

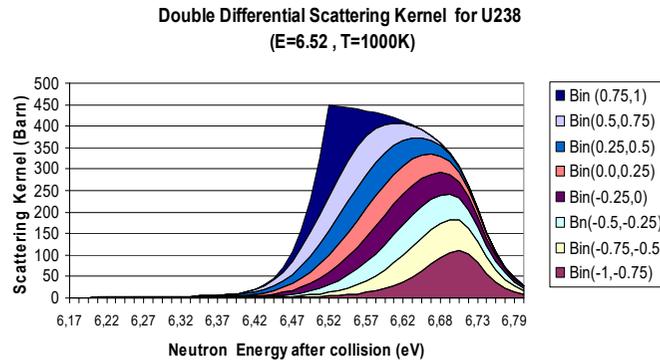
Observations of different thermal scattering models in view of graphite based materials

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The full free gas vs. the full solid state models (both energy dependent)



IDEAL GAS KERNEL: ENERGY DEPENDENT $\sigma_s(E_\tau)$

$$\sigma_s^{ii'}(E \rightarrow E', \Omega \rightarrow \Omega') =$$

$$\frac{1}{4\pi E} \sqrt{\frac{A+1}{A\pi}} \int_{\epsilon_{max}}^{\infty} d\xi \int_{\tau_0(\xi)}^{\tau_1(\xi)} d\tau$$

$$\left[\frac{(\xi + \tau)}{2} \right] \left(\frac{A+1}{A} \right)^2 \sigma_s \left[\left(\frac{A+1}{A^2} \right) \frac{[\xi + \tau]^2}{4} k_B T \right]$$

$$\exp \left\{ v^2 - \left[\frac{(\xi + \tau)^2}{4A} + \frac{(\xi - \tau)^2}{4} \right] \right\} \left[\frac{\epsilon_{max} \epsilon_{min} (\xi - \tau)^2}{B_0 \sin \phi} \right]$$

Solid state based expression of Word & Trammel (1980)

$$\frac{d^2 \sigma_R}{d\Omega dE} = \frac{v}{2\pi} \left(\frac{\Gamma_n}{2k_i} \right)^2 \left(\frac{E_f}{E_i} \right)^{1/2} \int_{-\infty}^{\infty} dT \exp(+iET) \int_0^{\infty} dt \exp[a(t)] \int_0^{\infty} dt'$$

$$\times \exp[a^*(t')] \times W(T, t, t').$$

$$W(T, t, t') = \left\langle \exp[-i\vec{k}_i \cdot \vec{r}(T-t)] \exp[i\vec{k}_f \cdot \vec{r}(T)] \right.$$

$$\left. \exp[-i\vec{k}_f \cdot \vec{r}] \exp[i\vec{k}_i \cdot \vec{r}(-t)] \right\rangle.$$

$$a(t) = -i(E_f - E_i - i\Gamma/2)t.$$

Solid state based phonon expansion scattering



Phonon Expansion: $S(\alpha, \beta) = e^{-\alpha\lambda} \sum_{n=0}^{\infty} \frac{1}{n!} [\alpha\lambda]^n \mathcal{T}_n(\beta)$

$$\# \mathcal{T}_0(\beta) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i\beta t} dt = \delta(\beta);$$

$$\# \mathcal{T}_1(\beta) = \frac{P(\beta)e^{-\frac{\beta}{2}}}{\lambda};$$

$$\# \mathcal{T}_n(\beta) = \int_{-\infty}^{\infty} \mathcal{T}_1(\beta') \mathcal{T}_{n-1}(\beta - \beta') d\beta';$$

$$\mathcal{T}_n(\beta) = e^{-\beta} \mathcal{T}_n(-\beta);$$

$$\# \beta = \frac{E - E'}{kT}; \quad \alpha = \frac{E' + E - 2(E E' \mu)^{\frac{1}{2}}}{AkT};$$

$$\# P(\beta) = \frac{\rho(\beta)}{2\beta \sinh\left(\frac{\beta}{2}\right)}; \quad \lambda = \int_{-\infty}^{\infty} P(\beta) e^{-\frac{\beta}{2}} d\beta;$$

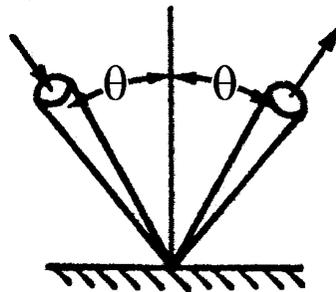
$$\# \rho(k, k', +t) = \frac{k \cdot k'}{2M} \int_0^{\infty} \frac{\rho(\omega)}{\omega} \left[\coth\left(\frac{\omega}{2kT}\right) \cos(\omega t) - i \sin(\omega t) \right] d\omega$$

■ *How important is the resonance treatment in low energies?*

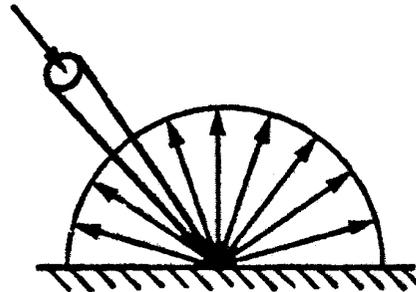
- *Should Optical Model based scattering replace this method? case in which the elastic coherent is dominant, like Graphite.*

Other scattering models ???

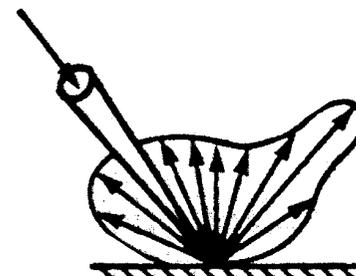
Duderstadt (1976): “Although **one introduces several horrifyingly brutal approximations**, at least to the solid state physicist, for nuclear engineer they are acceptable”.. This will mean that neutrons behave like light , or X-ray



Ideal mirroring reflection



Ideal diffuse reflection



Real surface reflection

How “horrifying” is the solid state physicists’ approximations?

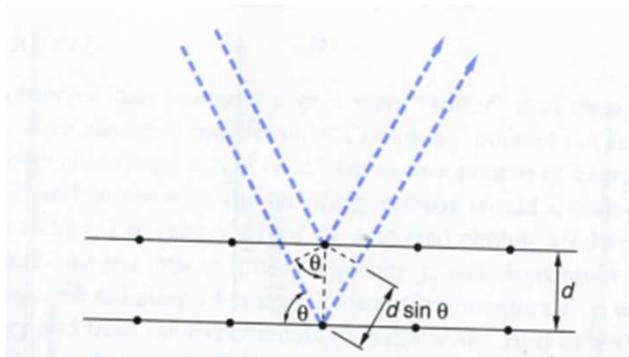
Two options for X ray based scattering

Bragg scattering:

$$2d \sin \theta = n\lambda$$

λ Wavelength

n an integer (depth of the layer)

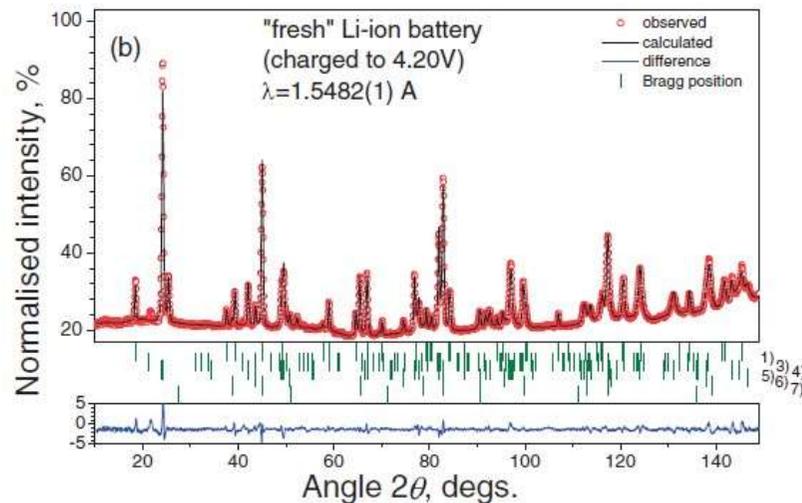
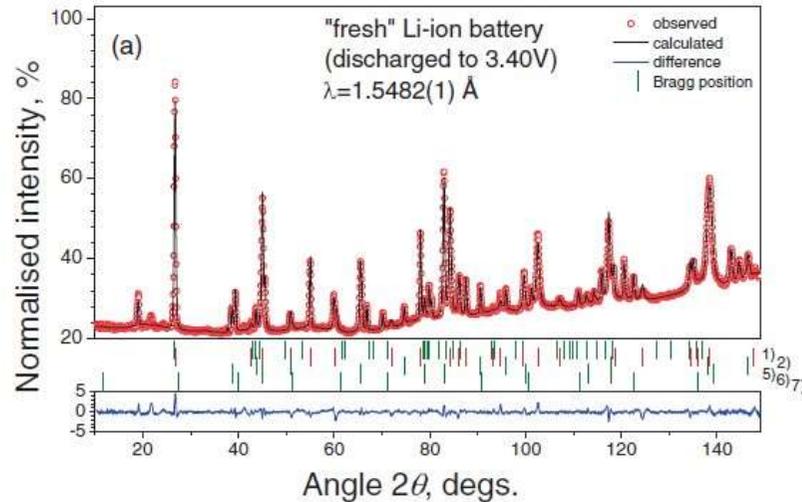


source : R. Tipler , R. Llewellyn “Modern Physics”

- Small angle Diffusive scattering kernel based on x-ray rough surface scattering
- “**assumingmolecular structure can be neglected...**”

Source: P. Müller-Buschbaum, Polymer Journal (2013) 45, 34–42

Scattering kernels for investigation Li-Ion Battery based on Bragg scattering (coherent- elastic)



- Calculated positions of Bragg peaks in charged and discharged state of Li-Ion Battery, when the anode is Graphite
- Inelastic-incoherent completely ignored

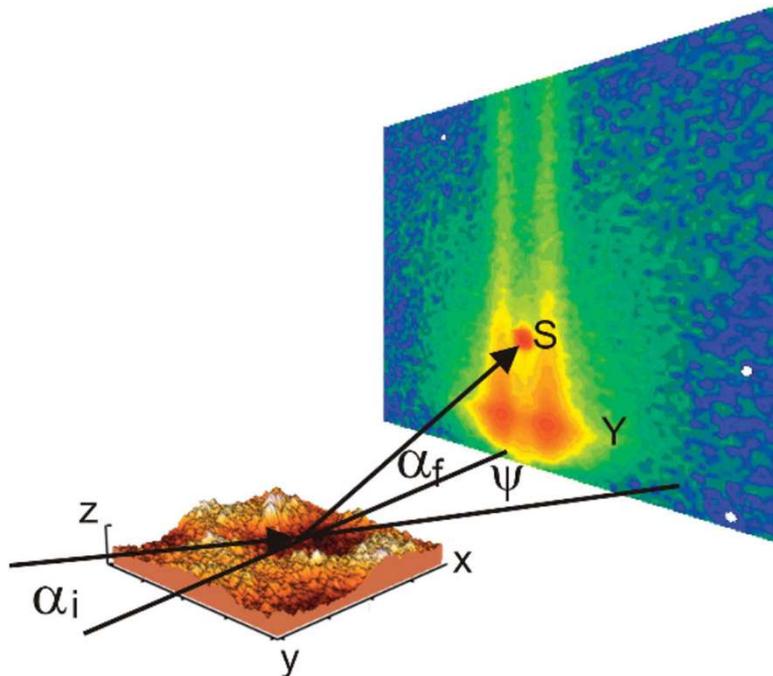
- Reference: Fatigue Process in Li-Ion Cells: An In Situ Combined Neutron Diffraction and Electrochemical Study, O. Dolotko et al., *Journal of The Electrochemical Society*, 159 (12) A2082-A2088 (2012)

$$2d \sin \theta = n\lambda$$

λ Wavelength

n an integer (depth of the layer)

Schematic of the scattering geometry used in GISANS (Grazing incidence small Angle neutron scattering)



Scattering vector $q = (q_x, q_y, q_z)$

$$q_x = 2\pi(\cos\Psi\cos\alpha_f - \cos\alpha_i) / \lambda$$

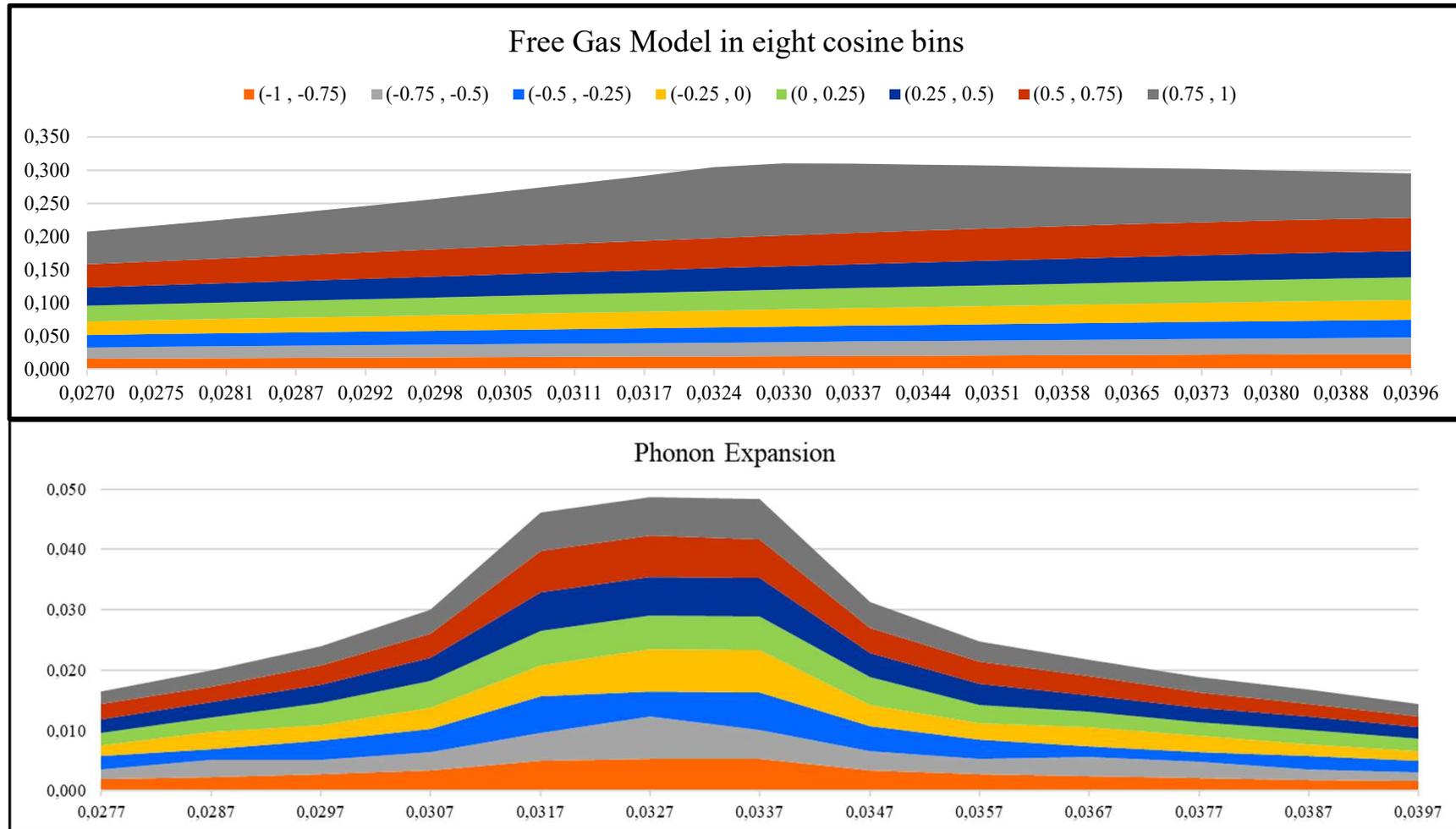
$$q_y = 2\pi(\sin\Psi\cos\alpha_f) / \lambda$$

$$q_z = 2\pi(\sin\alpha_i + \sin\alpha_f) / \lambda$$

- The sample is placed in the (x, y) plane, and the incident neutron beam is along the x axis.
- The resulting scattering pattern is anisotropic and typically exhibits a Yoneda peak (marked with Y) and a specular peak (marked with S)
- The scattering kernel approach is based upon [Fresnel transmission and reflection coefficients](#)
- Reference: Grazing incidence small-angle neutron scattering: challenges and possibilities,
[Source: P. Müller-Buschbaum, Polymer Journal \(2013\) 45, 34–42](#)

Coupling chemical binding and free gas model via the effective temperature: validity of the SCT approximation

Graphite: Comparison between SCT app. and phonon expansion at $E=0.0327$ eV ($n=12$) (via MATLAB)



Deriving the equation for agitating target: connection to the azimuth angle

$$\sigma_s^T(E \rightarrow E', \vec{\Omega} \rightarrow \vec{\Omega}') = \frac{1}{2\pi} \sigma_s^T(E \rightarrow E', \mu_0^{lab}) = \frac{1}{2\pi v} \left(\frac{A+1}{A} \right)^4 \left(\frac{A}{\pi} \right)^{3/2}$$

$$\int 2\pi u^2 du \int d\mu_u \int c^2 dc \int (u')^2 du' \int d\mu_{u'} \int \frac{2}{\sin \varphi} \frac{\delta(u'-u)}{(u')^2} \exp \left[v^2 - (A+1) \left(\frac{u^2}{A} + c^2 \right) \right]$$

$$\frac{4vv'c^2}{B_0'} \delta(\cos \varphi - \cos \hat{\varphi}) u \sigma_s(E_r) \frac{P(u, \mu_0^{cm})}{2\pi} d \cos \varphi$$

$$\frac{1}{uvc} \delta \left[\mu_u - \frac{(v^2 - c^2 - u^2)}{2uc} \right] \frac{1}{2u'ck_B T} \delta \left[\mu_{u'} - \frac{(v')^2 - (u')^2 - c^2}{2u'c} \right]$$

- The derivation of the equation is based on the fulfillment of all constraints marked by δ
- Note the azimuth angle is connected to the polar angle

$$(\vec{\Omega} \cdot \vec{\Omega}') = \mu_0^{lab} ; \varphi = \varphi_{u'} - \varphi_u$$

$$\cos \hat{\varphi} = R = \left[(4vv'c^2 \mu_0^{lab} - C_0') / B_0' \right]$$

- not incorporated in DBRC or $S(\alpha, \beta)$

Summary

- Thermal scattering analysis based on suggestions by J. Rowlands could help to better understanding of the different approaches
- More experiments in the thermal range for graphite based materials are under consideration.
- For thermal scattering: OMP can't comply with temperature, chemical binding or energy dependency, yet the idea of "optical model" is being used for Batteries and Photovoltaic cells.
- The azimuth angle, elastic coherent data should be considered for MC codes. In those cases the inelastic incoherent part could be presented by SCT approximation.
- **Response to Duderstadt (1976): The solid state physicists' approximations are " quite horrifying" as well.**