the reaction cross sections. If we are to support these, they naturally would go in the evaluation-wide \texttt{<derivedTransportData>}.  

## B. Grouped Transport Data

The largest use of data in ENDF format is modeling particle transport (usually neutrons, but users are sometime interested in transporting charged particles including electrons). To model the transport of particles such as neutrons one uses codes that solve the Boltzmann equation. The Boltzmann equation is also known as the transport equation and can be solved a variety of ways including with Monte-Carlo techniques, by discretizing it and solving the resulting matrix equation (also known as deterministic transport) or rarely by using the method of characteristics.

In a simplified 1-dimensional form with one target species and only elastic scattering, the Boltzmann equation is

\[
\frac{1}{v} \frac{\partial f(E)}{\partial t} + \frac{\partial f(E)}{\partial x} + \rho(E) \sigma(E) f(E) = \int dE' \rho(E') P(E' \rightarrow E) f(E')
\]  

where \( E \) is the kinetic energy of the projectile, \( f(E) = f(x, E, t) \) is the flux of a particle being transported (called the projectile), \( \rho \) is the target density, \( \sigma(E) \) is the cross section between the projectile and the target and \( P(E' \rightarrow E) \) is the probability density function (pdf) for scattering from energy \( E' \) to \( E \). In this simplified form, the angular dependence is ignored.

For both Monte Carlo and deterministic data, it is convenient to convert cross section resonance data to a linear tabulated form and pre-heat the tabulated cross sections to pre-defined temperatures as these are computationally intensive calculations.

For deterministic transport the energy variable \( E \) must be judiciously discretized by flux-weight averaging elements of the Boltzmann equation over energy bins (also called groups). For example, the cross section becomes for group \( i \) \((E_i < E < E_{i+1})\),

\[
\sigma(E) \Rightarrow \sigma_i = \frac{\int_{E_i}^{E_{i+1}} dE \sigma(E) \Phi(E)}{\int_{E_i}^{E_{i+1}} dE \Phi(E)}
\]  

and the \( P(E' \rightarrow E) \) factor in the right-hand-side of Eq. 57 becomes,

\[
\sigma(E) P(E' \rightarrow E) \Rightarrow \text{TM}_{io} = \frac{\int_{E_i}^{E_{i+1}} dE' \int_{E_i}^{E_{i+1}} dE \sigma(E') P(E' \rightarrow E) \Phi(E')}{\int_{E_i}^{E_{i+1}} dE' \Phi(E')}
\]  

where \text{TM}_{io} \) (called the transfer matrix) expresses the fact that a projectile in energy group \( o \) produces outgoing particles into energy group \( i \). Equation 59 is a simplified version; a more accurate \text{TM}_{io} includes Legendre expansion, a multiplicity for the number of outgoing particles \( M(E') \) and a conservation factor \( W(E') \). The conservation factor is either 1 or \( E' \), depending on whether the number or average energy of the outgoing particles is to be conserved for each group.

Ideally, the flux \( \Phi(E) \) that we weight with should be the flux that results from solving the transport equation, \( f(E) \). Clearly this presents a “chicken and egg” problem since we can’t weight with a flux we haven’t yet computed. Therefore, in practice various techniques have been developed to make judicious choices for the weighting fluxes. Bondareko and multi-band treatments are some of the more popular choices in the application community.

**Discussion point:**
We will need a follow on project that can determine the requirements for the various transport approaches beyond what we list here.
The more accurate transfer matrix is,

\[ \text{TM}_{io,l} = \frac{\int_{E_o}^{E_{o+1}} dE P_i(\mu) \int_{E_i}^{E_{i+1}} dE' M(E') \int_{E_i}^{E_{i+1}} dE W(E') \sigma(E') P(E' \rightarrow E, \mu) \Phi(E')}{\int_{E_o}^{E_{o+1}} dE' \Phi(E')} \]

(60)

The grouped deterministic transport representation of Eq. 57 with Legendre notation (albeit still simplified) is,

\[ \frac{1}{v_i} \frac{\partial f_i}{\partial t} + \frac{\partial f_i}{\partial x} + \rho \sigma_i f_i = \rho \sum_o \text{TM}_{io,l} f_o \]

(61)

In addition to \( \sigma_i \) and \( \text{TM}_{io,l} \), it is convenient to group other quantities; namely, product multiplicity, energy dependent \( Q \) data (i.e., \( Q(E) \)), projectile kinetic energy and momentum, product average kinetic energy and forward momentum, and inverse speed (i.e., \( 1/v_i \)). It is also convenient to have several of the quantities in linear tabulated form; namely, projectile momentum, and product kinetic energy and forward momentum.

All of these quantities can be classified by their dimension and location with the structure as,

**Requirement 64: Grouped transport data**

64.1 With the exception of \( \text{TM}_{io,l} \), all grouped quantities can be stored as a list (i.e., an array of dimension 1) while \( \text{TM}_{io,l} \) requires an array of dimension 3.

64.2 The following describes where each element shall be stored.

64.2.1 The inverse speed shall be stored within the <evaluation> element.

64.2.2 The cross section, energy dependent \( Q \) data, and projectile kinetic energy and momentum shall be stored within the appropriate <reaction> element.

64.2.3 The transfer matrix, multiplicity, and product average kinetic energy, forward momentum and average forward scattering angle shall be stored within the appropriate <product> element.

64.3 The group structure and flux weights are stored in the <styles> elements for the evaluation.

64.4 Grouped data must clearly state the flux weighting and group structure used to derive the data as well as the original data that was grouped.

In addition, both the number and energy conserving \( \text{TM}_{io,l} \), should be stored as separate elements (with unique names), as both may be needed.

In the following sections, the cross section weighted average \( \langle A \rangle_i \) of a quantity \( A(E) \) is defined as,

\[ \langle A \rangle_i = \frac{\int_{E_i}^{E_{i+1}} dE A(E) \sigma(E) \Phi(E)}{\int_{E_i}^{E_{i+1}} dE \Phi(E)} \]

(62)

Note, we may want to define it as

\[ \langle A \rangle_i = \frac{\int_{E_i}^{E_{i+1}} dE A(E) \sigma(E) \Phi(E)}{\int_{E_i}^{E_{i+1}} dE \sigma(E) \Phi(E)} \]

(63)

1. Inverse speed

The cross section weighted average inverse speed is defined as \( \langle 1/v \rangle_i \), where \( v(E) \) is the projectile’s velocity as a function of its energy.

2. Multiplicity

The cross section weighted average multiplicity is defined as \( \langle m \rangle_i \), where \( m(E) \) is the product’s multiplicity as a function of its energy.

3. Q-value

The cross section weighted average Q-value is defined as \( \langle Q \rangle_i \), where \( Q \) is a reaction’s Q-value as a function of its energy.

4. Projectile kinetic energy and momentum

The cross section weighted average projectile kinetic energy is defined as \( \langle E \rangle_i \).

The projectile’s momentum as a function of its kinetic energy \( E \) is written as \( p(E) \) and it is convenient to store it in a linear tabulated form. The cross section weighted average momentum is defined as \( \langle p \rangle_i \).

C. Production data

Production cross sections are used to store the total probability for producing a radioactive daughter, irrespective of what reaction or reactions were involved in
creating that daughter. They are often used for modeling the activation of a material following irradiation. In ENDF they are also often used when a reaction can produce an isomer, to give the portion of the total cross section going to that isomer.

The production cross section \( \sigma_p \) for a specific product is a derived quantity that can be computed by:

\[
\sigma_p = \sum_r \sigma_r \cdot M_r
\]

where for each reaction \( r \), \( \sigma_r \) is the cross section and \( M_r \) is the multiplicity of product ‘p’.

Production cross sections are redundant and may be deprecated eventually, but they should continue to be supported for backwards-compatibility. The current solution (in GND) is to store production cross sections inside a ‘production reaction’ element. This is a special type of reaction that only contains a cross section and a single outgoing particle (the product).

Gas production is one specific form of production data that is used to understand the total amount of hydrogen and helium produced. The hydrogen gas production cross section is a sum of the proton, deuteron and triton production cross sections. The helium gas production cross section is a sum of the \(^3\)He and \( \alpha \)-particle production cross sections. The total gas production cross section is a sum of the hydrogen and helium gas production cross sections. In ENDF, gas production cross sections are given in MT’s 203-207.

**D. Damage cross sections**

Radiation damage is obviously important in many nuclear applications. As energetic particle transits a material, it ionizes atoms and molecules and dislocates atoms. The modeling of the interactions is a complicated interplay of nuclear, atomic and materials physics. Radiation damage is also an under developed part of the ENDF format. While radiation damage data needs to be archived, how it will fit within the hierarchy is a subject of future investigation. The reason is that traditional measures of damage such as displacements per atom (DPA) are crude and new approaches are being developed that better incorporate the materials properties. These approaches are being examined as part of a recently formed IAEA Coordinated Research Project (Stoller 2012).
FIG. 39 The `<functionDef>` Outline of the `functionDef` scheme for prototyping new functional forms and interpolation schemes.

```xml
<documentation> element.

67.3.2 The definition of the `<functionDef>` in a `<definition>` element. This includes lists of expected inputs (including name, unit and range) and outputs (again including name, unit and range) as well as the expression(s) detailing the new `<functionDef>`, preferably in a common format such as MathML.

67.3.3 Unit tests so that the implementation `<functionDef>` can be tested against the authors’ expected results. The unit tests, in `<unitTest>` elements, contain the inputs, expected outputs and any documentation/comment describing each test.

ACKNOWLEDGMENTS

FIXME : We wish to acknowledge valuable review and feedback from all the contributors. For a full list, see appendix D.

AECL

BNL Under DOE contract DE-SC0012704

IAEA

LANL

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NNL

NRG

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Appendix A: Graphical notation

In order to illustrate points and various proposed layouts of the top level hierarchy, we used diagrams written in the Crow’s Foot notation for Entity Relationships (Barker 1990) as implemented in the yEd code (yEd 2015). The Crow’s Foot notation is one type of diagram used to show entity relationships in the Universal Modeling Language (UML) (UML.org 2015).

TABLE III Crow’s Foot notation elements. Any of these may be parent elements or contained within other parent elements.

<table>
<thead>
<tr>
<th>Element</th>
<th>Description</th>
</tr>
</thead>
<tbody>
<tr>
<td>High level element</td>
<td></td>
</tr>
<tr>
<td>Mid level element</td>
<td></td>
</tr>
<tr>
<td>Low level element</td>
<td></td>
</tr>
<tr>
<td>Link to an element</td>
<td></td>
</tr>
<tr>
<td>Element with attributes</td>
<td></td>
</tr>
<tr>
<td>Condition</td>
<td></td>
</tr>
<tr>
<td>Sample container</td>
<td></td>
</tr>
</tbody>
</table>

TABLE IV Connections denoting the number of child elements that are contained within a parent element within the Crow’s Foot notation.
Appendix B: List of <distribution> requirements

Here we list the various energy distributions, angular distributions and energy-angle distributions that are used in the ENDF format and that must be implemented in the new format. We do not prove requirements for any of these other than to state that each element listed below should include member data such that data in the new format is equivalent to the ENDF data. We assume the implementation of each of these is clear enough that this listing is sufficient.

FIXME NEED GND 1.6 NAMES

<table>
<thead>
<tr>
<th>Element names</th>
<th>Description</th>
<th>ENDF equivalent</th>
</tr>
</thead>
<tbody>
<tr>
<td>LegendrePointwise</td>
<td>Legendre moment data</td>
<td>MF=4, LTT=1, LI=0</td>
</tr>
<tr>
<td>LegendrePiecewise</td>
<td></td>
<td></td>
</tr>
<tr>
<td>Pointwise</td>
<td>table</td>
<td>MF=4, LTT=2, LI=0</td>
</tr>
<tr>
<td>Piecewise</td>
<td>isotropic</td>
<td>MF=4, LTT=0, LI=1</td>
</tr>
<tr>
<td>Isotropic</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

TABLE V Angular probability density functions $P(\mu|E)$ supported by ENDF that must be implemented in the new format.

<table>
<thead>
<tr>
<th>Element names</th>
<th>Description</th>
<th>ENDF equivalent</th>
</tr>
</thead>
<tbody>
<tr>
<td>Pointwise,</td>
<td>tabulated</td>
<td>MF=5, LF=1</td>
</tr>
<tr>
<td>Piecewise</td>
<td></td>
<td></td>
</tr>
<tr>
<td>GeneralEvaporation</td>
<td>general evaporation spectrum</td>
<td>MF=5, LF=5</td>
</tr>
<tr>
<td>SimpleMaxwellianFission</td>
<td>Maxwellian spectrum</td>
<td>MF=5, LF=7</td>
</tr>
<tr>
<td>Evaporation</td>
<td>evaporation spectrum (similar to a Maxwellian but different pre-factor)</td>
<td>MF=5, LF=9</td>
</tr>
<tr>
<td>Watt</td>
<td>Watt spectrum (a boosted Maxwellian)</td>
<td>MF=5, LF=11</td>
</tr>
<tr>
<td>NBodyPhaseSpace</td>
<td>N-body phase space</td>
<td>MF=6, LAW=6</td>
</tr>
<tr>
<td>MadlandNix</td>
<td>Madland-Nix model parameterization</td>
<td>MF=5, LF=12</td>
</tr>
</tbody>
</table>

TABLE VI Energy probability density functions $P(E'|E)$ supported by ENDF that must be implemented in the new format.
<table>
<thead>
<tr>
<th>Element names</th>
<th>Description</th>
<th>ENDF equivalent</th>
</tr>
</thead>
<tbody>
<tr>
<td>unknown</td>
<td>unknown or unspecified</td>
<td>MF=6, LAW=0</td>
</tr>
<tr>
<td>pointwise</td>
<td>tabular</td>
<td>MF=6, LANG=11-15</td>
</tr>
<tr>
<td>piecewise</td>
<td>N/A, stored as two-body, angular distribution</td>
<td>MF=6, LAW=2</td>
</tr>
<tr>
<td></td>
<td>N/A, treated as isotropic</td>
<td>MF=6, LAW=3</td>
</tr>
<tr>
<td></td>
<td>uncorrelated with isotropic angular distribution, delta function for energy</td>
<td></td>
</tr>
<tr>
<td>recoil</td>
<td>two-body recoil</td>
<td>MF=6, LAW=4</td>
</tr>
<tr>
<td>CoulombElastic</td>
<td>large angle part of charged particle elastic</td>
<td>MF=6, LAW=5</td>
</tr>
<tr>
<td>KalbachMann</td>
<td>Kalbach-Mann systematics for pre-equilibrium emission</td>
<td>MF=6, LANG=2</td>
</tr>
<tr>
<td></td>
<td>N/A, frame can be specified on most elements so this is redundant</td>
<td>MF=6, LAW=7</td>
</tr>
<tr>
<td>??</td>
<td>Legendre moments as a function of incoming and outgoing energies ( f_\ell(E', E) )</td>
<td>MF=6, LANG=1</td>
</tr>
<tr>
<td>uncorrelated</td>
<td>uncorrelated product of an energy and an angular distribution</td>
<td>MF=4 and MF=5</td>
</tr>
</tbody>
</table>

TABLE VII Energy-Angle probability density functions \( P(\mu, E'|E) \) supported by ENDF that must be implemented in the new format.
Appendix C: Terminology

**A:** The total number of protons and neutrons in a given nucleus.

**abundance:** For isotopes that occur naturally, the abundance values are proportional to the probability of finding these isotopes and normalized so that the sum of the abundances for all the isotopes of a given chemical element is equal to 100. The source is [Holden 2004].

**α decay:** The emission of a $^4\text{He}$ nucleus ($\alpha$ particle).

**α particle:** A $^4\text{He}$ nucleus, that is, a nucleus made up of 2 neutrons and 2 protons.

**AMPX:** AMPX ([Dunn 2002](http://www.ippe.obninsk.ru/podr/cjd/)) is a modular system of computer programs developed at ORNL with primary emphasis on processing neutron and photon evaluations to produce cross-section libraries for nuclear systems analysis.

**application programming interface (API):** The set of routines, protocols, and tools for building software applications. The API specifies how software components should interact.

**atom:** An atom is the smallest unit of matter that defines the chemical elements. Every solid, liquid, gas, and plasma is made up of neutral or ionized atoms.

**atomic mass unit (amu):** The atomic mass unit (amu) is defined so that 1 amu is equal to the mass of a $^{12}\text{C}$ atom divided by 12.

**attribute, XML:** These normally are used to describe XML elements, or to provide additional information about elements.

**Auger electrons:** Electrons that are produced when a vacancy in an orbit A is filled by an electron from the orbit B and an electron from an orbit C is ejected. These electrons are labeled by the three orbits that are involved in the production. For instance, AugerKL2L3 means that the original vacancy was in the K orbit, which was filled by an electron in the L2 orbit and the ejected electron came from the L3 orbit. Coster-Kronig transitions are a special type of Auger electrons where the last two orbit are part of the same shell. Following nuclear decay, vacancies in the electron orbits are produced, which are filled by the emission of X-rays and electrons. Often, instead of listing the energy and intensity for each Auger electron, average intensities and sum intensities are given. For instance, the intensity of the Auger K electrons is the sum of the intensities for all the KBC Auger electrons.

**β− decay:** The transformation of one neutron inside a nucleus into a proton plus an electron and an antineutrino: $n \rightarrow p + \bar{\nu}_e$.

**β+ decay:** The transformation of one proton inside a nucleus into a neutron plus a positron and a neutrino: $p \rightarrow n + \nu_e$.

**β−delayed particle emission:** The emission of a nucleon, nucleons or a nucleus following β− decay. For proton rich nuclei, the emission of a proton following β+ decay and electron capture has been observed. For neutron rich nuclei, the release of one or two neutrons following β− decay is possible. The emission of α particles has been observed for some nuclei in all types of β decay. Also, for a few nuclei, fission can take place following β+ decay and electron capture.

**branching ratio:** The probability of a certain event occurring when multiple events are possible.

**bremstrahlung:** Literally “breaking radiation”. The process of electromagnetic radiation when a charged particle is accelerated or decelerated.

**BROND:** The Russian Evaluated Neutron Data Library developed at Center Jadernykh Dannykh (CJD) in Obninsk, Kaluga Region, Russia. Further information available at [http://www.ippe.obninsk.ru/podr/cjd/](http://www.ippe.obninsk.ru/podr/cjd/).

**CALEND:** The CALENDF Nuclear Data Processing System is used to convert the evaluation defining the cross-section in ENDF format (i.e. the point-wise cross-sections and/or the resonance parameters, both resolved and unresolved) into forms useful for applications. Those forms used to describe neutron cross-section fluctuations correspond to “cross-section probability tables”, based on Gauss quadratures and effective cross-sections. CALENDF also provides capabilities for group collapsing, for merging of several nuclei and for temperature interpolation; these calculations are based on data probability table description. CALENDF is developed by the Commissariat a l’Energie Atomique, Centre de Saclay.

**CENDL:** Chinese Evaluated Data Library is an evaluated nuclear reaction data library developed at the Chinese Nuclear Data Center (CNDC), Beijing, China in support of Chinese nuclear applications.

**channel:** is context sensitive concept. In resonance region (and anywhere else where we are using the R-matrix formalism), a channel has a specific meaning as all the quantum numbers needed to uniquely

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[Dunn 2002](http://www.ippe.obninsk.ru/podr/cjd/)
denote a quantum state. For a two-body reaction, that usually means \( c = \{\alpha, s, \ell, J \} \). For N-body reactions, we use the channel more loosely since in these cases many processes can lead to the same reaction products. For this reason, we will try to avoid using the term “channel” when discussing N-body reactions.

**cluster decay:** The emission of a nucleus heavier than an \( \alpha \) particle, for instance \( ^{14}\text{C} \). Branching ratios for this decay tend to be very small, due to the large Coulomb barrier encountered by the cluster and its very small pre-formation factor, that is, the probability of finding the cluster formed inside the nucleus.

**coherence, quantum mechanical:** Two quantum mechanical wave functions are coherent if they have a constant phase difference and the same frequency (and therefore energy).

**Compton scattering:** Compton scattering is the inelastic scattering of a photon by a quasi-free charged particle, usually an electron. It results in a decrease in energy (increase in wavelength) of the photon (which may be an X-ray or gamma ray photon), called the Compton effect. In the ENDF format, is is termed incoherent elastic photoatomic scattering and the data stored is an incoherent scattering function which modifies Klein-Nishina formula.

**conversion electrons (CE):** An electron released from the atomic shell by transferring the energy of a gamma quantum emitted from the same nucleus to this electron. The kinetic energy of the conversion electron is equal to the energy of the gamma quantum reduced by the binding energy of the electron.

**correlation matrix:** A matrix that describes the correlations between parameters in a covariance matrix. The correlation matrix is defined in section VI.A, equation (17).

**Coster-Kronig transition:** The Coster-Kronig transition is a special case of the Auger process in which the vacancy is filled by an electron from a higher subshell of the same shell. If, in addition, the electron emitted (the “Auger electron”) also belongs to the same shell, one calls this a super Coster-Kronig transition.

**Coulomb scattering:** See Rutherford scattering.

**covariance matrix:** A covariance matrix (also known as dispersion matrix or variance-covariance matrix) is a matrix whose element in the \( i \), \( j \) position is the covariance between the \( i^{th} \) and \( j^{th} \) elements of a random vector (that is, of a vector of random variables).

**covariance matrix, relative:** The covariance matrix, scaled by the original data. See equation (14) for a precise definition.

**cross section:** If a beam of particles (or photons) enters a thin layer of material (thickness \( dx \)), then the particle number \( N \) will be reduced by \( dN = -\mu N dx \), where \( \mu \) is the attenuation coefficient. To describe the attenuation coefficient in a way independent of the material density, one introduces the cross section \( \sigma = \mu/n \), where \( n \) is the numerical density (number of atoms per volume) of the material. \( \sigma \) has the dimension of an area; it expresses the likelihood of interaction between particles.

**cross section, capture:** The capture cross section is the reaction cross section for the process in which a projectile is captured by a target and one or more photons is emitted.

**cross section, differential:** FIXME

**cross section, elastic:** FIXME

**cross section, fission:** FIXME

**cross section, gas production:** The sum of the production cross sections for all reactions which produce any isotope of hydrogen and helium nuclei.

**cross section, inelastic:** FIXME

**cross section, infinite dilution:** FIXME

**cross section, integral:** FIXME

**cross section, macroscopic:** FIXME

**cross section, non-elastic:** The total cross section minus the elastic cross section.

**cross section, production:** FIXME

**cross section, sum rule:** Any addition rule that relates different cross sections for the same target–projectile system.

**cross section, total:** The sum of all reaction cross sections and the elastic cross section.

**Crow’s Foot Notation:** A notation used in Entity-Relationship modeling. Crow’s foot diagrams represent entities as boxes, and relationships as lines between the boxes. Different shapes at the ends of these lines represent the cardinality of the relationship.
cumulative distribution function (CDF): In probability theory and statistics, the cumulative distribution function (CDF), or just distribution function, describes the probability that a real-valued random variable $x$ with a given probability distribution will be found to have a value less than or equal to $x$. In the case of a continuous distribution, it gives the area under the probability density function from minus infinity to $x$.

delayed fission gamma spectrum (DFGS): The energy spectrum of emitted gammas from the beta (and other) decays of the prompt fission fragments.

delayed fission neutron spectrum (DFNS): The energy spectrum of emitted neutrons from the beta (and other) decays of the prompt fission fragments.

Delbrück scattering: Delbrück scattering is the coherent elastic scattering of photons in the Coulomb field of heavy nuclei.

deuterium: A neutral atom with 1 electron and a deuteron as a nucleus. An isotope of the hydrogen atom.

deuteron: A $^2\text{H}$ nucleus, that is, a nucleus made up of 1 neutron and 1 proton. Not to be confused with deuterium.

Digital Object Identifier (DOI): A digital object identifier (DOI) is a character string (a “digital identifier”) used to uniquely identify an object such as an electronic document. Metadata about the object is stored in association with the DOI name and this metadata may include a location, such as a URL, where the object can be found. The DOI for a document remains fixed over the lifetime of the document, whereas its location and other metadata may change. Referring to an online document by its DOI provides more stable linking than simply referring to it by its URL, because if its URL changes, the publisher need only update the metadata for the DOI to link to the new URL.

displacements per atom (DPA): Displacements per Atom (DPA) is a phenomenological measure of radiation damage. DPA = $0.8E_{\text{avail}}/(2E_d)$ where $E_{\text{avail}}$ is the total energy available in a reaction and $E_d$ is an element dependent parameter, usually equal to 25 eV, but can range as high as 90 eV.

distribution, energy: Probability density function that gives the probability of a particle scattering with an outgoing energy $E'$ as a function of incident energy $E$, $P(E'|E)$

distribution, energy-angle distributions: Probability density function that gives the probability of a particle scattering in direction $\mu = \cos(\theta)$ with an outgoing energy $E'$ as a function of incident energy $E$, $P(E',\mu|E)$

dose: The product of a radiation energy times the probability per disintegration, the resulting unit is MeV$x\times\text{Bq-s}$.

double $\beta$ decay: The simultaneous transformation of two neutrons inside a nucleus into two protons, or alternatively, the simultaneous transformation of two protons into two neutrons.

EADL: Evaluated Atomic Data Library, a library developed at LLNL to store atomic relaxation data. Atomic relaxation data describes the emission of electrons and photons as an atom relaxes back to neutrality following an ionization event. EADL is stored in the ENDL format and the contents of EADL have been translated into the ENDF format and are equivalent to the ENDF atomic relaxation sublibrary.

EAF: The European Activation File (EAF) is the collection of nuclear data that is required to carry out inventory calculations of materials that have been activated following exposure to neutrons or charged particles.

EDA: R-matrix fitting code developed at LANL for fitting reactions on light nuclei.

EEDL: Evaluated Electron Data Library, a library that describes the interaction of electrons with matter. EEDL is stored in the ENDL format and the contents of EEDL have been translated into the ENDF format and are equivalent to the ENDF electron sublibrary.

EFF: The European Fusion File (EFF) Project is a collaborative project with work funded by the European Fusion Development Agreement (EFDA). The emphasis is on the pooling of resources and removal of duplication of effort, leading to the efficient development of two types of nuclear data libraries for use in fusion power plant design and operation studies. The two branches consist of, on the one hand, a transport file for modeling and design capabilities and, secondly, an activation file for the calculation and simulation of dose rates and energy release during operation of a future power plant.
elastic scattering: A scattering event in which the scattering particle is reemitted with an outgoing energy such that the scattering particles’s incident and outgoing energies are equal in the system center of mass frame.

elastic scattering, compound: Elastic scattering in which the scattering particle is absorbed into the target nucleus, forming a compound nucleus which later decays, (re)emitting the scattering particle. Compound elastic scattering angular distributions are typically nearly isotropic in the rest frame of the colliding system.

elastic scattering, shape: Elastic scattering in which the scattering particle scatters off the target nucleus without being absorbed. Shape elastic scattering angular distributions are typically anisotropic in the rest frame of the colliding system.

elastic scattering, TSL: Elastic scattering in which a low energy/thermal neutron scatters either coherently or incoherently off many nuclei in a material.

electron: A negatively charged fundamental particle (lepton). It has a mass of $0.5109989(4)$ MeV, a charge of $-1.60217646(6) \times 10^{-19}$ Coulombs and $J^P = 1/2^+$. Electrons bind to atomic nuclei, occupying specific atomic levels.

electron capture (EC): The process where one of the protons inside a nucleus following the interaction with one of the orbiting electrons transforms into a neutron plus a neutrino: $p + e^- \rightarrow n + \nu_e$.

element, chemical: Materials with identical chemical properties. There is a one-to-one correspondence between the element name and the number of protons. For instance, all carbon atoms have 6 protons in their nuclei.

element, XML: Everything from (including) the element’s start tag to (including) the element’s end tag. An element can contain other elements, text, attributes, or a mix of these.

ENDF/A: Formerly the development library of the ENDF project. Historically partial and in development evaluations were maintained in the ENDF/A library.

ENDF/B: The main release library of the ENDF project. The ENDF/B library is a product of the Cross Section Evaluation Working Group (CSEWG), a long standing collaboration of institutions within and without the United States.

ENDF format: A nuclear data format used by many nuclear data projects including the namesake library, the ENDF/B library. The ENDF format is maintained by the Cross Section Evaluation Working Group (CSEWG).

ENDL: LLNL’s Evaluated Nuclear Data Library, an evaluated nuclear reaction data library comparable in scope with the ENDF library. ENDL is used to support transport calculations at LLNL.

ENDL format: The format of data in ENDL. It is much simpler than the ENDF file, but is much less expressive than the ENDF format as all data must be stored as a multi-dimensional linear interpolation table.

end point energy: The maximum kinetic energy that an electron in $\beta^-$ decay or a positron in $\beta^+$ decay can have, obtained when the kinetic energy of the neutrino/anti-neutrino is equal to zero.

energy group:FIXME

energy loss: See stopping power.

entity–relationship model: An entity–relationship model (ER model) is a data model for describing the data or information aspects of a business domain or its process requirements, in an abstract way that lends itself to ultimately being implemented in a database such as a relational database.

ENSDF: The Evaluated Nuclear Structure Data File. ENSDF contains evaluated nuclear structure and decay data in a standard format. For each nuclide, all known experimental data used to deduce nuclear structure information are included. Each type of experiment is presented as a separate dataset. In addition, there is a dataset of “adopted” level and gamma-ray transition properties, which represent the evaluator’s determination of the best values for these properties, based on all available experimental data.

ENSDF format: The format used to store the ENSDF database.

EPDL: Evaluated Photon Data Library, a library developed at LLNL to store photo-atomic data. EPDL is stored in the ENDL format and the contents of EPDL have been translated into the ENDF format and are equivalent to the ENDF photo-atomic sub-library.

equiprobable bins:FIXME
Ericson fluctuations: Ericson fluctuations are the statistical fluctuations observed in cross sections in the overlapping resonance region (ORR).

evaluate: The process of digesting experimental data, combining it with the predictions of nuclear model calculations and attempting to extract the true value of a cross section is referred to as an evaluation.

evaluation: One projectile and one target material and all the data needed to describe the interactions between the two over some incident energy range. The projectile is usually assumed to be impingement upon a stationary target and is usually assumed to be less massive than the target material. However, the format must be flexible enough to store data in “inverse kinematics” where the lighter particle is at rest relative to the heavier one or in the center of mass frame. The format must also be flexible enough that the target material is actually a nontrivial collection of nuclei such as in thermal neutron scattering.

evaluator: One who evaluates.

EXDL: The Evaluated Excitation Data Library, a library that describes the excitation of atoms due to photon interaction. ENDF does not yet have a format to support this data.

eXtensible Markup Language (XML): A markup language that defines a set of rules for encoding documents in a format which is both human-readable and machine-readable.

flux weighting: [FIXME]

fission: Nuclear fission is either a nuclear reaction or a radioactive decay process in which the nucleus of an atom splits into smaller parts (lighter nuclei). The fission process often produces free neutrons and photons (in the form of gamma rays), and releases a very large amount of energy even by the energetic standards of radioactive decay.

fission, spontaneous: A radioactive decay process which results in fission.

fission product yield (FPY): The yield of a particular isotope some time after a fission event. As there are two fission fragments following every (binary) fission event, the yield sums to two.

fission product yield, independent (IFPY): The yield of a particular fission fragment immediately after a fission event. As there are two fission fragments following every (binary) fission event, the yield sums to two.

Fudge: An open source processing and data modification code developed at LLNL. Fudge is the first processing code capable of manipulating data in both the ENDF and GND format. Fudge originally stood for “For UpDating and Generating Evaluations”.

gamma ray: The term “gamma rays” is used here for electromagnetic waves (photons) that have a nuclear origin, that is, photons are emitted following the rearrangement of the protons and neutrons in a nucleus. In contrast, the term “X rays” is used for photons emitted following the rearrangement of the electrons orbiting an atomic nucleus. Gamma rays are one type of radiation.

gamma ray emission: Atomic nuclei are quantum system with a discrete set of energies. Gamma ray emission can take place when a nucleus rearranges its protons and neutrons into a lower energy state.

GND: Generalized Nuclear Data format. The name for the hierarchical arrangement of nuclear data developed at LLNL. It was originally developed in XML, but can be serialized into any hierarchical low level data format such as HDF5. GND is the prototype of the format being developed by WPEC SubGroup 38.

group: See “energy group”.

half-life ($T_{1/2}$): The length of time for a given radioactive species to reduce its activity in half. The number of decays as a function of time is given as:

$$N(T) = N_0 \exp \left( -\ln(2) \times (T - T_0)/T_{1/2} \right)$$  \hspace{1cm} (C1)

$T_{1/2}$ is related to the life-time $\tau$ by:

$$T_{1/2} = \ln(2) \times \tau$$  \hspace{1cm} (C2)

and to the width $\Gamma$ by

$$T_{1/2} = \ln(2) \times \left( h/2\pi \right)/\Gamma$$  \hspace{1cm} (C3)

where $h$ is Planck’s constant. Please note that when the half-life of a given nuclear level is smaller than $10^{-15}$ seconds, it is customary to list the width ($\Gamma$) of the level instead.

HDF5: Hierarchical Data Format, version 5 (HDF5) is a data model, library, and file format for storing and managing data. It supports an unlimited variety of datatypes, and is designed for flexible and efficient I/O and for high volume and complex data.
hindrance factor (HF): The ratio between the alpha decay partial experimental half-life and a calculated half-life.

incoherence, quantum mechanical: Two quantum mechanical wave functions are incoherent if they either do not have a constant phase difference or the same frequency (and therefore energy). Quantum mechanical wavefunctions (or parts of the same wavefunction) are said to be incoherent if their phases are not equal.

intensity, radiation: Radiation intensities indicate the probability of observing the corresponding radiations. Two different conventions are used:
- For decay radiation, intensities are listed per 100 decays of the parent nucleus. For instance, in the decay of $^{232}$Th, the alpha particle with 4012 keV is listed as having an intensity of 78.2%, which means that this alpha particle will be emitted 78.2 times for every 100 decays of 232Th.
- For the gamma rays, intensities correspond to gamma branching ratios for each level, assigning 100 to the strongest gamma ray.

IBAN DL: The Ion Beam Analysis Nuclear Data Library developed and formerly maintained by A. Gurbich under the IAEA auspices. It contains available experimental nuclear cross-sections relevant to Ion Beam Analysis and is available at https://www-nds.iaea.org/exfor/ibandl.htm

inelastic scattering: A scattering event in which the scattering particle is reemitted with an outgoing energy different from the scattering particles’s incident energies in the system center of mass frame. Usually the scattering particle’s energy is lost in the transition, but if the target is already in an excited state it is possible for the scattering particle to gain energy from the de-exciting target nucleus. This is occasionally referred to as “super-elastic scattering.”

inelastic TSL scattering: Inelastic scattering in which a low energy/thermal neutron scatters either coherently or incoherently off many nuclei in a material.

internal conversion: A process where the transition from one nuclear level to another level in the same nucleus is carried out by transferring the excess energy to an orbiting electron. The alternate process is gamma emission. The electrons are ejected from the atom with an energy equal to the transition energy minus the electron binding energy. These electrons are called conversion electrons and are labeled by the orbit the electron had occupied. For instance, CE-L means conversion electron from the L shell.

internal conversion coefficient (ICC): The ratio of the number of internal conversion decays to the number of gamma decays is the internal conversion coefficient, denoted $\alpha$. The ICC can be decomposed into the ICC for individual transitions (electric vs. magnetic and by multipolarity). The probability of internal conversion $P_{IC}$ is related to the probability of gamma emission $P_{\gamma}$ by $P_{IC} = \alpha P_{\gamma}$.

inverse kinematics: A two-body collision in the lab frame in which the projectile is the more massive of the two bodies and the target is the less massive of the two bodies.

isobar: A number of nuclei with the same number of protons plus neutrons are called isobars, such as $^{144}$Sm, $^{144}$Nd and $^{144}$Gd.

isomer: A nuclear isomer is a metastable excited state of an atomic nucleus.

isomeric Transition (IT): The process where a long-lived excited nuclear level decays by gamma emission or internal conversion.

isotone: A number of nuclei with the same number of neutrons are called isotones, such as $^{144}$Sm, $^{142}$Nd and $^{146}$Gd.

isotope: A number of nuclei with the same number of protons are called isotopes, such as $^{144}$Sm, $^{142}$Sm and $^{146}$Sm.

JEFF: Joint Evaluated File, predecessor of the JEFF project.

JEFF: The JointEvaluated Fission and Fusion File (JEFF) project is a collaboration between NEA Data Bank member countries. The JEFF library combines the efforts of the JEFF and EFF/EAF Working Groups to produce a common sets of evaluated nuclear data, mainly for fission and fusion applications. Available at https://www.oecd-nea.org/dbdata/jeff/

JENDL: The Japanese Evaluated Nuclear Data Library, developed by the Japan Atomic Energy Agency (JAEA) to support nuclear application development in Japan. Available at http://wwwndc.jaea.go.jp/jendl/j40/j40.html

$J^P$: The angular momentum ($J$) and parity ($P$) associated with a nuclear level or a particle. For instance, the ground states of nuclei with even number of protons and neutrons have $J^P = 0^+$. The intrinsic $J^P$ of the proton is equal to $1/2^+$. 
KERMA: An acronym that means “Kinetic Energy Released in MAterials” and is defined as the sum of the initial kinetic energies of all the charged particles liberated by uncharged ionizing radiation (i.e., indirectly ionizing radiation such as photons and neutrons) in a sample of matter, divided by the mass of the sample. It is defined by the quotient \( K = \frac{dE_{ic}}{dm} \).

Klein-Nishina formula: The Klein–Nishina formula gives the differential cross section of photons scattered from a single free electron in lowest order of quantum electrodynamics. At low frequencies (e.g., visible light) this is referred to as Thomson scattering; at higher frequencies (e.g., x-rays and gamma-rays) this is referred to as Compton scattering.

large angle Coulomb scattering (LACS): Charged particle scattering at large angles (typically \( \geq 20^\circ \)). At low angles or at low energies, charge particle scattering must be treated with an alternate approach.

lethargy: Neutron lethargy, or logarithmic energy decrement, \( u \), is a dimensionless logarithm of the ratio of the energy of source neutrons to the energy of neutrons after a collision:

\[
 u = \ln(E_o/E), \text{ or } u_2 - u_1 = \ln(E_1/E_2).
\]

Also, the feeling one has after working on a requirements document for far too long.

level: A quantum mechanical system or particle that is bound can only take on certain discrete values of energy.

logft: For each decay branch in \( \beta \) decay, the decay probability, \( T_{1/2} \) and energy released can be combined in a single quantity known as logft, which is defined as:

\[
 \text{logft} = \log(f(Z,E_0)T_{1/2})
\]

where \( E_0 \) is the end-point energy for the transition and \( f(Z,E_0) \) is the Fermi integral. Logft values increase with increasing \( T_{1/2} \), decay probability and \( E_0 \) values. There is a correlation between the type of transition and its logft value.

MAT: The designator used in the ENDF format to distinguish materials.

material: In nuclear data, a material generalizes the concept of a target and may include a single nucleus or nucleon or a macroscopic collection of nuclei that a projectile scatters off of.

mass: Mass is a property of a material which determines resistance to being accelerated by a force and the strength of its mutual gravitational attraction with other bodies.

mass, nuclear: The mass of a nucleus with \( Z \) proton and \( N \) neutrons in a neutral-atom state is:

\[
\begin{align*}
\text{Mass}(Z,N) &= Z \times \text{Mass}(\text{hydrogen atom}) + N \times \text{Mass}(\text{free neutron}) - \text{BE}(Z,N) \\
\end{align*}
\]

where \( \text{BE}(Z,N) \) is the Binding Energy, that is, the energy needed to dissociate the nucleus into free nucleons. Note in this product, masses are given in energy units.

mass excess (\( \Delta \)): The mass excess \( \Delta(Z,N) \) is defined as:

\[
\Delta(Z,N) = (\text{Mass}(Z,N)(\text{amu}) - A) \times \text{amu}
\]

where \( \text{Mass}(Z,N)(\text{amu}) \) is mass in atomic mass units and amu = Mass\((6,6)\)/12.

MF: In the ENDF format, the MF designator denotes the observable of interest. For example, MF=3 denotes cross section data.

Mott scattering: The modification of Rutherford scattering of an electron and a nucleus to include the effects of nuclear recoil and the nuclear magnetic moment.

MT: In the ENDF format, the MT designator denotes the reaction of interest. For example, MT=18 denotes fission.

multi-band treatment: The modification of Rutherford scattering of an electron and a nucleus to include the effects of nuclear recoil and the nuclear magnetic moment.

multifragmentation: Nuclear (multi-)fragmentation, defined as the nuclear decay mechanism in which at least three intermediate mass fragments (\( Z \geq 3 \)) are produced, is the disassembly phenomenon specific to hot nuclear matter produced in nuclear collisions at beam energies of 20-100 MeV/nucleon.

multiplicity: The average number of particles emitted from a reaction per unit incident energy. For an \( (n,2n) \), the multiplicity of neutrons is 2 for all energies, but the photon multiplicity is variable and depends on the nuclear structure of the residual nucleus.

multipolarity: A measure of the angular momentum carried away by a photon during an electromagnetic transition or decay. An \( \ell = 1 \) transition is a dipole transition, \( \ell = 2 \) transition is a quadrupole transition, etc.
NJOY: LANL’s nuclear data processing code, see http://t2.lanl.gov/nis/codes/NJOY12/

nubar, $\bar{\nu}$: The average number of neutrons emitted during a fission reaction.

nubar, delayed, $\bar{\nu}_d$: The average number of neutrons emitted from the (usually beta) decay of the fission fragments. In the ENDF format, the delayed nubar is broken out into six time groups.

nubar, prompt, $\bar{\nu}_p$: The average number of neutrons emitted immediately after fission, due either to emission before or during the fission event or evaporated from the fission products immediately after fission.

nucleus: The very dense region consisting of protons and neutrons at the center of an atom.

pair production: Pair production is the creation of an elementary particle and its antiparticle, for example an electron and its antiparticle, the positron.

particle: A particle is any small, localized object that can be attributed properties such as mass, charge, spin, parity, and half-life. This definition is deliberately broad to include fundamental particles such as electrons or photons, composite particles such as atomic nuclei or atoms, and also excited states of composite particles. We recognize that this definition may be surprising to nuclear physicists who tend to think of nuclei as a separate family from fundamental particles.

photoexcitation: Photoexcitation is the physical process in which an electron in an atom or molecule is promoted to an excited state from the interaction of a photon.

photoionization: Photoionization is the physical process in which an ion is formed from the interaction of a photon with an atom or molecule.

pointwise: A function represented by an interpolation table (usually with lin-lin interpolation) is a pointwise function.

primary knock-on atom (PKA): Primary Knock-on Atom

principal component analysis (PCA): A statistical procedure that uses an orthogonal transformation to convert a set of observations of possibly correlated variables into a set of values of linearly uncorrelated variables called principal components. For a multivariate Gaussian probability distribution, the orthogonal transform is constructed from the matrix of eigenvectors of the covariance matrix and the eigenvalues of the covariance matrix are the uncorrelated variables known as the principal components.

probability density function (PDF): In probability theory, a probability density function (PDF), or density of a continuous random variable, is a function that describes the relative likelihood for this random variable to take on a given value.

probability density function, Gaussian: The Gaussian PDF is given by:

$$f(x, \mu, \sigma) = \frac{1}{\sigma \sqrt{2\pi}} \exp \left(-\frac{(x - \mu)^2}{2\sigma^2}\right)$$  \hspace{1cm} (C8)

The parameter $\mu$ in this definition is the mean or expectation of the distribution (and also its median and mode). The parameter $\sigma$ is its standard deviation; its variance is therefore $\sigma^2$. A random variable with a Gaussian distribution is said to be normally distributed and is called a normal deviate.

probability density function, log-normal: A log-normal (or lognormal) distribution is a continuous probability distribution of a random variable whose logarithm is normally distributed. Thus, if the random variable $x$ is log-normally distributed, then $y = \ln(x)$ has a normal distribution. Likewise, if $y$ has a normal distribution, then $x = \exp(y)$ has a log-normal distribution. A random variable which is log-normally distributed takes only positive real values.

probability density function, normal: See probability density function, Gaussian.

projectile: In a two-body collision in the laboratory frame, the projectile is the body in motion. The projectile is typically the less massive of the two colliding particles.

prompt fission gamma spectrum (PFGS): The spectrum of gammas emitted during and immediately after a fission event, but before the fission fragments undergo weak decays.

prompt fission neutron spectrum (PFNS): The spectrum of neutrons emitted during and immediately after a fission event, but before the fission fragments undergo weak decays.

PURM: The module in the Oak Ridge National Laboratory processing code AMPX (Dunn, 2002) responsible for producing cross section probability tables in the URR.

PURR: The module in the Los Alamos National Laboratory processing code NJOY (MacFarlane, 2012) responsible for producing cross section probability tables in the URR.
**Q matrix**: The matrix that connects the Cumulative Fission Yields with the Independent Fission Yields.

**quantum number**: Quantum numbers describe values of conserved quantities in the dynamics of a quantum system. Typical quantum numbers encountered in nuclear physics are spin, parity, orbital and total angular momenta, mass, charge and isospin.

**Q value**: The Q value for a reaction is the amount of energy released by that reaction.

**Raleigh scattering**: Rayleigh scattering named after the British physicist Lord Rayleigh, is the (dominantly) elastic scattering of light or other electromagnetic radiation by particles much smaller than the wavelength of the radiation. When the scattering is coherent one uses Thompson scattering formula modified by (anomalous) form factors.

**reaction**: A nuclear reaction is a process in which two (or more) nuclei or nuclear particles collide, producing a different set of products than the initial particles. From this perspective, elastic scattering is not a “reaction” while inelastic scattering is a “reaction” as the energy of the inelastically scattered particle has changed.

**reaction, exclusive**: A reaction with well defined non-gamma reaction product multiplicities (e.g., \((n,2n)\)). Inelastic reactions to discrete states are considered exclusive since the residual nucleus is left in a well defined state before it gamma cascades.

**reaction, inclusive**: A reaction that is a sum of exclusive reactions (e.g., total or fission). Inelastic from the continuum is not considered here since the residual is still well defined but total inelastic is inclusive since it is a sum of inelastic continuum and discrete reactions.

**reaction, photonuclear**: A nuclear reaction induced by a high energy photon (typically with incident energy \(\geq 1\) MeV)

**Reference Input Parameter Library (RIPL)**: A common library of input parameters for use in nuclear reaction modeling codes. Much of the data in RIPL is derived from other sources (e.g. ENSDF) or from systematics. More information is available at [https://www-nds.iaea.org/RIPL-3/](https://www-nds.iaea.org/RIPL-3/).

**resonance**: In the context of nuclear data, a resonance is a compound state formed when a nucleus absorbs a projectile. The compound state is unstable against particle emission and so is characterized by emission widths.

**resonance region (RR)**: The region in a cross section where resonances are observed is called the resonance region (RR).

**resonance region, resolved (RRR)**: The resolved resonance region is the portion of the resonance region where the resonance widths are small enough and the resonances are spaced far enough apart in energy that they can be individually resolved experimentally.

**resonance region, overlapping (ORR)**: The overlapping resonance region is the portion of the resonance region at the highest energies where the resonance widths are larger or comparable to the inter-resonance spacing. Thus, the cross sections exhibit large interference effects and sizable fluctuations. These fluctuations are known as Ericson fluctuations.

**resonance region, unresolved (URR)**: In the unresolved resonance region, the resonances of a cross section are so close together that, although they are not overlapping, they cannot be experimentally resolved.

**R-matrix**: A tool in computational quantum mechanics for studying two-body scattering. R-matrix theory begins by placing a reaction zone inside a spherical box. Outside the box, asymptotic (and thus calculable) wave functions are used to describe the pair of scattering particles. Using the continuity of flux on the boundary, an R-matrix can be defined which can be used to parameterize things such as the scattering cross section. The R-matrix method was originally formulated for studying resonances in nuclear scattering by Wigner and Eisenbud.

**Rutherford scattering**: The scattering of two charged particles purely by the static electric force. The Rutherford cross section is calculable either quantum mechanically and classically.

\[ S_{\alpha\beta} \]: \( S_{\alpha\beta} \) is the scattering kernel of the double differential elastic scattering cross section for thermal neutrons in the Thermal Scattering Law formalism.

**SAMMY**: A resonance fitting code based on R-matrix theory developed by ORNL. It is capable of fitting resonances on nearly all nuclei and is capable of simultaneously fitting hundreds if not thousands of resonances simultaneously.

**self-shielding**: Self-shielding occurs when the neutron flux in one part of a material is shielded from another part of the same material.

**shell**: A nuclear or atomic shell is a state or a collection of states with the same well defined quantum numbers. Electronic shells are usually denoted with the
X-ray notation (e.g. K, L, M, N, ...) while nuclear shells are denoted using a different notation (e.g. s, p, d, f, ...).

**shell model**: The nuclear shell model is a model of the atomic nucleus which uses the Pauli exclusion principle to describe the structure of the nucleus in terms of energy levels.

**spallation**: The process in which a nucleus emits a large number of nucleons as a result of being hit by a high-energy particle, thus greatly reducing its atomic weight.

**spin group**: The name of the group of resonances with the same total spin and parity. Organizing resonances by spin groups results in much smaller tables of resonance parameters as the needless storing of zero widths is avoided.

**stopping power**: The retarding force acting on charged particles due to interaction with matter, resulting in loss of particle energy. The stopping power of the material is numerically equal to the loss of energy E per unit path length, x: \( S(E) = -dE/dx \)

**subshell**: Each shell in a shell model is composed of one or more subshells, which are themselves composed of atomic or nuclear orbitals.

**target**: In a two-body collision in the laboratory frame, the target is the body at rest. The target is typically the more massive of the two colliding particles.

**TENDL**: TALYS Evaluated Nuclear Data Library is an evaluated nuclear data library generated using the TALYS reaction code. TENDL is available at [http://www.talys.eu/tendl-2012/](http://www.talys.eu/tendl-2012/).

**thermal scattering law (TSL)**: Thermal scattering law data describe the situation in which the de Broglie wavelength of an incident neutron is so large that the neutron wave function cannot resolve individual nuclei, but rather sees the macroscopic material.

**Thompson scattering**: Thomson scattering is the elastic scattering of electromagnetic radiation by a free charged particle, as described by classical electromagnetism. It is just the low-energy limit of Compton scattering.

**time group**:FIXME

**transfer matrix**: The transfer matrix is a matrix that arises during the discretization of the Boltzmann transport equation. See subsection XV for a precise definition.

**UKNDL**: United Kingdom Nuclear Data Library, the nuclear data library and format developed by UK’s Atomic Weapons Establishment. UKNDL format was the prototype of both the ENDF and ENDL formats.

**Universal Modeling Language (UML)**: The Unified Modeling Language is a general-purpose modeling language in the field of software engineering, which is designed to provide a standard way to visualize the design of a system.

**Universal Resource Locator (URL)**: A URL is one type of Uniform Resource Identifier (URI); the generic term for all types of names and addresses that refer to objects on the World Wide Web. The term “Web address” is a synonym for a URL that uses the HTTP / HTTPS protocol.

**uncertainty**: The uncertainty is the square root of expected deviation of a measurement from the expected mean.

**vacancy, shell**: A (sub)shell vacancy occurs when a particle is knocked out of a nuclear or atomic shell, leaving a “hole” or unoccupied state in the shell.

**Working Party on Evaluation Cooperation (WPEC)**: A framework provided by the Nuclear Energy Agency so that other institutions can co-operate and conduct multi-year projects which promote the exchange of information on nuclear data evaluations, measurements, nuclear model calculations, validation, and related topics.

**xlink**: XML Linking Language, or XLink, is an XML markup language and W3C specification that provides methods for creating internal and external links within XML documents, and associating metadata with those links.

**xpath**: XPath, the XML Path Language, is a query language for selecting nodes from an XML document. In addition, XPath may be used to compute values (e.g., strings, numbers, or Boolean values) from the content of an XML document.

**xsd**: XSD (XML Schema Definition), a recommendation of the World Wide Web Consortium (W3C), specifies how to formally describe the elements in an Extensible Markup Language (XML) document. It can be used by programmers to verify each piece of item content in a document.

**xsdir**: The xsdir file, serves as a table of contents for the transport code MCNP with information on where and how the data are stored for each target material can be found.
**xsl:** Extensible Stylesheet Language (XSL) is used to refer to a family of languages used to transform and render XML documents.

**ZA:** A common nucleus designator usually computed as $ZA = Z \times 1000 + A$. 
Appendix D: Contributors

Here we attempt to list everyone who contributed to this document and the general Subgroup 38 discussions. We apologize if we have omitted anyone. Below we list everyone in English alphabetical order by family name. Each contributor's affiliation is given in the form “institute, country”. Where an institute has multiple locations, the form of the affiliation is “institute, city, country”.

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