



Atomic scale Monte-Carlo simulations of neutron diffraction experiments on UO_2 up to 1664 K

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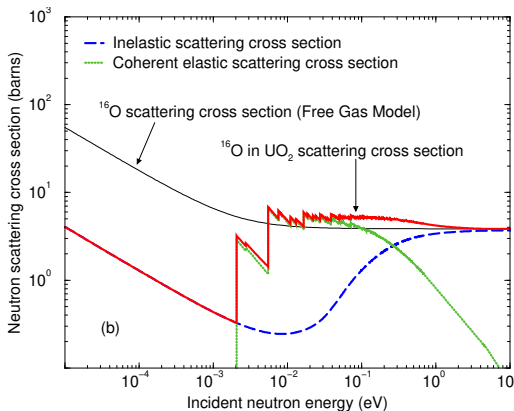
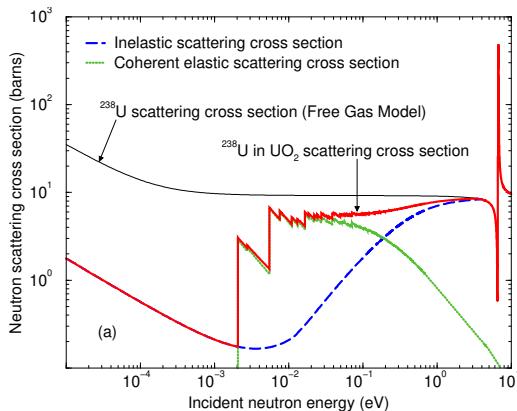
- 1 Context & thermal neutron scattering formalism
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$$\sigma_{t_u}(E) = \sigma_{\gamma_u}(E) + \sigma_{n_u}(E) + \sigma_{f_u}(E)$$

$$\sigma_{t_o}(E) = \sigma_{\gamma_o}(E) + \sigma_{n_o}(E) \quad (1)$$



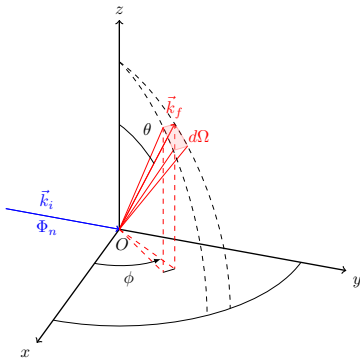


Figure: Schematic representation of low energy neutron scattering process.

At low neutron energy range (usually less than a few eV), the scattering of neutron in materials can be described via the double differential cross section (DDXS)

Cai and Kittelmann (2020); Schober (2014); Squires (2012)

$$\frac{d^2\sigma_{\vec{k}_i \Rightarrow \vec{k}_f}}{d\Omega dE_f} \equiv \frac{n(d\Omega, dE_f)}{\Phi_n d\Omega dE_f} = \frac{k_f}{k_i} s(\vec{Q}, \omega) \quad (2)$$

$s(\vec{Q}, \omega)$ is the scattering function defined by

$$s(\vec{Q}, \omega) \equiv \frac{1}{2\pi\hbar} \sum_{j,j'=1}^N \overline{b_{j'}} b_j \int_{-\infty}^{\infty} \langle j', j \rangle \exp(-i\omega t) dt. \quad (3)$$



Harmonic approximation:

$$\langle j', j \rangle = \exp(-i\vec{Q} \cdot (\vec{d}_{j'} - \vec{d}_j)) \exp(-W_{j'}(\vec{Q})) \exp(-W_j(\vec{Q})) \exp(\langle (\vec{Q} \cdot \vec{u}_{j'}(0))(\vec{Q} \cdot \vec{u}_j(t)) \rangle). \quad (4)$$

Incoherent approximation:

$$\langle (\vec{Q} \cdot \vec{u}_{j'}(0))(\vec{Q} \cdot \vec{u}_j(t)) \rangle = 0 \quad \text{if } j \neq j'. \quad (5)$$

Taylor expansion (phonon expansion [Sjolander \(1958\)](#)):

$$\exp(\langle (\vec{Q} \cdot \vec{u}_j(0))(\vec{Q} \cdot \vec{u}_j(t)) \rangle) = \sum_{n=0}^{\infty} \frac{1}{n!} (\langle (\vec{Q} \cdot \vec{u}_j(0))(\vec{Q} \cdot \vec{u}_j(t)) \rangle)^n. \quad (6)$$

Cubic approximation:

Phonon density of states (PDOS)

$$\langle (\vec{Q} \cdot \vec{u}_j(0))(\vec{Q} \cdot \vec{u}_j(t)) \rangle = \frac{\hbar Q^2}{2M_j} \int_0^\infty \frac{\rho_j(\omega)}{\omega} \left(\coth\left(\frac{\hbar\omega}{2k_B T}\right) \cos(\omega t) + i \sin(\omega t) \right) d\omega. \quad (7)$$

For uranium dioxide: $\sigma_{n_j}(E) = \sigma_{\text{coh}}^{\text{el}}(E) + \sigma_j^{\text{inel}}(E)$ (8)

where $\sigma_{\text{coh}}^{\text{el}}(E)$ represents the coherent elastic cross section (MF=7, MT=2 in ENDF-6 format [Trkov and Brown \(2018\)](#))

$$\sigma_{\text{coh}}^{\text{el}}(E) = \frac{\pi^2 \hbar^2}{m N V_{\text{unit cell}} E} \sum_{hkl}^{E \geq E_{hkl}} d_{hkl} |F(\vec{\tau}_{hkl})|^2 \mathcal{P}_{hkl}(\vec{PO}, \vec{\tau}_{hkl}) \quad (9)$$

where

$$F(\vec{\tau}_{hkl}) = \sum_{j=1}^{N_{\text{unit cell}}} \bar{b}_j \exp\left(-\frac{\hbar^2 \tau_{hkl}^2}{4 M_j k_B T} \Lambda_j(T)\right) \exp\left(i \vec{\tau}_{hkl} \cdot \vec{p}_j - i \delta_{jO} c_{123}^O \tau_{hkl}^3\right) \quad (10)$$

and the inelastic cross section $\sigma_j^{\text{inel}}(E)$ can be obtained by integrating its double differential form

$$\frac{d^2 \sigma_j^{\text{inel}}}{d\Omega dE_f} = \frac{\sigma_{b,j}}{4\pi k_B T} \sqrt{\frac{E_f}{E_i}} S_j^{\text{inel}}(\alpha_j, \beta) \quad (11)$$

where $S_j^{\text{inel}}(\alpha_j, \beta)$ is dimensionless thermal scattering laws (TSLs) (MF=7, MT=4).



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Codes or platforms capable of calculating the thermal scattering laws (TSLs) of crystalline materials

- ▶ LEAPR [MacFarlane \(1994\)](#) module of NJOY code [Macfarlane, Muir, Boicourt, Kahler III, and Conlin \(2017\)](#)
- ▶ Monte-Carlo based library NCrystal [Cai and Kittelmann \(2020\)](#); [Cai, Kittelmann, Klinkby, and Márquez Damián \(2019\)](#)
- ▶ OCLIMAX platform [Cheng, Daemen, Kolesnikov, and Ramirez-Cuesta \(2019\)](#); [Cheng and Ramirez-Cuesta \(2020\)](#)
- ▶ Full law analysis scattering system hub (FLASSH) [Zhu and Hawari \(2018\)](#)

Table: Comparison of codes or platforms which enable to calculate the TSLs of crystalline materials

Codes/Platforms	LEAPR+NCrystal	OCLIMACX	FLASSH	CINEL (this work)
Harmonic approximation	Yes	Yes	Yes	Yes
Incoherent approximation	Yes	1-p correct	1-p correct	Yes
Cubic approximation	Yes	No	No	Yes
Support any material	Yes	Yes	Yes	Yes
SCT approximation	Yes	No	No	No
GPU speedup	No	not reported	not reported	Yes



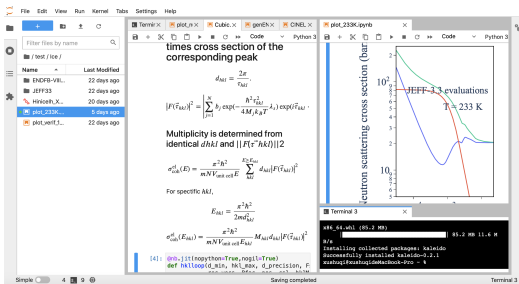


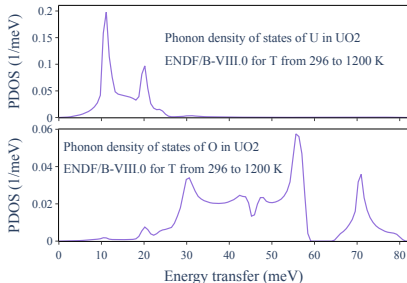
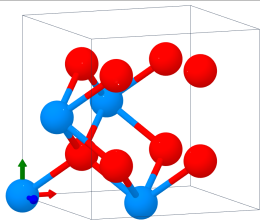
Table: Comparison of the computational time of the calculations of the TSL of **H in H₂O** at room temperature with phonon expansion order $N_{\text{phonon}} = 2000$. (Special thanks to our colleague P. TAMAGNO for providing his GPUs for test.)

	CPU	GPU K20	GPU K6000
GPU memory (gigabyte)	N/A	2	12
CUDA cores	N/A	384	2280
Computational time	~ 3 h	~ 15 min	< 2 min

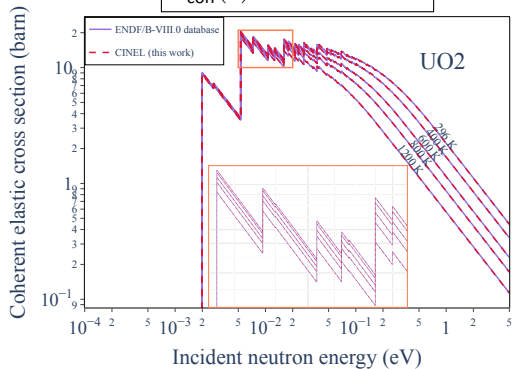
Figure: **JupyterLab** interface. Live code, texts, formatted mathematical equations and interactive graphics are mixed in Jupyter Notebooks which are integrated in JupyterLab together with blocks like terminal and text editor.

Ideal fluorite structure, $\text{Fm}\bar{3}\text{m}$ symmetry

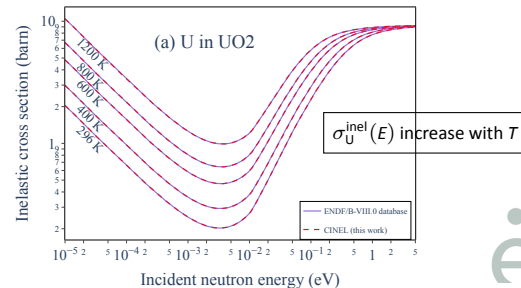
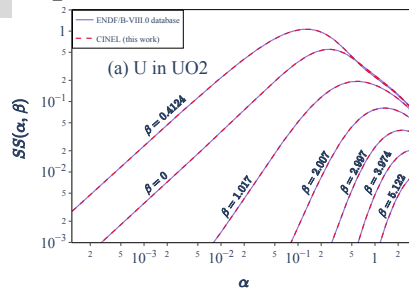
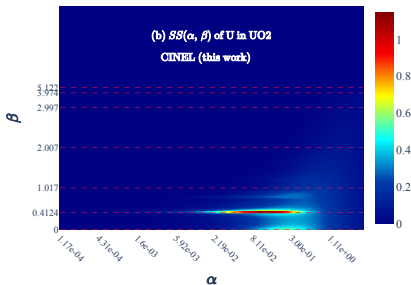
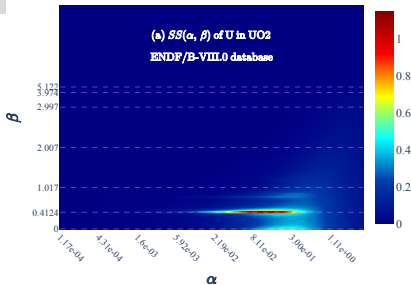
Ong et al. (2013)

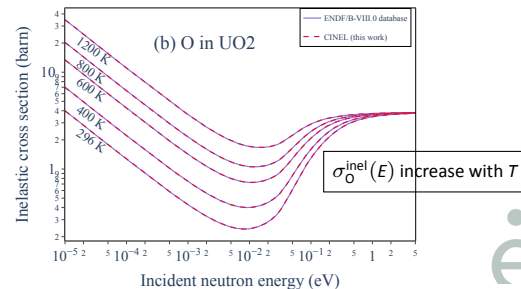
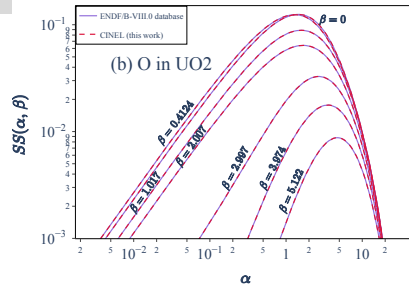
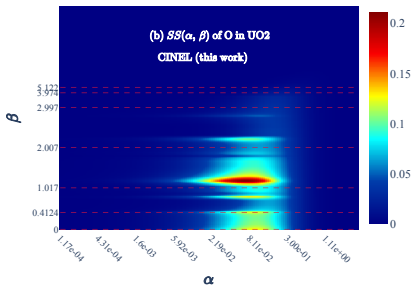
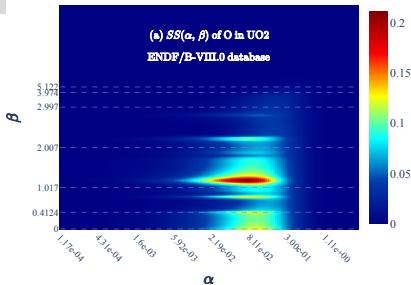


$\sigma_{\text{coh}}^{\text{el}}(E)$ decrease with T



Wormald, Fleming, Hawari, and Zerkle (2021)





Numerical validations of CINEL: ^{27}Al (FCC), ^9Be (HCP), ^{56}Fe (BCC)

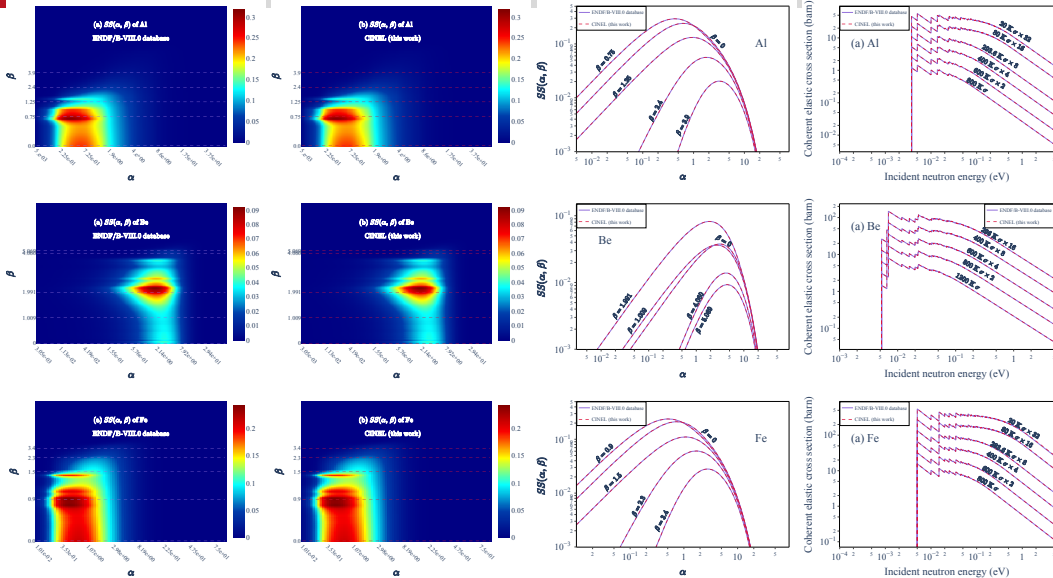


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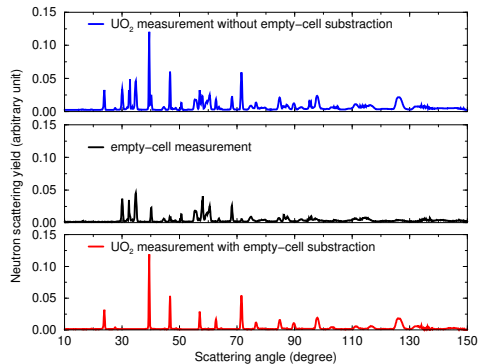
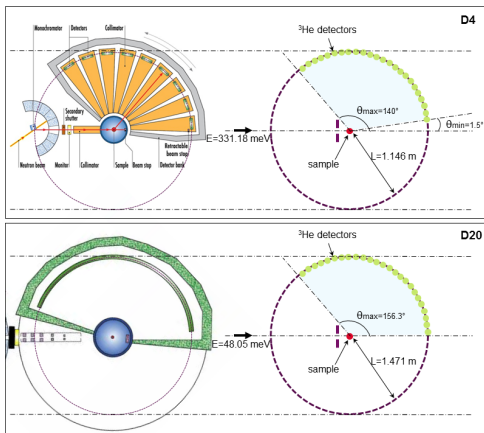
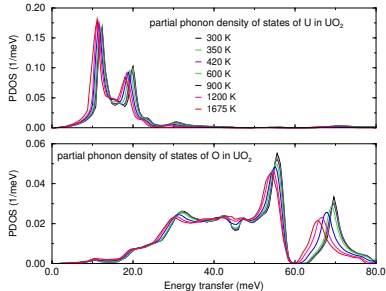
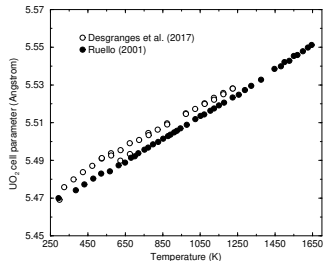
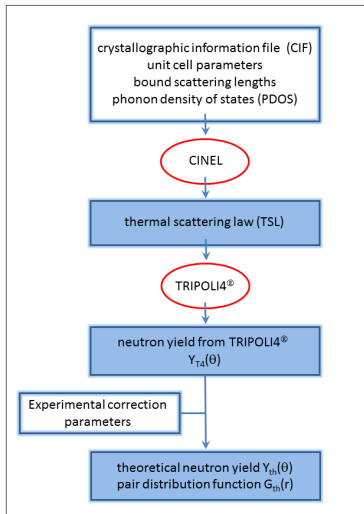


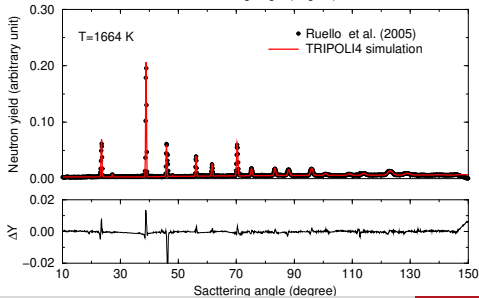
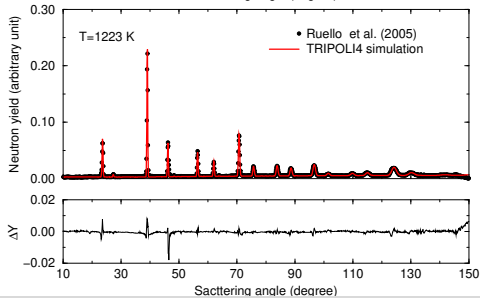
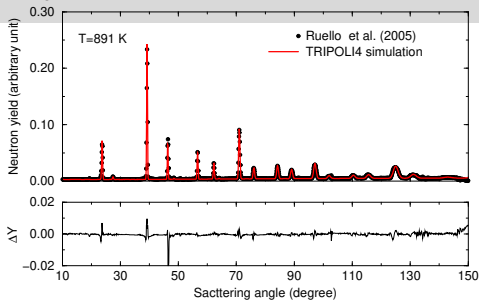
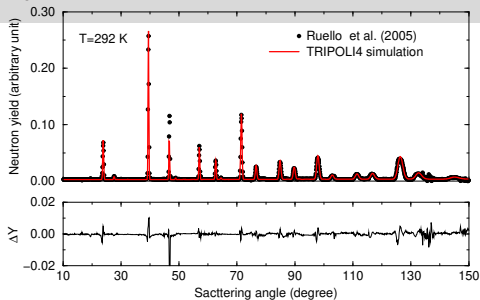
Figure: Examples of neutron diffraction patterns measured on D20 at $T = 292 \text{ K}$ ($E = 48.05 \text{ meV}$).

Figure: The left hand drawings represent simplified top view of the D4 and D20 diffractometers of the Institute Laue–Langevin (ILL). Those on the right hand side show the geometries introduced in the Monte-Carlo calculations.



Noguere et al. (2020)

Experimental & theoretical diffraction patterns of D2O



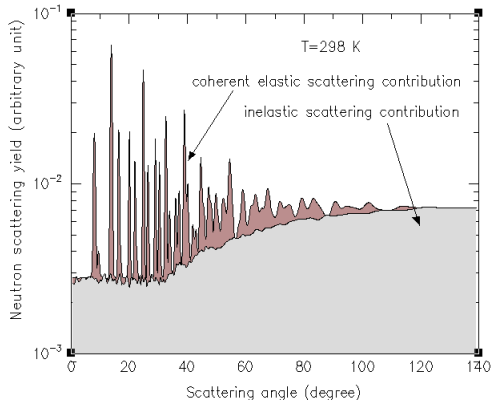


Figure: Inelastic and coherent elastic scattering contributions simulated with the TRIPOLI4[®] code at $T = 298$ K for D4 ($E = 331.18$ meV).

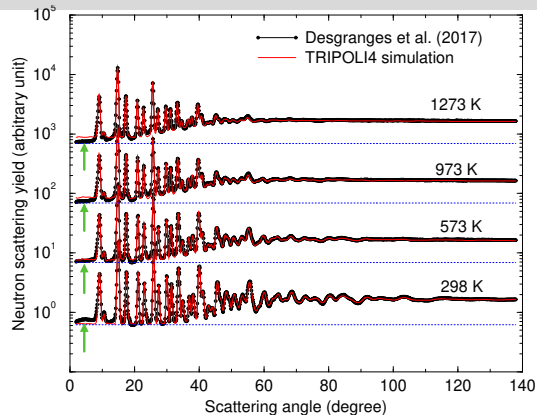
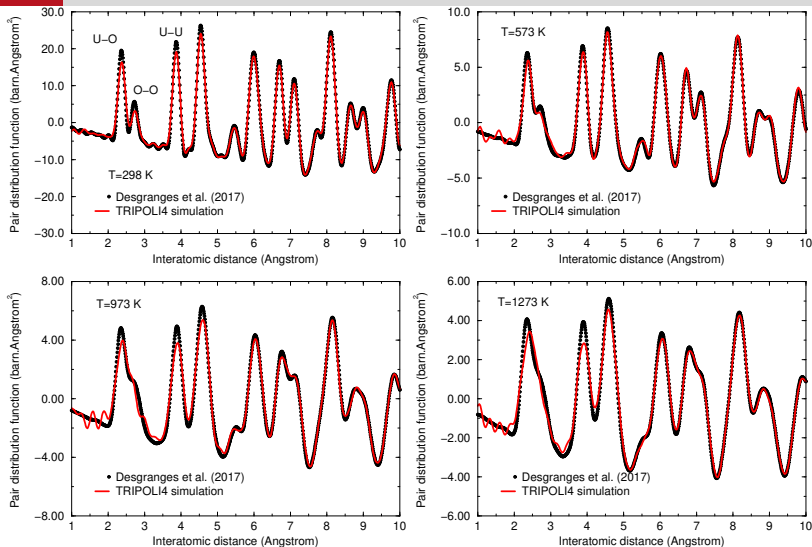


Figure: Experimental and theoretical UO_2 diffraction patterns for D4 from 298 to 1273 K ($E = 331.18$ meV). The green arrows highlight the differences observed between the experiment and the theory at the forward scattering angles. The dotted blue lines show the increasing inelastic scattering contribution with the scattering angle.



Confirmation of unexpected shortening of the U-O distance with increasing temperature reported in the literature.

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Conclusions

- ▶ Numerical validations of the thermal scattering law (TSL) processing code `CINEL` by comparing with the ENDF/B-VIII.0 database for various crystalline materials;
- ▶ Performances of the Monte-Carlo neutron transport code `TRIPOLI4`[®] with TSLs provided by `CINEL`, to simulate the neutron diffraction experiments on UO_2 measured at ILL up to 1664 K;
- ▶ Confirmation of the local deviation of the oxygen atoms from the average positions for UO_2 in ideal fluorite structure ($\text{Fm}\bar{3}\text{m}$ symmetry) at elevated temperatures by comparing the experimental and theoretical atomic pair distribution functions (PDF);

The neutron cross sections for UO_2 , calculated with the processing code `CINEL`, will be delivered at the JEFF project of the OECD/NEA databank.

Perspectives

- ▶ Introduction of one-phonon correction to improve the incoherent approximation and elimination of the cubic approximation in `CINEL`;
- ▶ Interpretation of the background correction to be investigated.



Thank you for your attention!

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