Self-consistent adiabatic description of the fission: automatic production of class-II PES

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P(ND)^2-2

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Microscopic description of the fission

Two-step approach :

1. Production of microscopic potential energy surfaces (PES)
   - Hartree-Fock-Bogoliubov code using a two-center oscillator basis
   - effective nucleon-nucleon interaction Gogny D1S
   - *N*-dimensional PESs
   - results : static properties of the fragments
2. Wave packet propagation
   - TDGCM method with GOA
   - initial state : eigenstate of an extrapolated first well
   - microscopic inertia tensor (GCM)
   - results : statistic properties of the fragments

Formalism - HFB2CT

$$\delta\langle\varphi|\hat{H} - \lambda_N\hat{N} - \lambda_Z\hat{Z} - \sum_i \lambda_i\hat{Q}_{i0}|\varphi\rangle = 0$$

Constrained Hartree-Fock-Bogoliubov method:

- D1S Gogny parametrization
- self-consistent mean and pairing fields
- two-center harmonic oscillator basis

Constraints:

- neutron and proton numbers N et Z

$$\begin{aligned}\langle\varphi|\hat{N}|\varphi\rangle &= N \\ \langle\varphi|\hat{Z}|\varphi\rangle &= Z\end{aligned}$$

- q_{10} to avoid spurious center of mass motion
- multipolar moments q_{i0}

$$\langle\varphi|\hat{Q}_{i0}|\varphi\rangle = q_{i0}$$

Formalism - TDGCM + GOA

General GCM state with N different degrees of freedom $\{q_1, \ldots, q_N\}$:

$$|\psi(t)\rangle \equiv \left(\prod_i^N \int dq_i\right) f(q_1, \ldots, q_N, t)|\phi(q_1, \ldots, q_N)\rangle$$

Variational principle:

$$\frac{\partial}{\partial f^*} \int_{t_1}^{t_2} \langle\psi(t)| \left(\hat{H} - i\hbar\frac{\partial}{\partial t}\right) |\psi(t)\rangle = 0$$

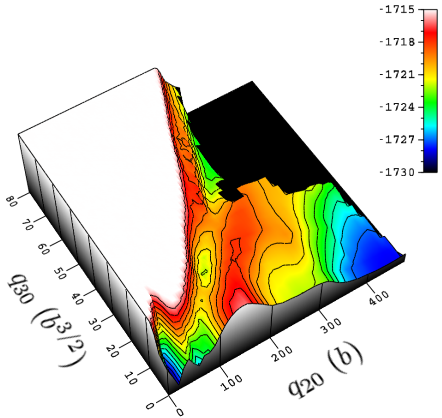
Using the Gaussian Overlap Approximation (GOA), we obtain a Schrödinger-like equation:

$$\hat{H}_{\mathrm{coll}} g(t) = i\hbar\frac{\partial}{\partial t} g(t)$$

with

$$\hat{H}_{\mathrm{coll}} = -\frac{\hbar^2}{2}\sum_{i,j}^{N} \frac{\partial}{\partial q_i} B^{ij} \frac{\partial}{\partial q_j} + \hat{V}$$

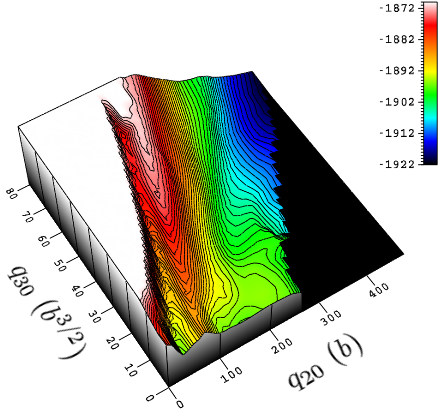
Results - PESs



$^{226}_{90}\mathrm{Th}_{136}$

N. Dubray et al., Phys. Rev. C **77**, 014310 (2008).

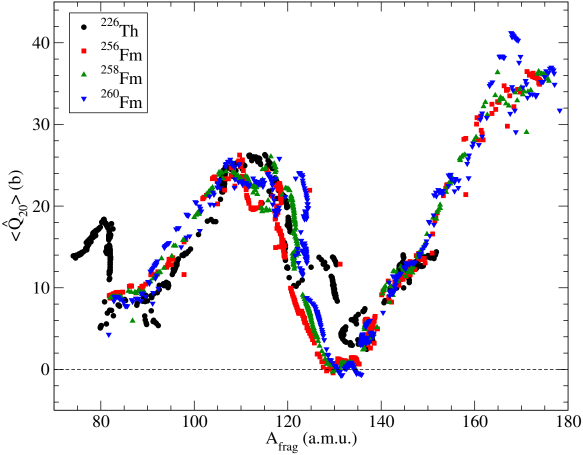
Results - PESs



$^{256}_{100}\mathrm{Fm}_{156}$

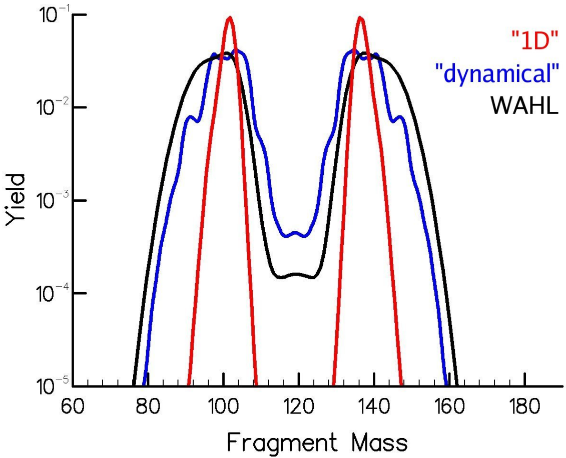
N. Dubray et al., Phys. Rev. C **77**, 014310 (2008).

Results - Fragment deformation $\langle \hat{Q}_{20} \rangle$



N. Dubray et al., Phys. Rev. C **77**, 014310 (2008).

Results - Fragment mass distribution for ^{238}U



H. Goutte et al., Phys. Rev. C **71**, 024316 (2005).

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Problem 1 - Convergence

- Problem: the HFB solver does not converge to a solution.
- Consequence: the HFB solution is bad.
- Symptom: the convergence quantity is too high.

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Problem 3 - Discontinuity

- Problem: two solutions close in the constraint deformation subspace are not close in the full deformation space.
- Consequence: a dynamic description using these points is missing a part of the physics (wrong barrier, saddle point, . . .).
- Symptom: none easily visible.

PES classes

Class	Convergence	Minimization	Discontinuity
0	problem ?	problem ?	problem ?
1	OK	problem ?	problem ?
2	OK	OK	problem ?
3	OK	OK	OK

The density-distance operator

We define the density-distance operator

$$D_{\rho\rho'}(|\psi\rangle, |\psi'\rangle) \equiv \int \mathrm{d}\tau^3 |\rho(\vec{r}) - \rho'(\vec{r})|$$

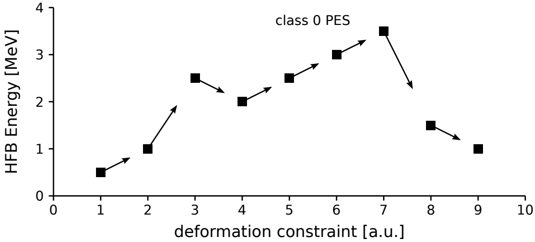
where $\rho(\vec{r})$ and $\rho'(\vec{r})$ are the total local densities of the states $|\psi\rangle$ and $|\psi'\rangle$.

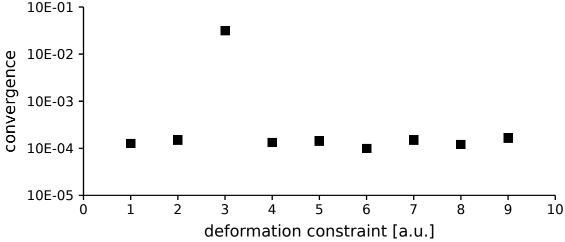
4-lines long algorithm to clean a PES

- all solutions with a too high convergence value are marked as "bad",
- all solutions with a too high maximum density distance value AND with their energy being higher than the corresponding partner's energy are marked as "bad",
- all solutions marked as "bad" are recalculated from the neighboring solution with the lowest energy that has not been used for the same calculation before,
- recalculate the density distances and restart.

- this algorithm can be used during or after the production of a N-dimensional PES.
- if there is no fatal convergence problem and if all valleys have been discovered, the result is at least a class 2 PES.

Example of a PES cleaning





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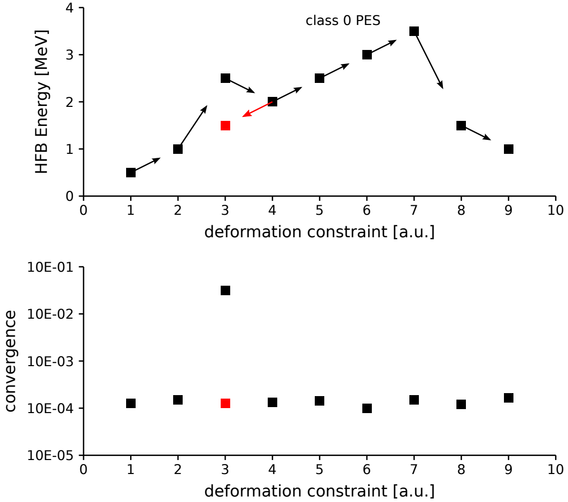

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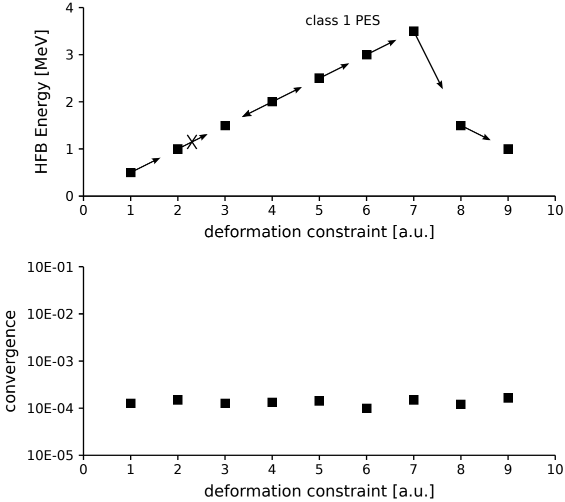

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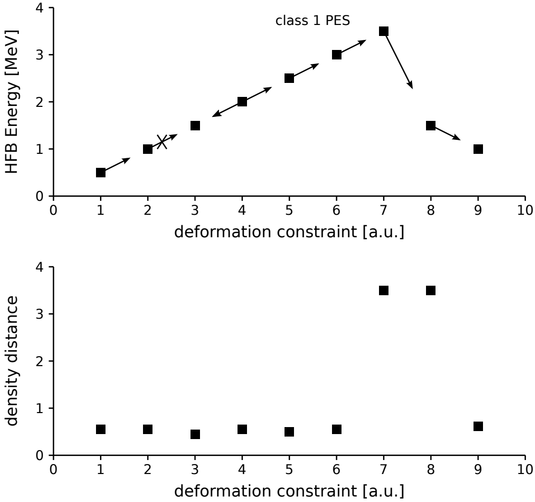

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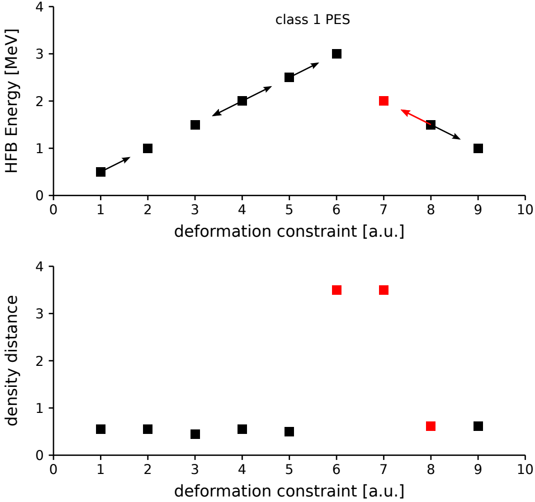

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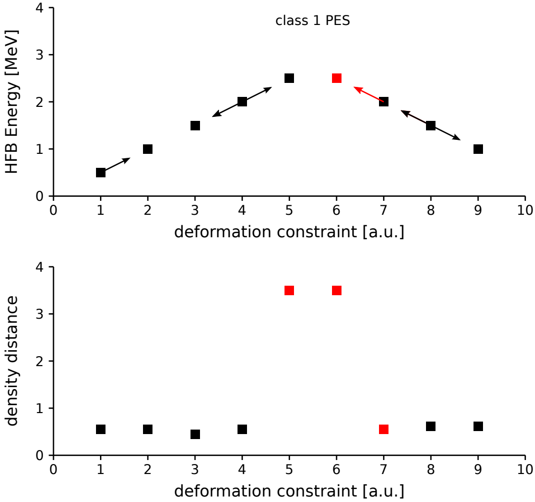

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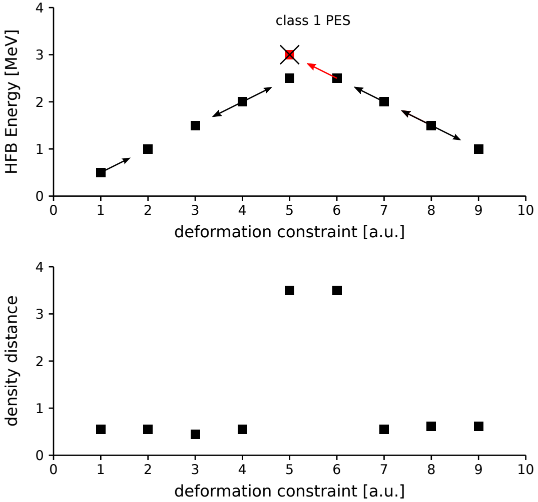

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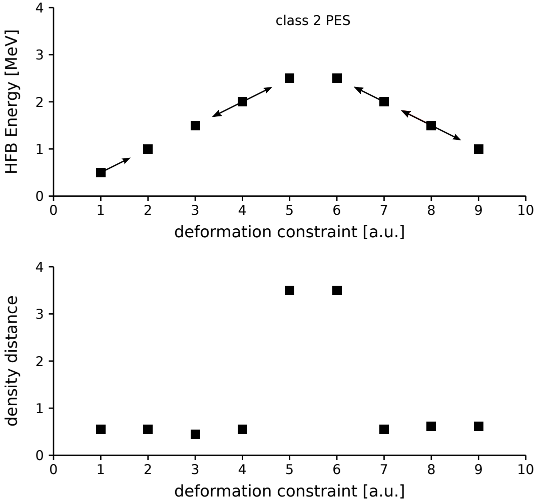

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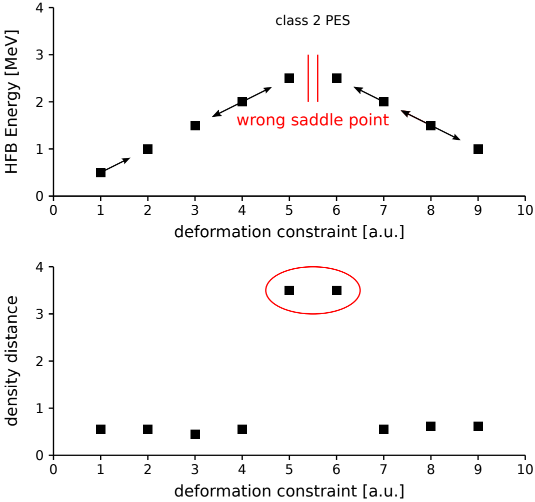


Important remark

We call a final point a solution that does not change when taking a bigger deformation subspace with the same symmetries.

If there has been no fatal convergence problem and if all valleys have been discovered,

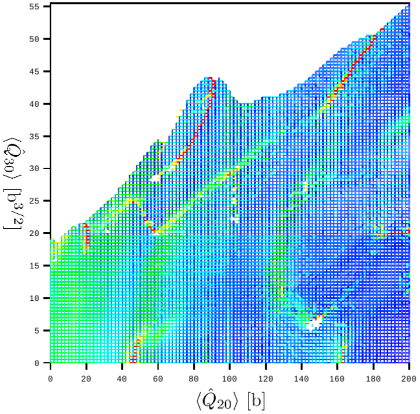
- a class 2 PES can have wrong or missing saddle points,
- a class 3 PES has only final points (minima, saddle points, etc. . .).

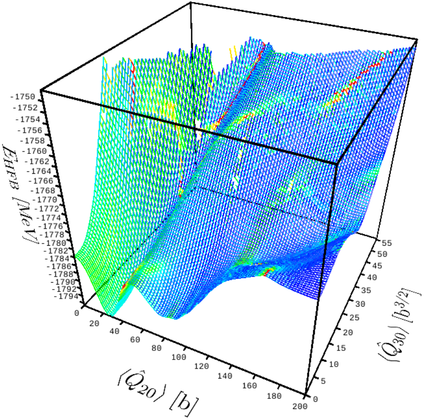
N. Dubray and D. Regnier, Comp. Phys. Comm. **183**, 2035 (2012).

Realistic example of automatic production

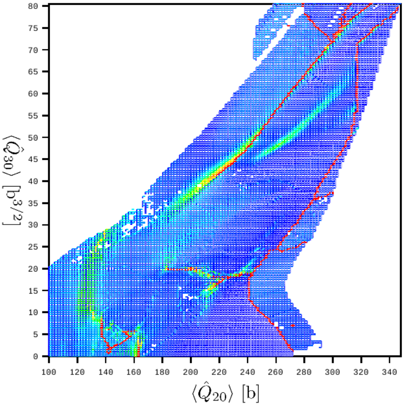
- 1- and 2-center PESs for ^{238}U.
- fully automatic production.
- class 2 PES enforced.
- basis-distance used in conjunction with density-distance.

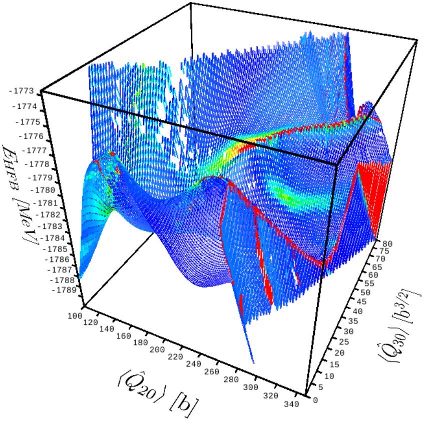
1-center PES of ^{238}U - class 2 / 3?





2-center PES of ^{238}U- class 2 / 3?





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Class 3 PESs

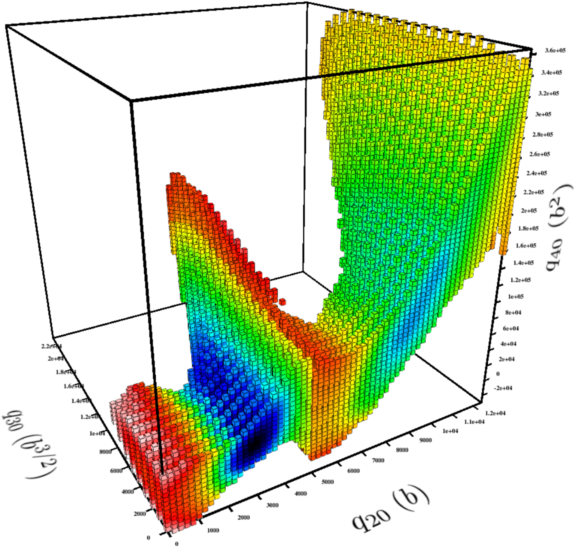
Continuous regular mesh ?

- dimension of a class 3 PES: $N = 1, 2, 3, 4, 5, 6, \ldots$?
- regular mesh + hypercube + $N > 2$ = huge number of points N_p
- dimension of the TDGCM + GOA hamiltonian matrix: N_p^2

Use a sparse mesh !

- no hypercube, focus on the physics in any dimension
- optimal number of points
- solve the TDGCM + GOA equation with FEM

Class 3 PES (q_{20}, q_{30}, q_{40}) for ^{240}Pu



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Conclusions

- Producing class 2 PESs is easy and can be automatic.
- Class 0,1,2 PESs can have non-final points.
- A class 3 N-dimensional PES has the same saddle points, paths... as any extended $(N + x)$-PES (since all points are final points).

Perspectives - fission description

In a near future, we plan to have

- a fully-automatic production of class 3 sparse N-PESs,
- a code to solve the TDGCM + GOA equation on sparse N-PESs with FEM (WIP by D. Regnier).

Next step: find the lowest N value for a given class 3 N-PES.